

# Systematics of strength functions of isoscalar dipole and monopole modes

J. Terasaki (Univ. of North Carolina at Chapel Hill)

- Quasiparticle random-phase approximation
- Systematics of strength functions: IS  $1^-$  Ni and  $0^+$  Sn
- Incompressibility of nuclear matter

July. 14, 2005, Notre Dame

## Quasiparticle random phase approximation (QRPA)

What approximation?

1. Excitations are represented by linear combinations of 2-quasiparticle creation and annihilations (if there is no pairing correlations, 1p-1h creation and annihilation).
2. Good for collective vibrations (the more harmonic, the better).

## Scheme of solution

HFB ← Dobaczewski et al. Nucl. Phys. A422 (1984) 103  
q.p. wave functions are obtained.

QRPA ← matrix formulation

Cf. e.g. D.J. Rowe, Nuclear Collective Motion

**HFB**



← self-consistent mean field and pair field  
w.f. of canonical basis, uv factors  
two-body interaction

**QRPA**

**Self-consistent  
Spherical symmetry**

J.T. et al., P.R.C 71, 034310 (2005)

## Interaction used

Skyrme interaction for particle-hole matrix elements

SkM\*, SLy4, SkP, SkO'

For particle-particle and hole-hole matrix elements

Volume-type pairing interaction  
(no density dependence)

## Definitions of the transition operators used ( $J^\pi=0^+$ , IS)

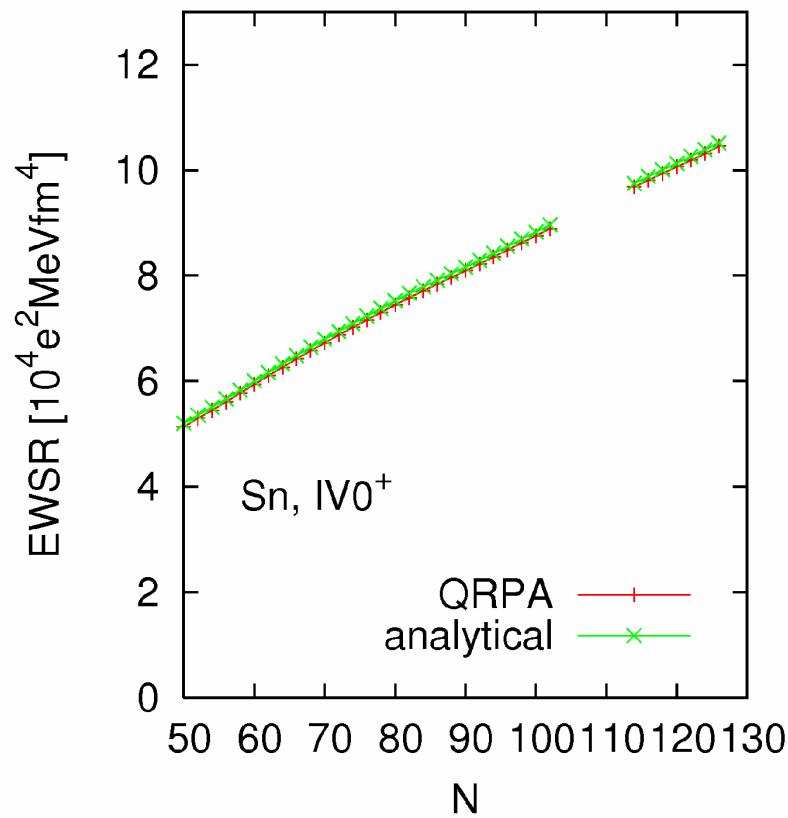
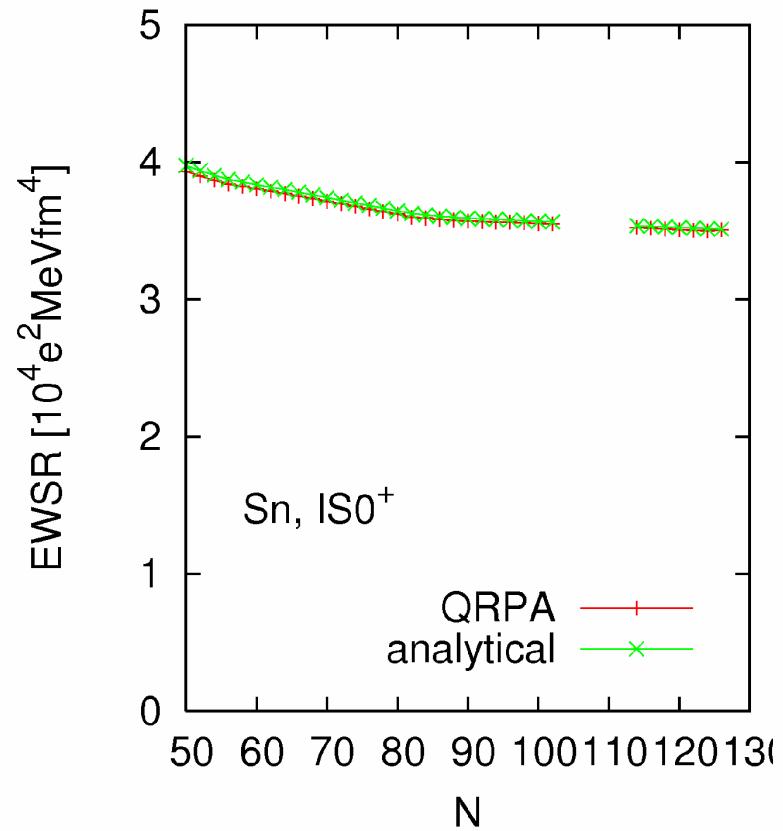
$$F_{00}^{IS} = \frac{eZ}{A} \sum_{i=1}^A \dot{\mathbf{a}} r_i^2$$

used by G.F. Bertsch, S. Kamerdzhiev, S. Sagawa (not recent ones) ...

