

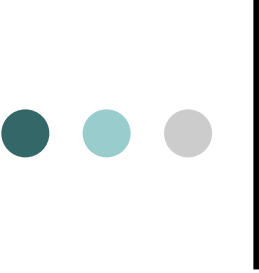


Shell model method for Gamow-Teller transitions in deformed nuclei

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ND, Aug 19-21



β -decay & electron-capture in stars

- Stellar weak-interaction rates are important for studying astrophysical problems
 - for nucleosynthesis calculations
 - for core collapse supernova modeling
- Evaluation of nuclear matrix element
 - essentially a nuclear structure problem
- Theoretical models
 - Independent particle model (FFN)
 - Spherical shell model (SM)
 - Quasiparticle random phase approximation (QRPA)

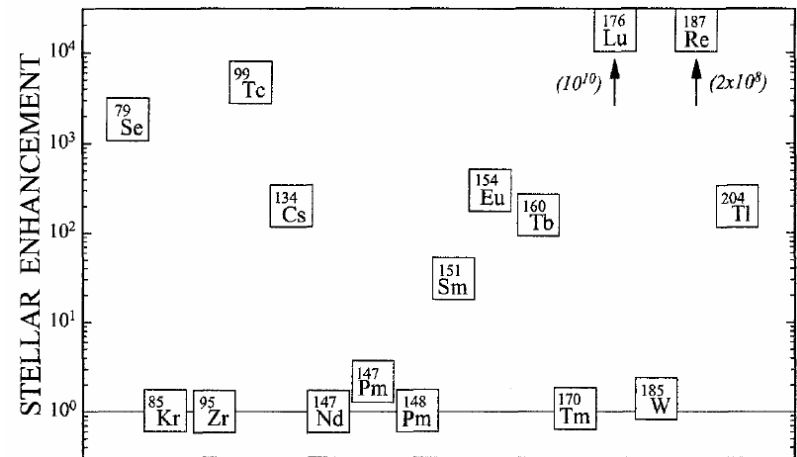
Stellar enhancement of β -decay rate

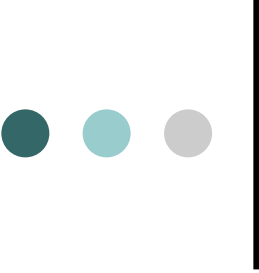
- A stellar enhancement can occur due to thermal population of excited states

$$\lambda_{\beta} = \sum_i \left(p_i \times \sum_j \lambda_{\beta ij} \right)$$
$$p_i = \frac{(2I_i + 1) \times \exp(-E_i / kT)}{\sum_m (2I_m + 1) \times \exp(-E_m / kT)}$$

