

# Exposing the Nuclear Burning Ashes of Radius Expansion Type I X-ray Bursts

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# Punch-line



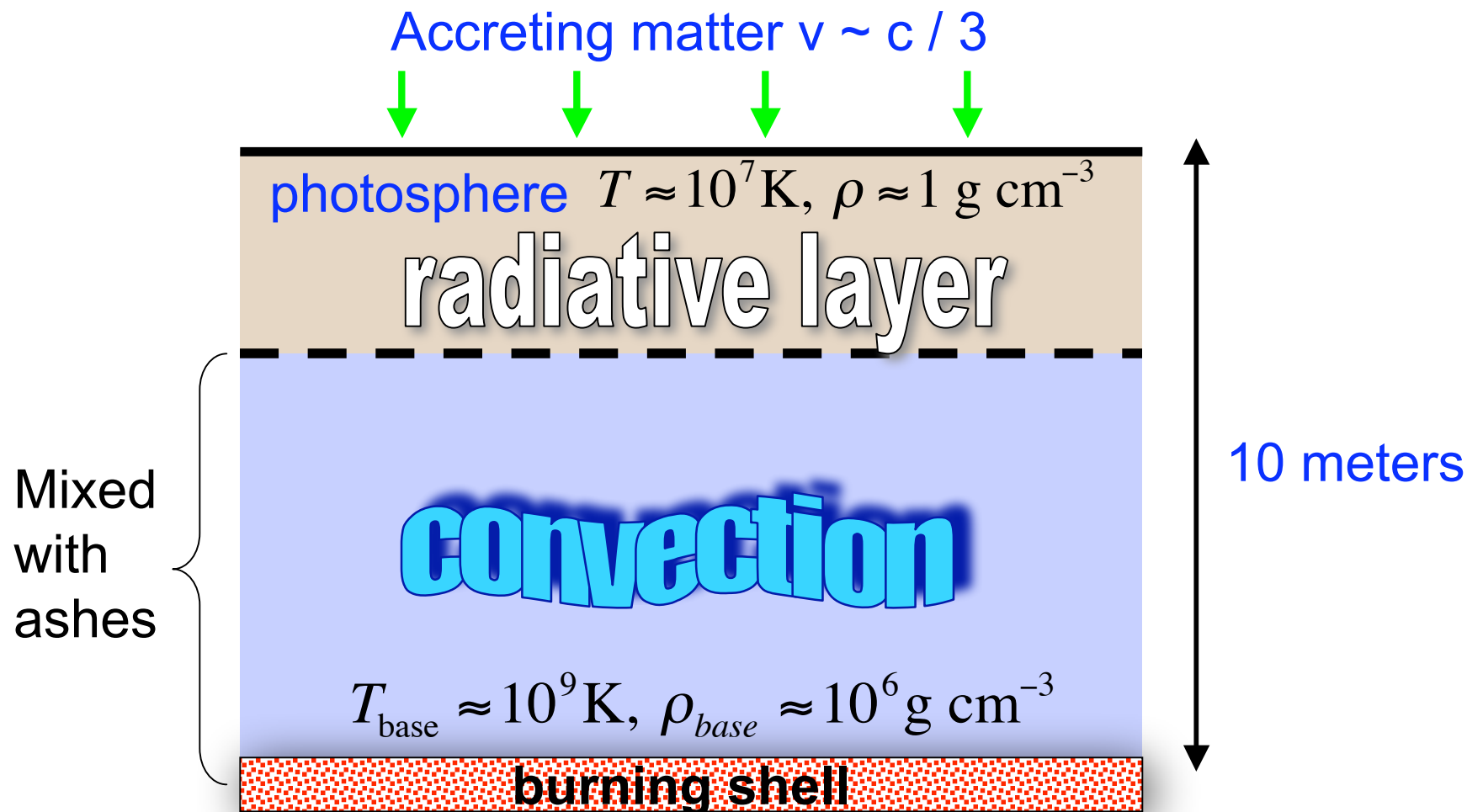
During radius expansion X-ray bursts ( $L > L_{\text{eddington}}$ ) **ashes** of nuclear burning can be **ejected** by the radiative wind.