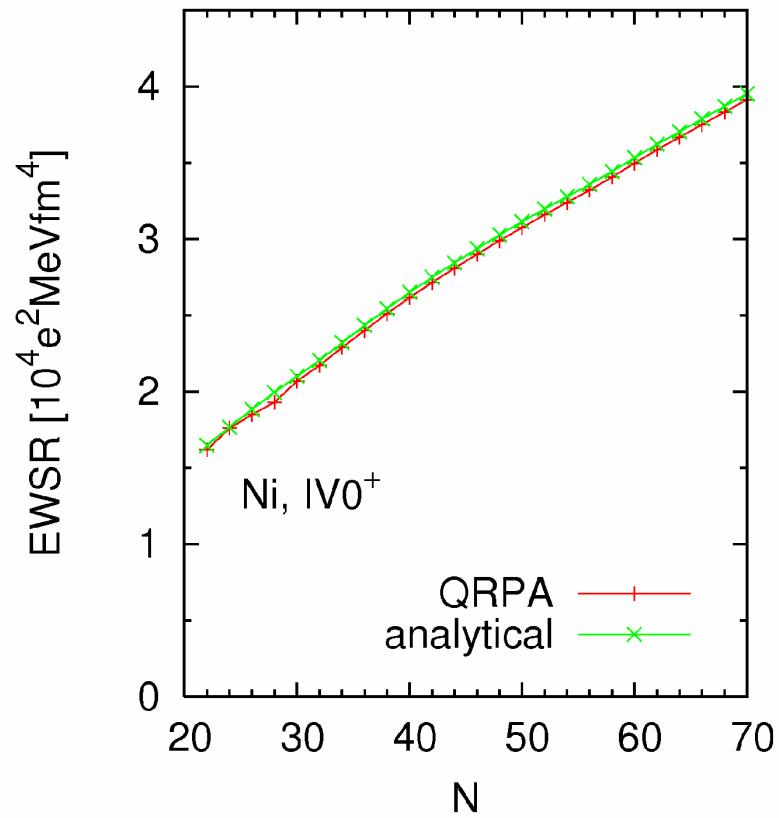
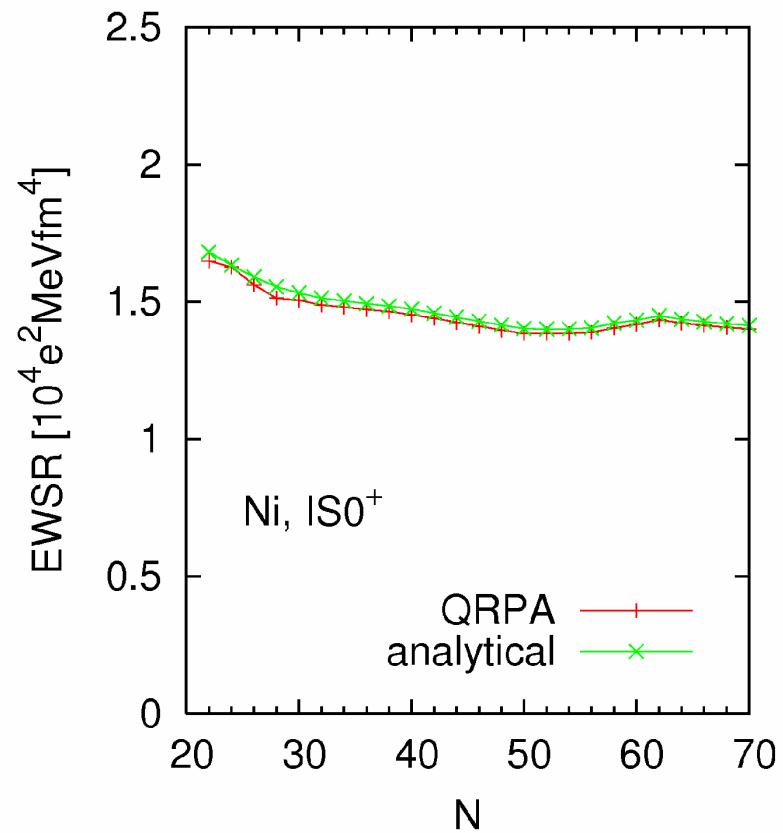
## Energy-weighted sum rule (EWSR)

$$\sum_k \dot{\mathbf{a}} E_k | \langle k | F_{00}^{IS} | 0 \rangle |^2 = 2 \frac{e^2 \hbar^2}{m} \frac{Z^2}{A} \langle \mathbf{r}^2 \rangle$$