- Examples in the s-process

F. Kaeppler,
Prog. Part. Nucl. Phys. 43 (1999) 419



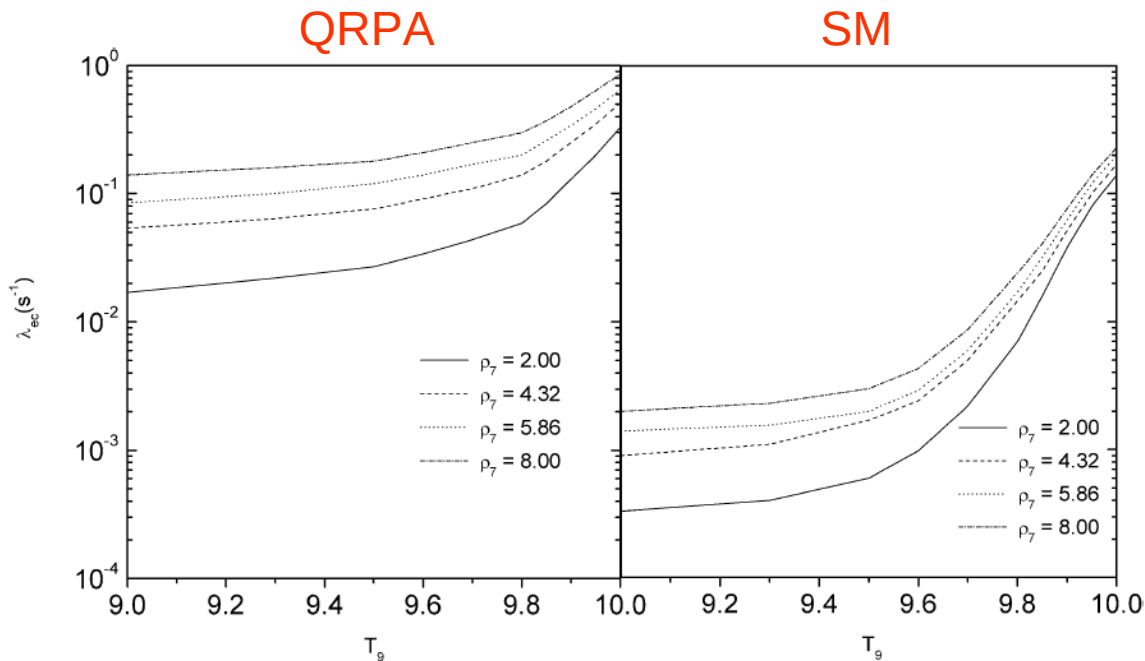


Electron-capture in core collapse supernovae

- Electron-capture rates are needed in simulations of core collapse supernovae
 - In prior simulations, electron capture has been treated in a highly parametrized fashion
 - With realistic treatment of electron capture on heavy nuclei, significant changes in the hydrodynamics of core collapse are found
 - W.R. Hix *et al.*, *Phys. Rev. Lett.* 91 (2003) 201102
- ^{55}Co was ranked as the most important nuclei with respect to the importance for the electron capture process
 - M.B. Aufderheide *et al.*, *Astrophys. J. Suppl.* 91 (1994) 389

Supernova electron capture rates on ^{55}Co by SM and QRPA

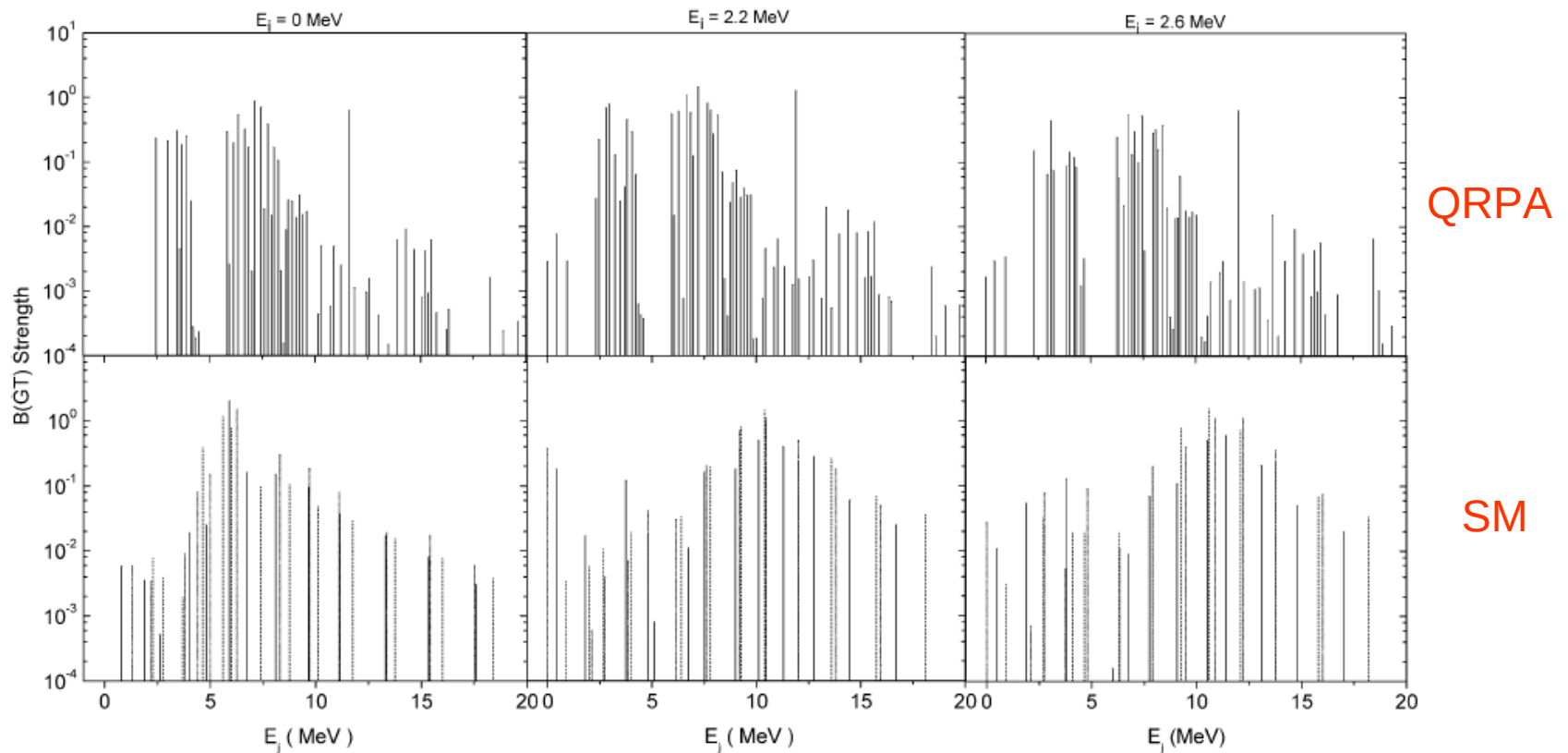
- QRPA rates by two orders of magnitude stronger than SM
 - K. Langanke & G. Martinez-Pinedo, *Phys. Lett. B* 436 (1998) 19
 - J.-U. Nabi & M.-U. Rahman, *Phys. Lett. B* 612 (2005) 190