Heavy element ashes not fully ionized  Detect spectral lines of ashes

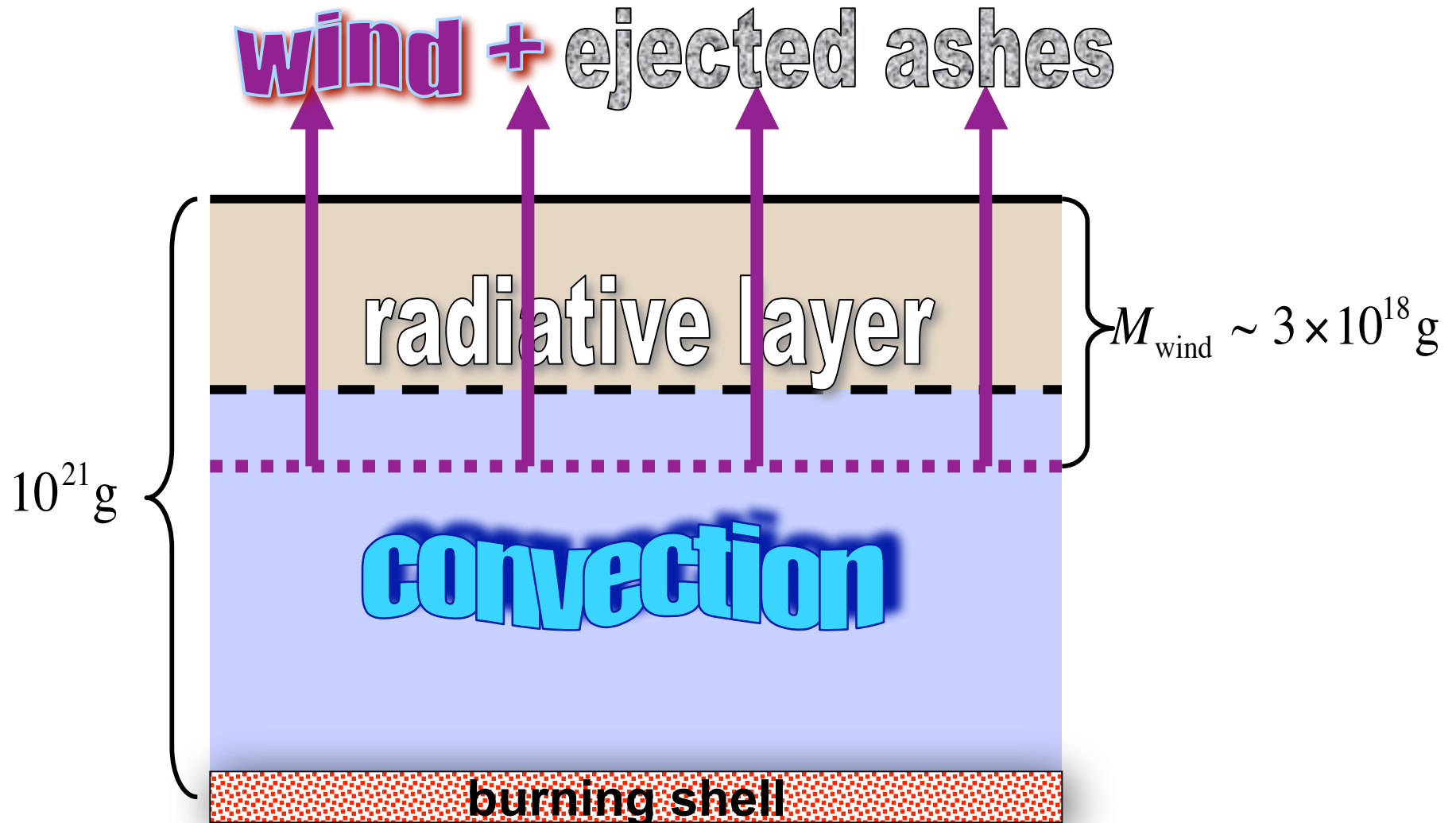
We find that the spectral features of ashes should be detectable with **current high resolution X-ray telescopes** (e.g., XMM, Chandra, ~~Astro E2~~):

1. In the **wind** during the radius expansion phase and
2. At the **neutron star surface** during the burst decay/cooling.

# Neutron star atmosphere during a burst



# Atmosphere during a radius expansion burst



# Vertical extent of convective zone

During burst rise, convective zone reaches out to pressures where the entropy matches that at the base of the burning layer.