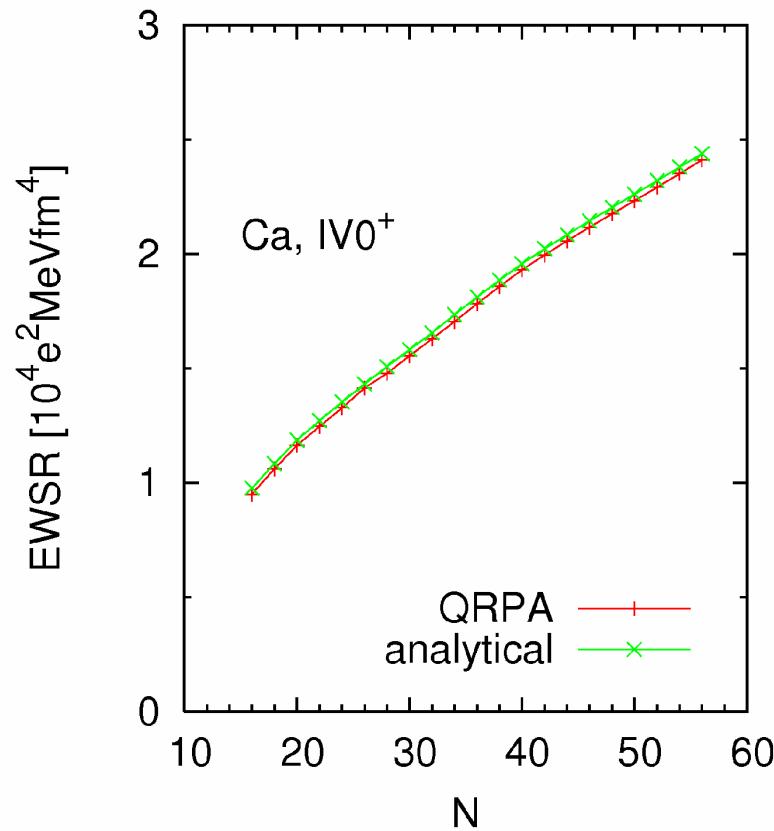
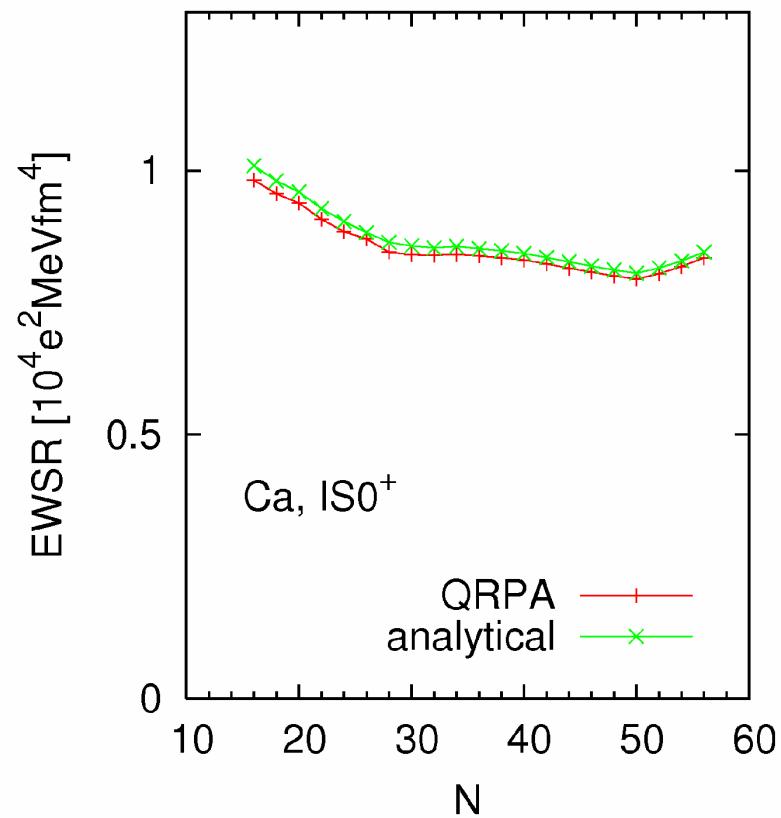
## Energy-weighted sum rule, $0^+$ (SkM\*)