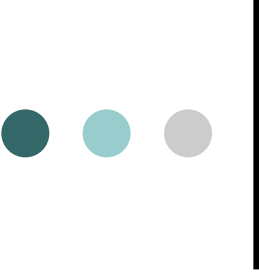


This can have a significant astrophysics impact because rate change of lepton-to-baryon ratio changes by 50% due to electron capture on ^{55}Co

B(GT) distribution in SM and QRPA

- G-T strength more fragmented in RQPA than in SM





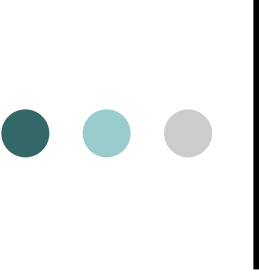
Comparison of these models

○ SM:

- K. Langanke & G. Martinez-Pinedo, *Phys. Lett. B* 436 (1998) 19
- Small model space in one major pf shell, allowing only a few particles excited from $f_{7/2}$, renormalization introduced
- For electron-capture rates at finite temperature, G-T strength is not calculated, but the Brink-hypothesis is used
- Deficiency: insufficient model space

○ QRPA:

- J.-U. Nabi & M.-U. Rahman, *Phys. Lett. B* 612 (2005) 190
- Calculated nuclear states are not angular momentum states, but K -states (which are mixtures of I -states)
- Deficiency: not a shell model, may contain spurious in the wavefunction



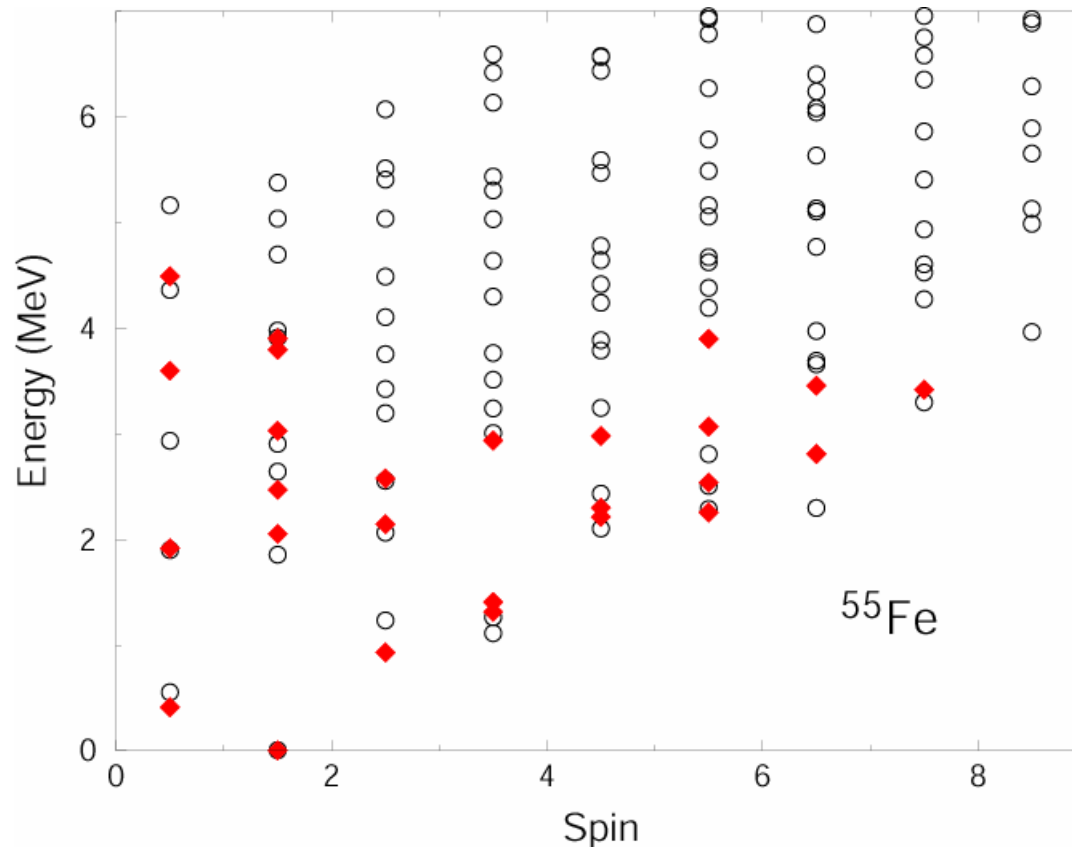
The projected shell model (PSM)