**In convective zone:**  $T \sim T_{\text{base}}(P/P_{\text{base}})^{2/5}$   
**In radiative zone:**  $T \propto P^{1/4}$

} Convection stops near intersection of the two profiles

How far out convection reaches depends on:

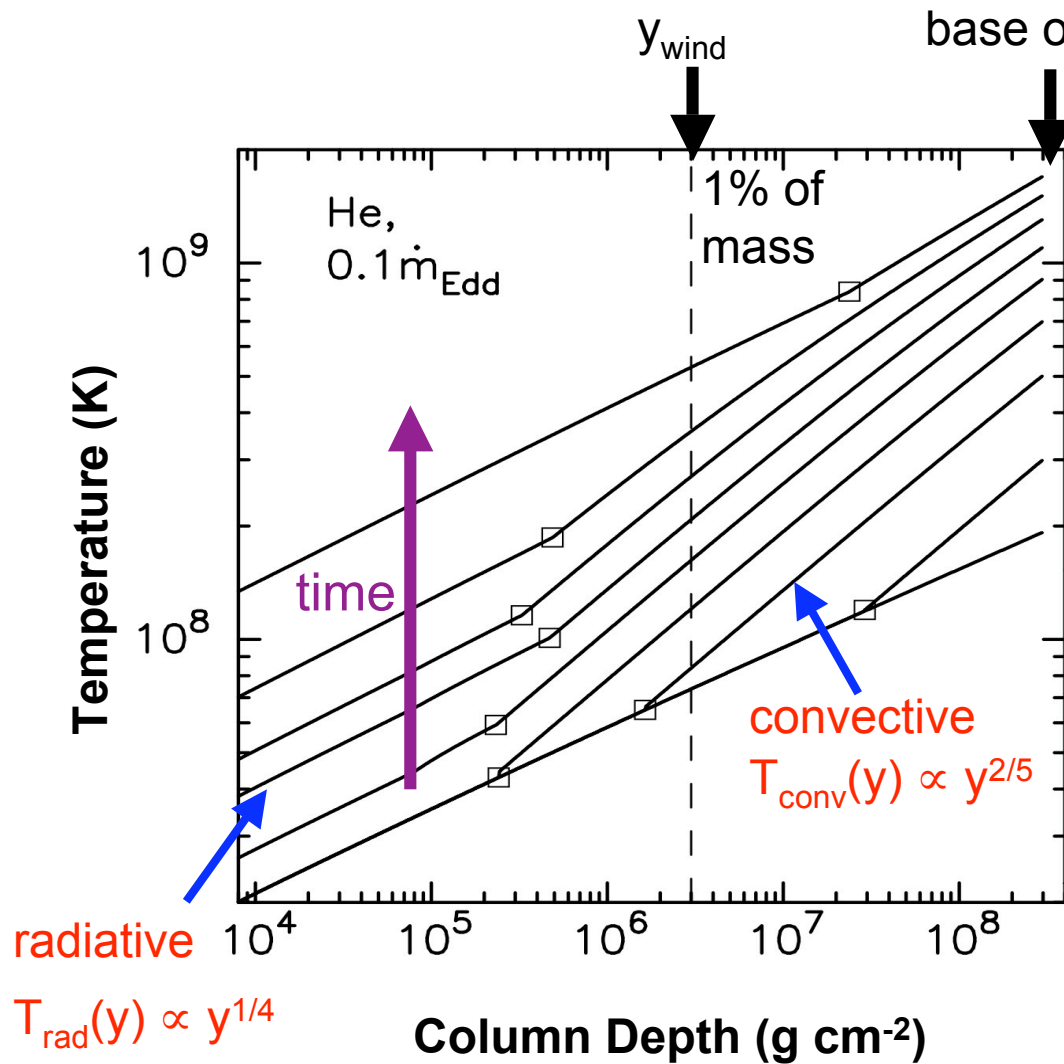
1. The **thermal state of the outer envelope prior to burst** ---the colder it is, the farther out convection reaches.
2. The **rate at which the radiative region heats up** due to radiation escaping out the top of the convective zone:

$$F_{\text{loss}} = \frac{4acT^4}{3\kappa y} \left( \frac{d \ln T}{d \ln P} \right)_{ad}, \text{ and a heating time } t_{\text{th}} = \frac{c_p y T}{F_{\text{loss}}}$$

3. The **rate at which the temperature at the base of the convective zone grows** due to nuclear burning:  $t_{\text{gr}} \approx \frac{c_p T_b}{\epsilon}$