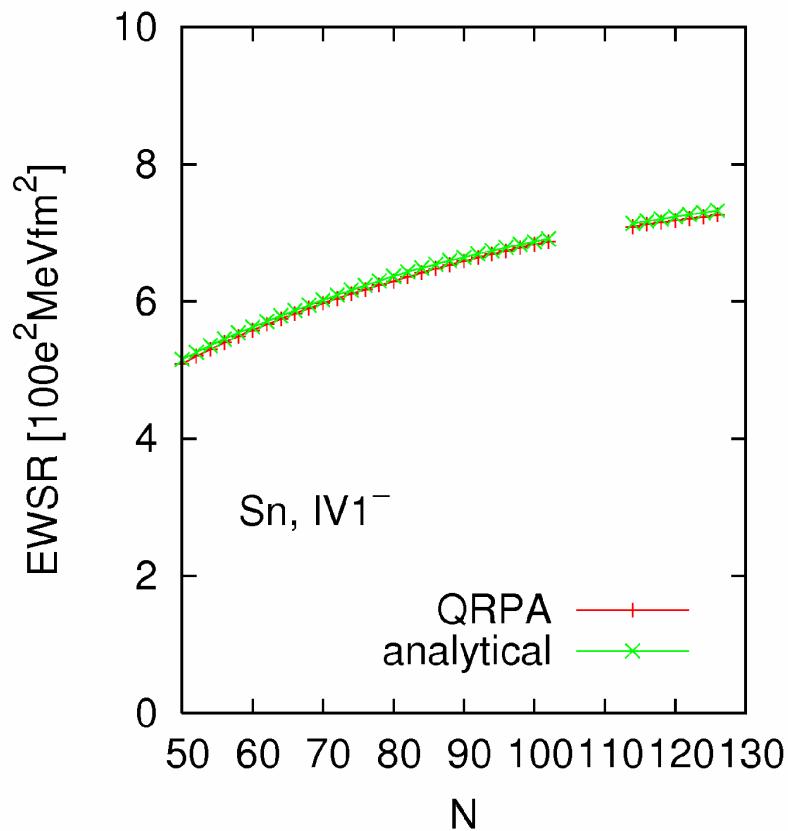
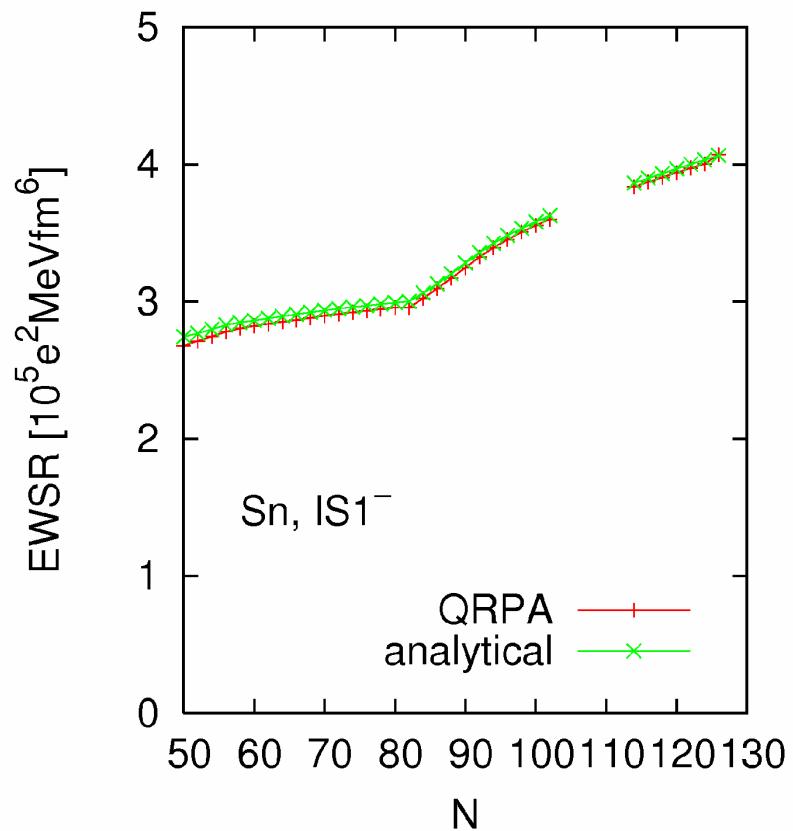
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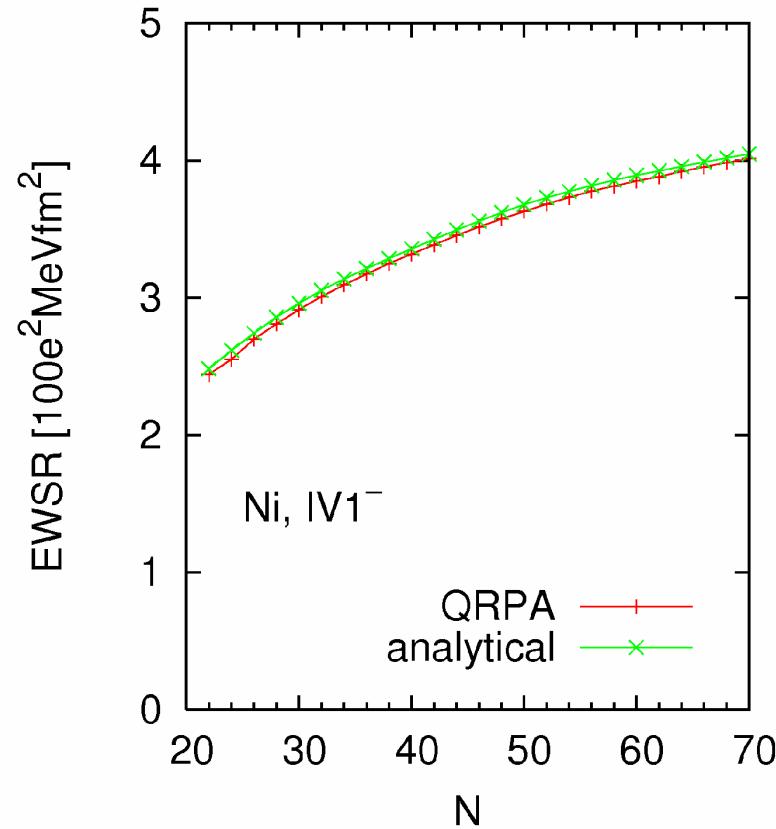
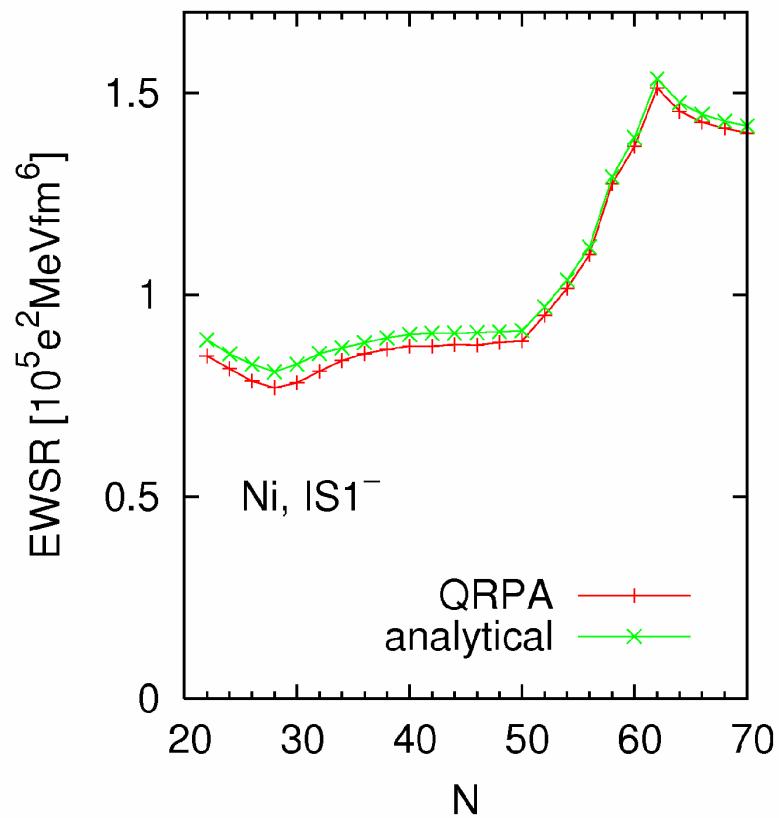