- Take a set of deformed (quasi)particle states (e.g. solutions of HF, HFB or HF + BCS)
- Build up configurations (qp vacuum + multi-qp states)
- Project them onto good angular momentum (if necessary, also parity, particle number) to form a basis in lab frame
- If necessary, superimpose configurations belonging to different qp representations (the GCM-concept)
- Diagonalize a two-body Hamiltonian in projected basis

Hara & Sun, *Int. J. Mod. Phys. E*4 (1995) 637

Energy spectrum for ^{55}Fe : validating the PSM application

- Projected shell model calculation for ^{55}Fe energy levels



Nuclear matrix elements in the projected basis

- Gamow-Teller rate $B(GT) = \frac{2I_f + 1}{2I_i + 1} \langle \psi_{I_f} | \hat{\beta}^\pm | \psi_{I_i} \rangle^2$
- PSM wavefunction $\psi_M^I = \sum_K f_K \hat{P}_{MK_K}^I | \phi_K \rangle$
- e-e system $| \phi_e(\varepsilon_e) \rangle \Rightarrow \{ | \varepsilon_e \rangle, b_\nu^\dagger b_\nu^\dagger | \varepsilon_e \rangle, b_\pi^\dagger b_\pi^\dagger | \varepsilon_e \rangle, b_\nu^\dagger b_\nu^\dagger b_\pi^\dagger b_\pi^\dagger | \varepsilon_e \rangle, \dots \}$
- o-o system $| \phi_o(\varepsilon_o) \rangle \Rightarrow \{ a_\nu^\dagger a_\pi^\dagger | \varepsilon_o \rangle, a_\nu^\dagger a_\nu^\dagger a_\nu^\dagger a_\pi^\dagger | \varepsilon_o \rangle, a_\nu^\dagger a_\pi^\dagger a_\pi^\dagger a_\pi^\dagger | \varepsilon_o \rangle, \dots \}$
- Overlapping matrix element (Gao, Sun, Chen, *PRC* 74 (2006) 054303).

$$\langle \phi_o(\varepsilon_o) | \hat{O} \hat{P}_{K_o K_e}^I | \phi_e(\varepsilon_e) \rangle \sim \int d\Omega D_{K_o K_e}^I(\Omega) \langle \phi_o(\varepsilon_o) | \hat{O} \hat{R}(\Omega) | \phi_e(\varepsilon_e) \rangle$$

The interactions

- Total Hamiltonian $\hat{H} = \hat{H}_0 + \hat{H}_{QP} + \hat{H}_{GT}$

- Quadrupole + monopole-pairing + quadrupole-pairing

$$\hat{H}_{QP} = -\frac{1}{2}\chi_{QQ} \sum_{\mu} \hat{Q}_{2\mu}^{\dagger} \hat{Q}_{2\mu} - G_M \hat{P}^{\dagger} \hat{P} - G_Q \sum_{\mu} \hat{P}_{2\mu}^{\dagger} \hat{P}_{2\mu}$$