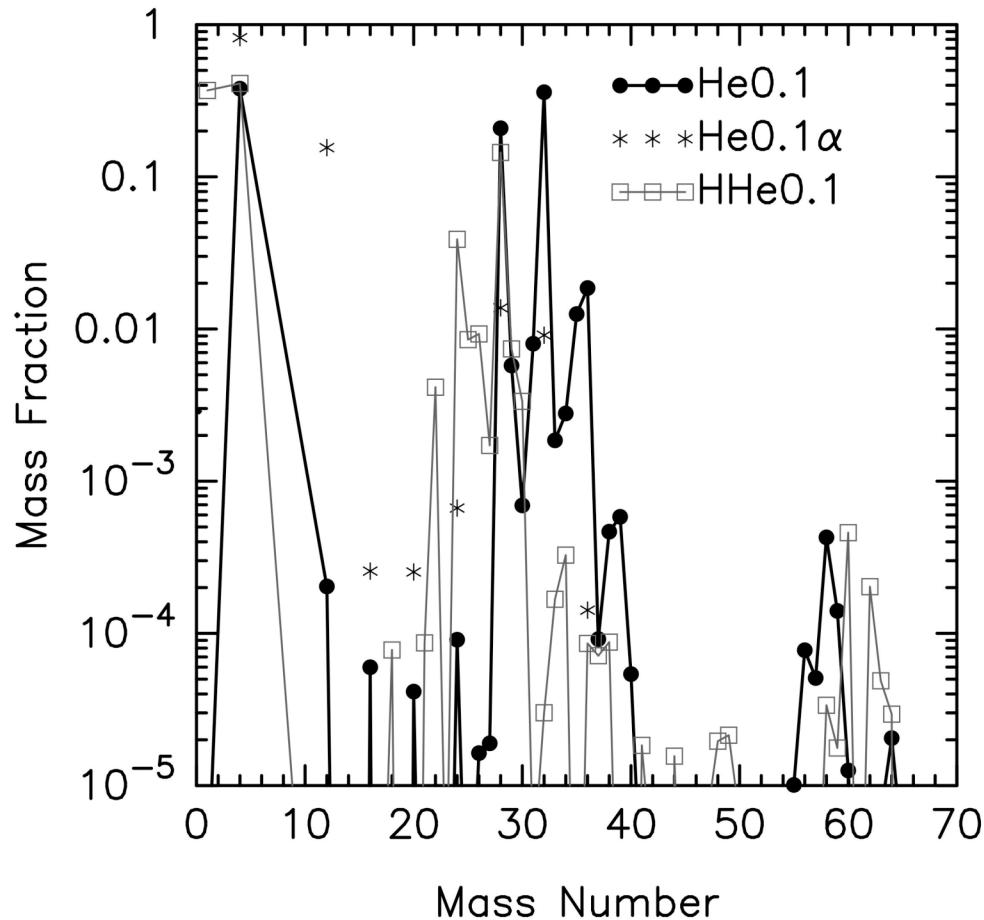
We calculate energy generation rate  $\epsilon$  and nucleosynthesis using a full nuclear reaction network.

# Evolution of temperature profile



- Convective zone extends beyond 1% pressure coordinate => ashes mixed in with wind ejecta
- Explored burst parameter space: E.g., accretion rate, composition of accreted material, column depth of base, crustal heat flux.
- For pure He accretors or low accretion rate ( $\sim 0.01 \text{ Edd}$ ) H/He accretors,  $y_{\text{conv,min}} < y_{\text{wind}}$ .