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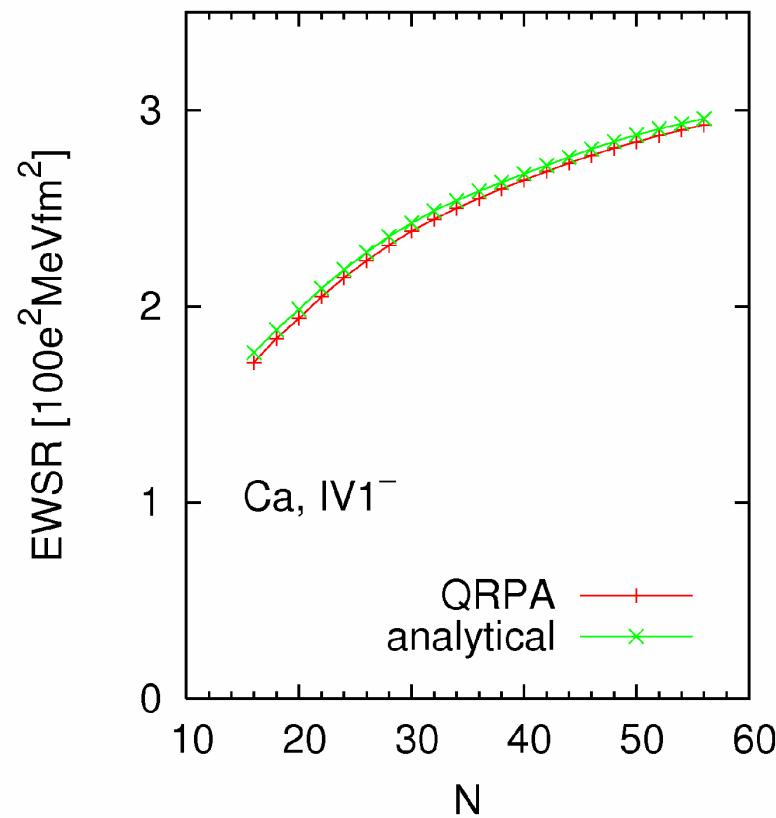
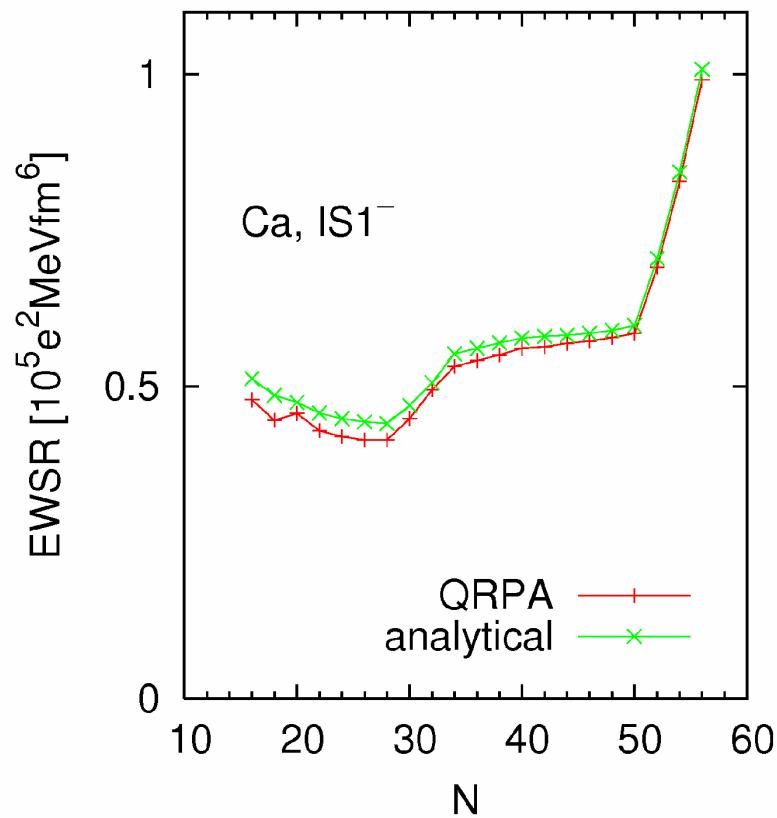
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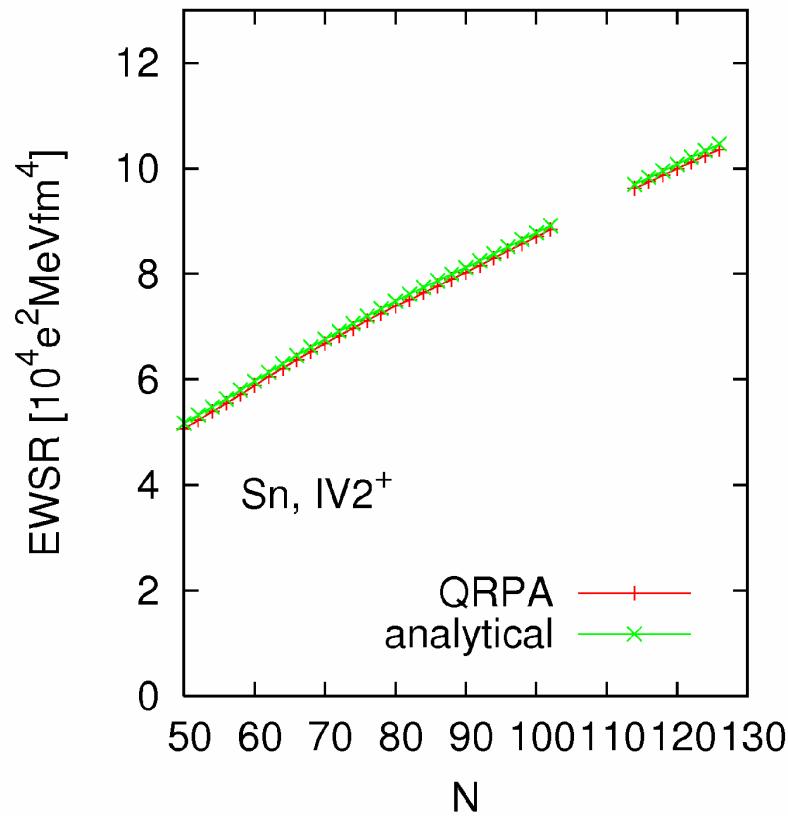