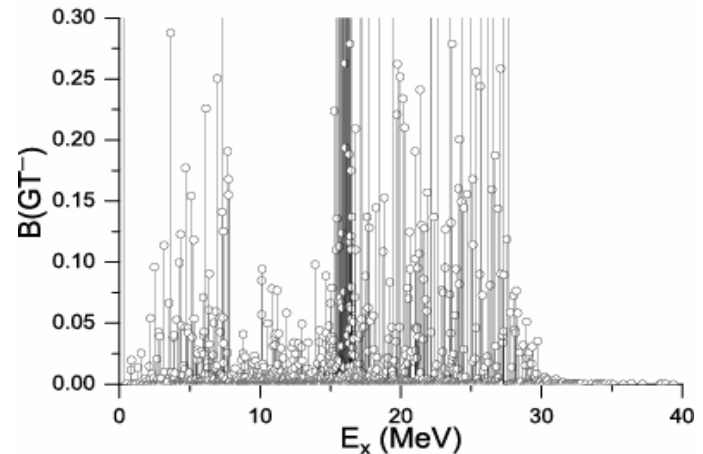
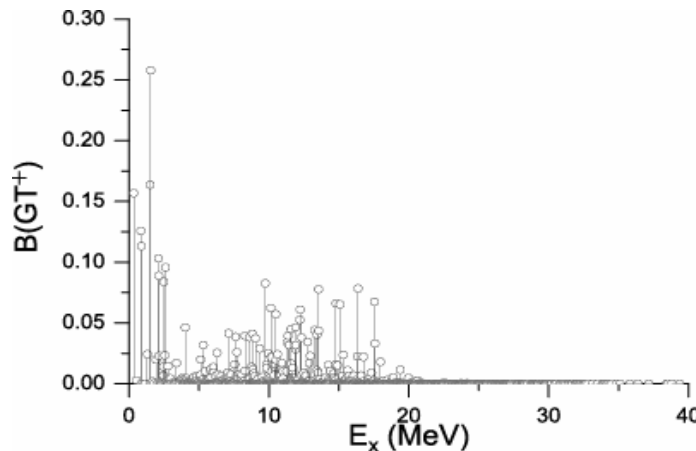
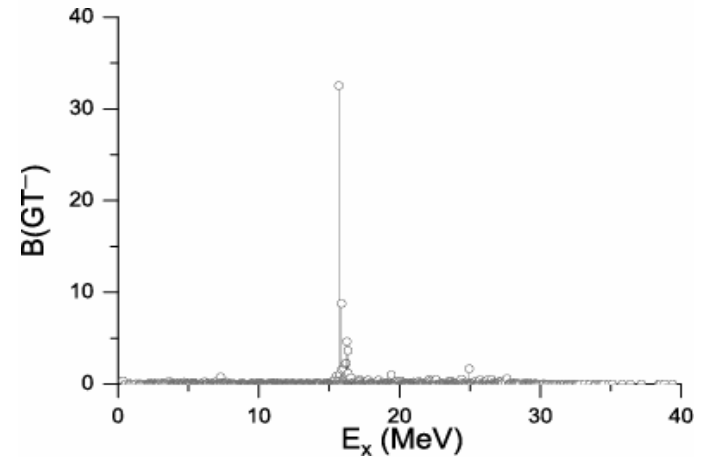
- Charge-exchange (Gamow-Teller)

$$\hat{H}_{GT} = +2\chi_{GT} \sum_{\mu} \hat{\beta}_{1\mu}^{-} (-1)^{\mu} \hat{\beta}_{1-\mu}^{+} - 2\kappa_{GT} \sum_{\mu} \hat{\Gamma}_{1\mu}^{-} (-1)^{\mu} \hat{\Gamma}_{1-\mu}^{+}$$