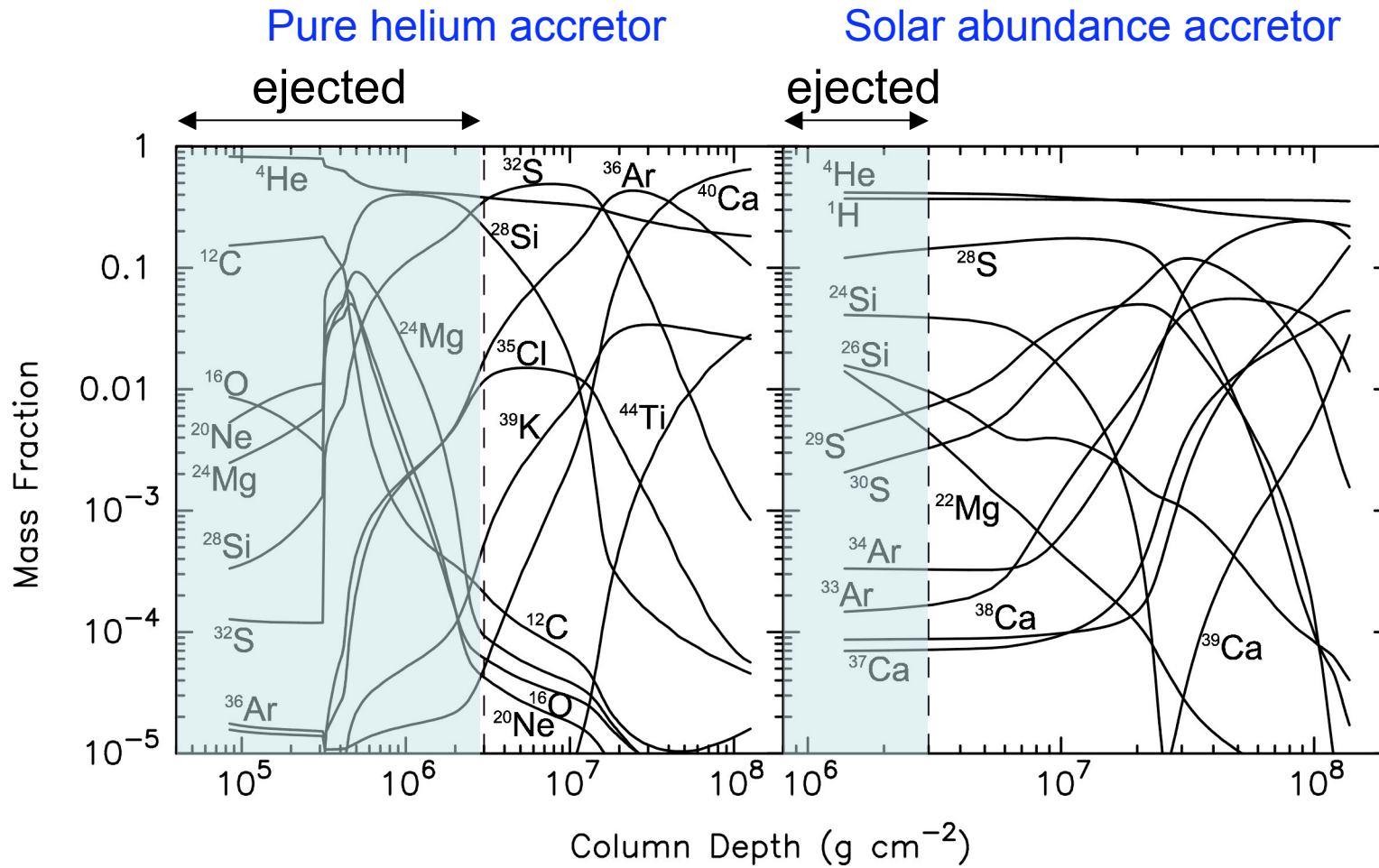
# Composition of ejected ashes



- **Si and S** are the most abundant ejected ashes.
- In burning layers that are initially pure He, protons from ( $\alpha$ , p) reactions enable the rapid onset of the  $^{12}\text{C}(p, \gamma)^{13}\text{N}(\alpha, p)^{16}\text{O}$  reaction sequence.
- This sequence **bypasses** the relatively slow  $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$  reaction and leads to a **spike in energy generation and broadening of abundance pattern beyond  $\alpha$ -nuclei**.