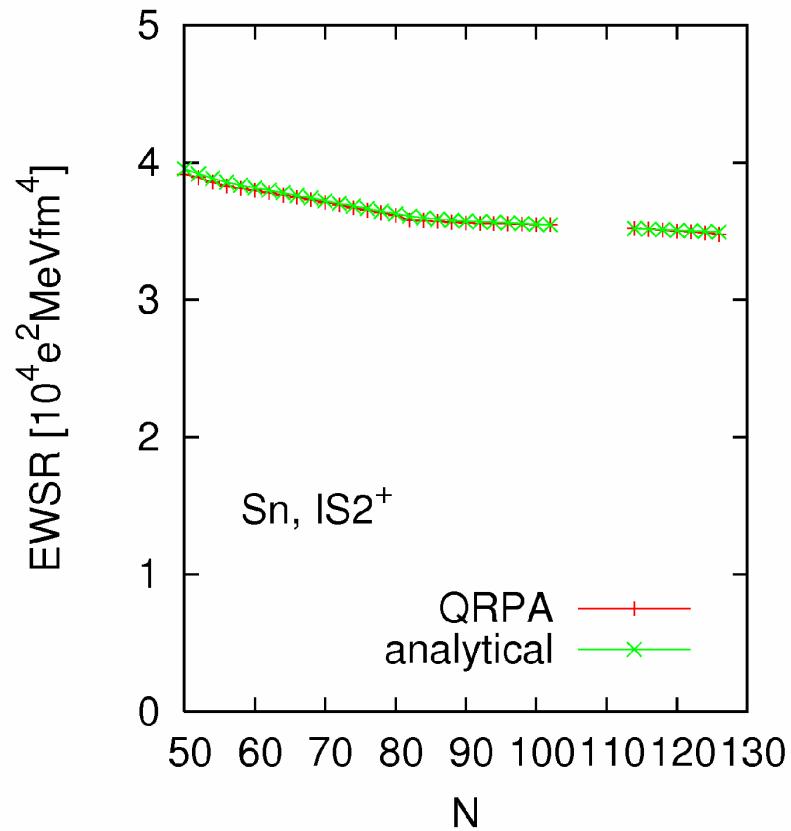
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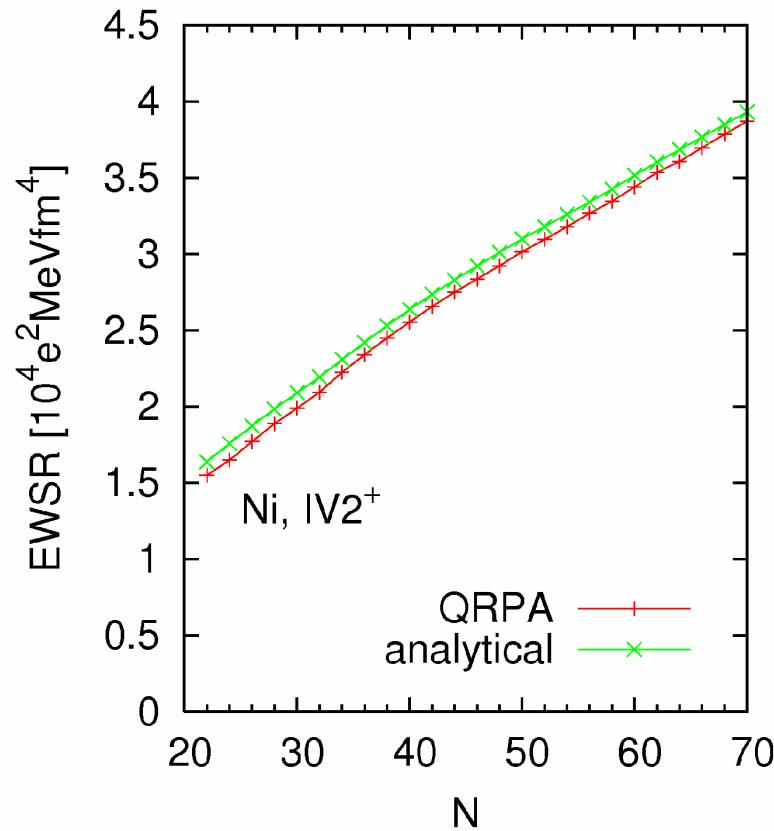
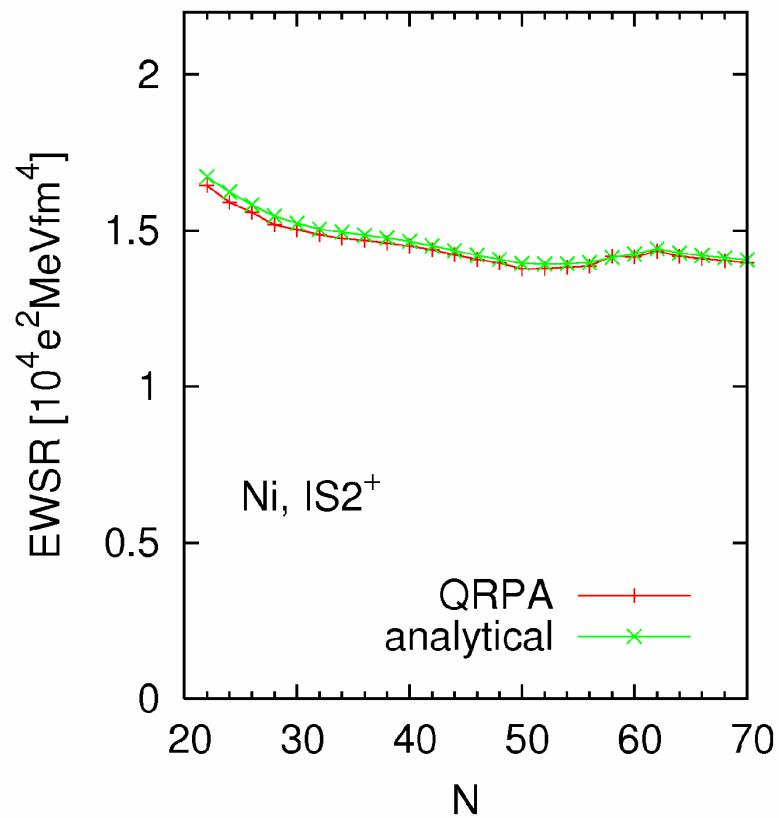