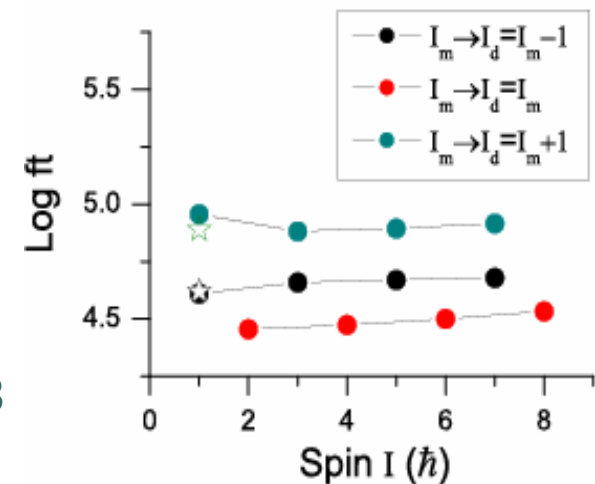
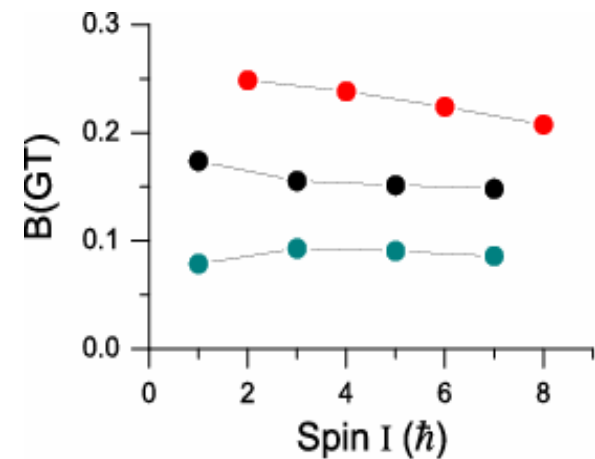
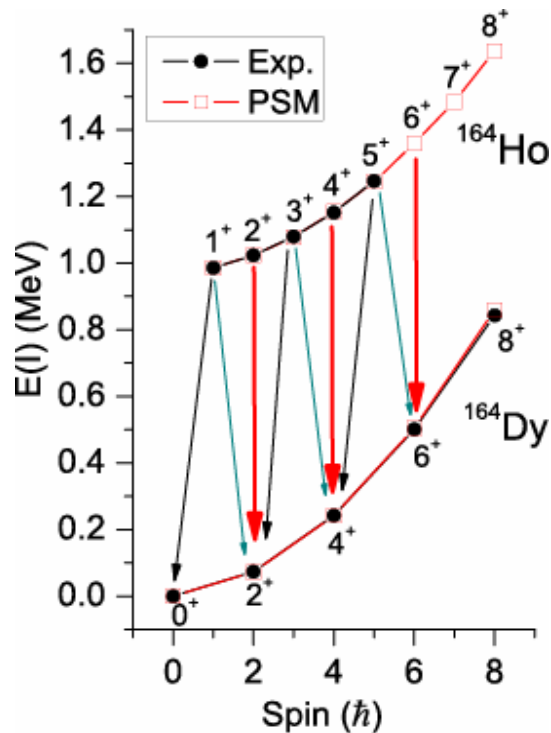
- Kuz'min & Soloviev, *Nucl. Phys. A* 486 (1988) 118

B(GT) distribution

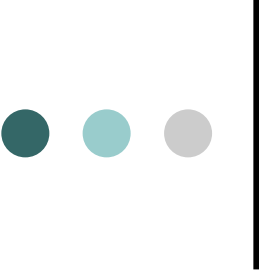
- Initial state: ground state of even-even nucleus
- Final states: all 1^+ states in odd-odd nucleus
- Ikeda sum-rule fulfilled



B(GT) and $\log ft$ in $^{164}\text{Ho} \rightarrow ^{164}\text{Dy}$

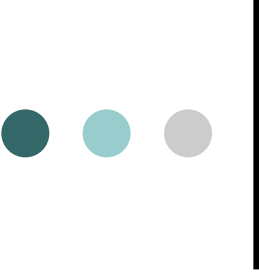


Gao, Sun, Chen, *Phys. Rev. C* 74 (2006) 054303



B(GT) in Projected Shell Model

- PSM wave functions contain correlations beyond mean-field and are written in the laboratory frame, having definite quantum numbers
- A state-by-state evaluation of B(GT) is feasible (no Brink-hypothesis needed)
 - The PSM single-particle state is big, but dimension of the shell model space is small (usually in the range of 10^2 – 10^4)
- PSM is a multishell shell model and can calculate forbidden transitions
 - Calculation of forbidden transitions is not possible for most of conventional shell models working in one-major shell basis



Conclusion

- Angular momentum projection is an efficient way of truncating shell model space to perform shell model calculations for heavy, deformed nuclei.
- The projection technique is well developed. Projected shell model is a practical example.
- A method of using projected wave functions to calculate weak interaction rates is developed, and can be applied to nuclear astrophysics.
- The method may overcome the deficiencies in SM and QRPA