**Aside:** Spike should also be **directly observable as a sharp rise in flux** during burst rise.

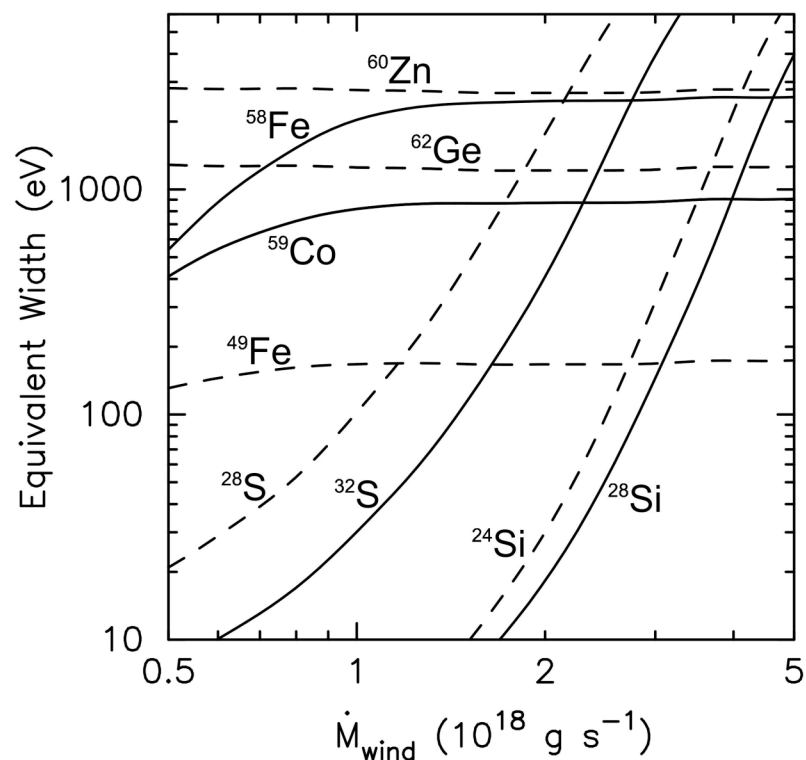
# Composition profile of ashes



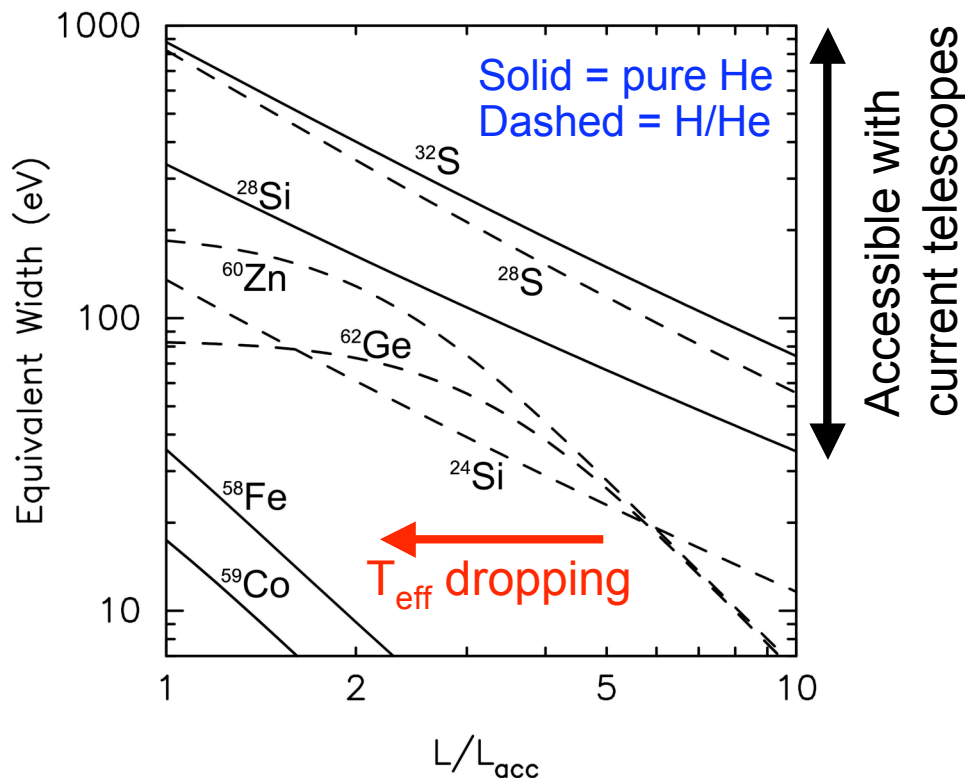


# Detecting the ashes

Photoionization edges in **winds** ...and in the **cooling photosphere**



**wind mass loss rate**



**cooling luminosity relative to accretion luminosity**

- These ejected nuclei imprint photo-ionization edges in wind out at radii of 100 - 1000 km.

- Once the photosphere recedes back to  $R_{\text{NS}}$ , exposed ashes may yield lines useful for redshift measurement.

# Summary

- Calculations of the convective zone evolution suggest that **ashes of nuclear burning are ejected by the radiative wind during radius expansion bursts** in systems accreting pure He or mixed H/He at low accretion rates.
- Large abundance of ashes with mass number  **$A = 30 - 60$** .
- The column density of ashes in hydrogen-like states is high enough that **current high resolution X-ray telescopes** should be able to **detect spectral features of ashes**:
  1. in the **wind** and,
  2. in the **photosphere at  $R_{NS}$**  during burst cooling.
- Detecting spectral features would **probe nuclear burning processes** during bursts and may enable a **measurement of the neutron star gravitational redshift**.