## Energy-weighted sum rule, $2^+$ (SkM\*)

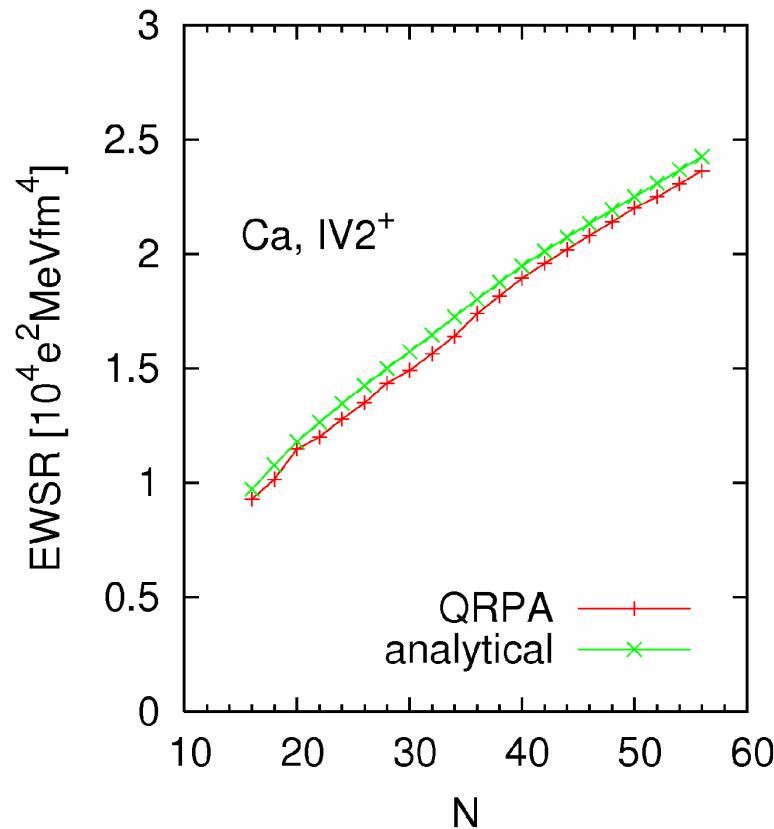
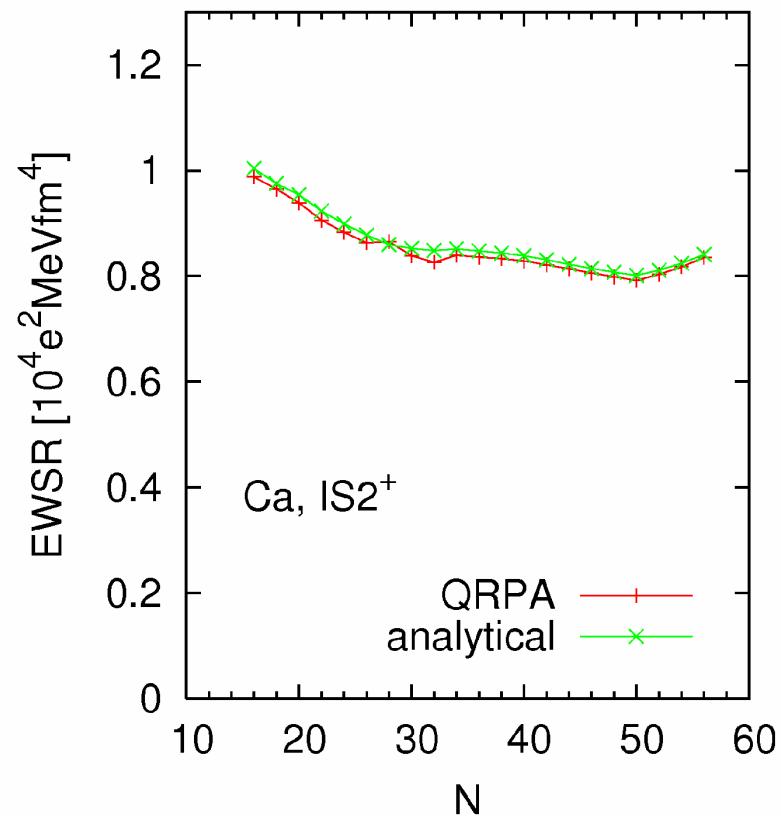


$$\text{EWSR}(0^+)/\text{EWSR}(2^+) = \frac{8p}{25} @ 1$$

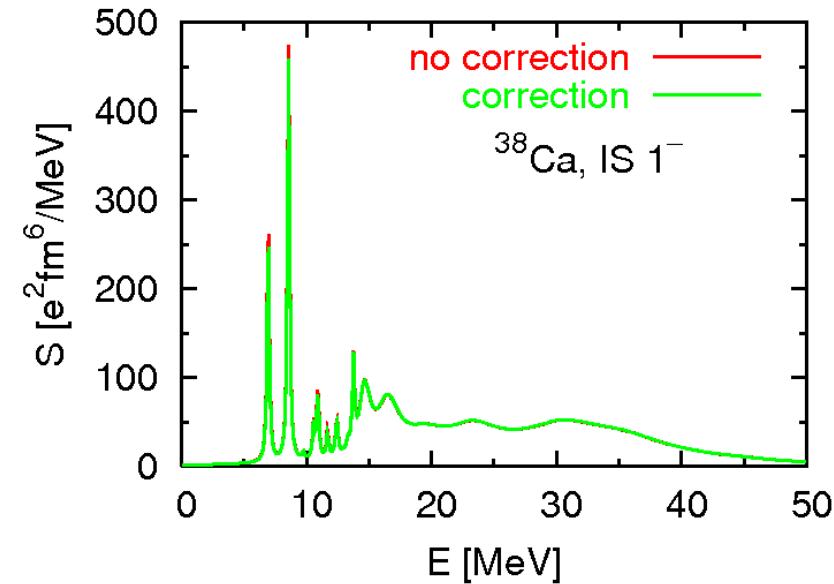
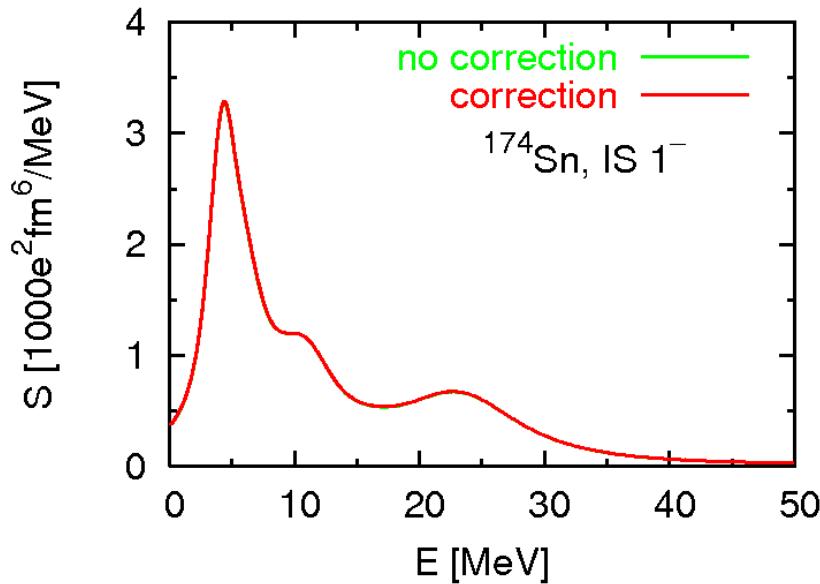
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## Separation of the spurious component; $J^\pi = 1^-$ mode

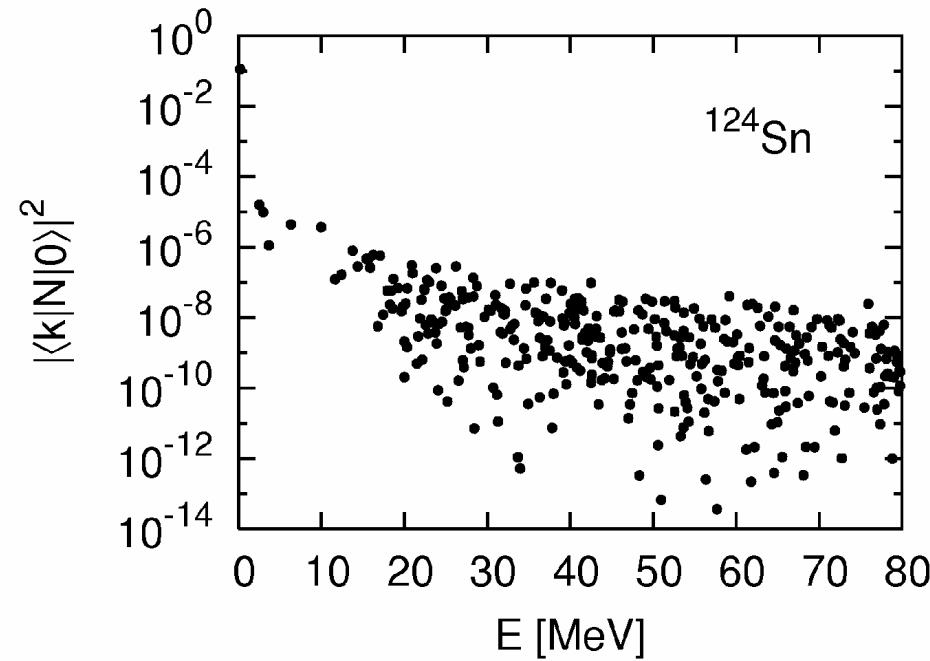


**no correction**     $F_{1M}(\underline{r}) = \underline{r}^3 Y_{1M}(O)$

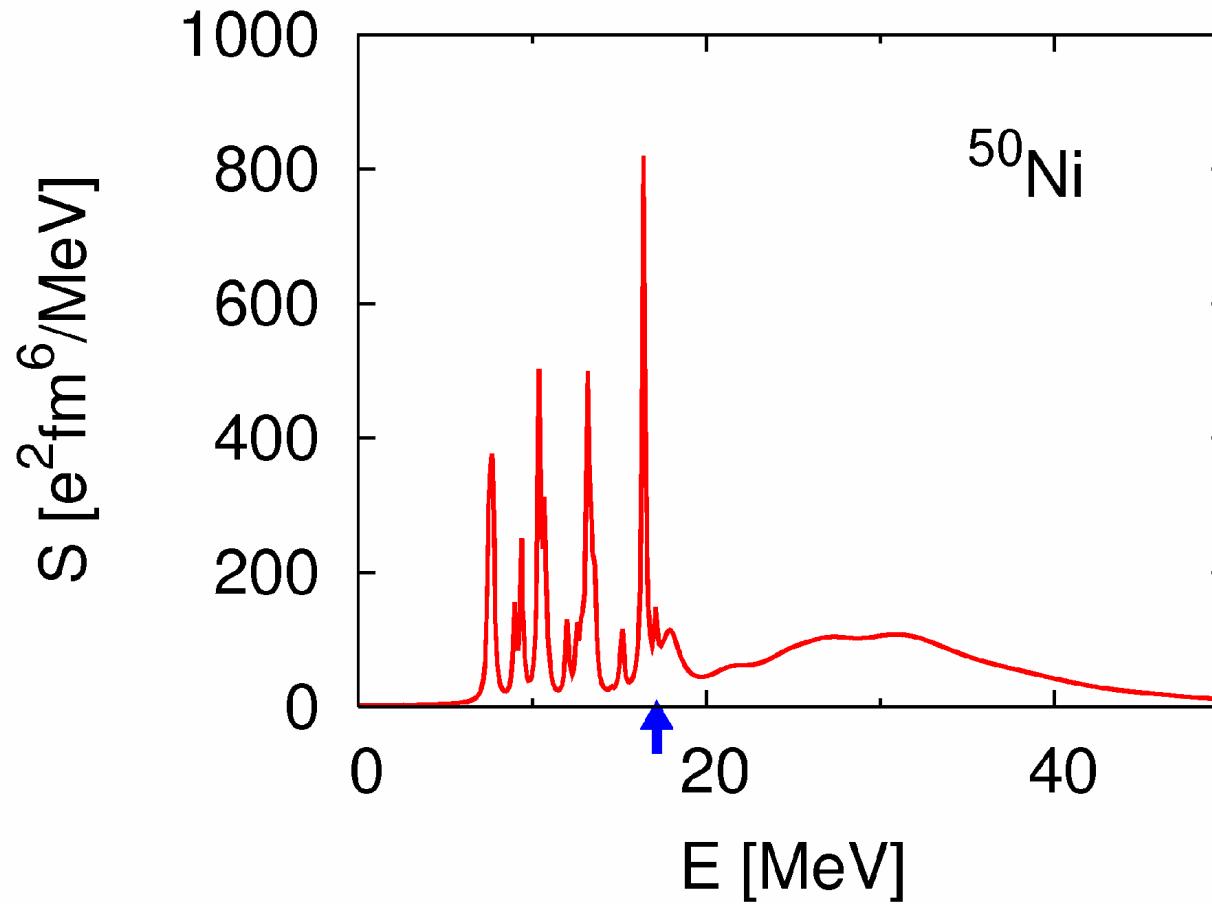
**correction**     $F_{1M}(\underline{r}) = \frac{\alpha}{\epsilon} \underline{r}^3 - \frac{5}{3} \langle \underline{r}^2 \rangle \underline{r} \overset{\circ}{\div} Y_{1M}(O)$

## Separation of the spurious component; $J^\pi = 0^+$ mode

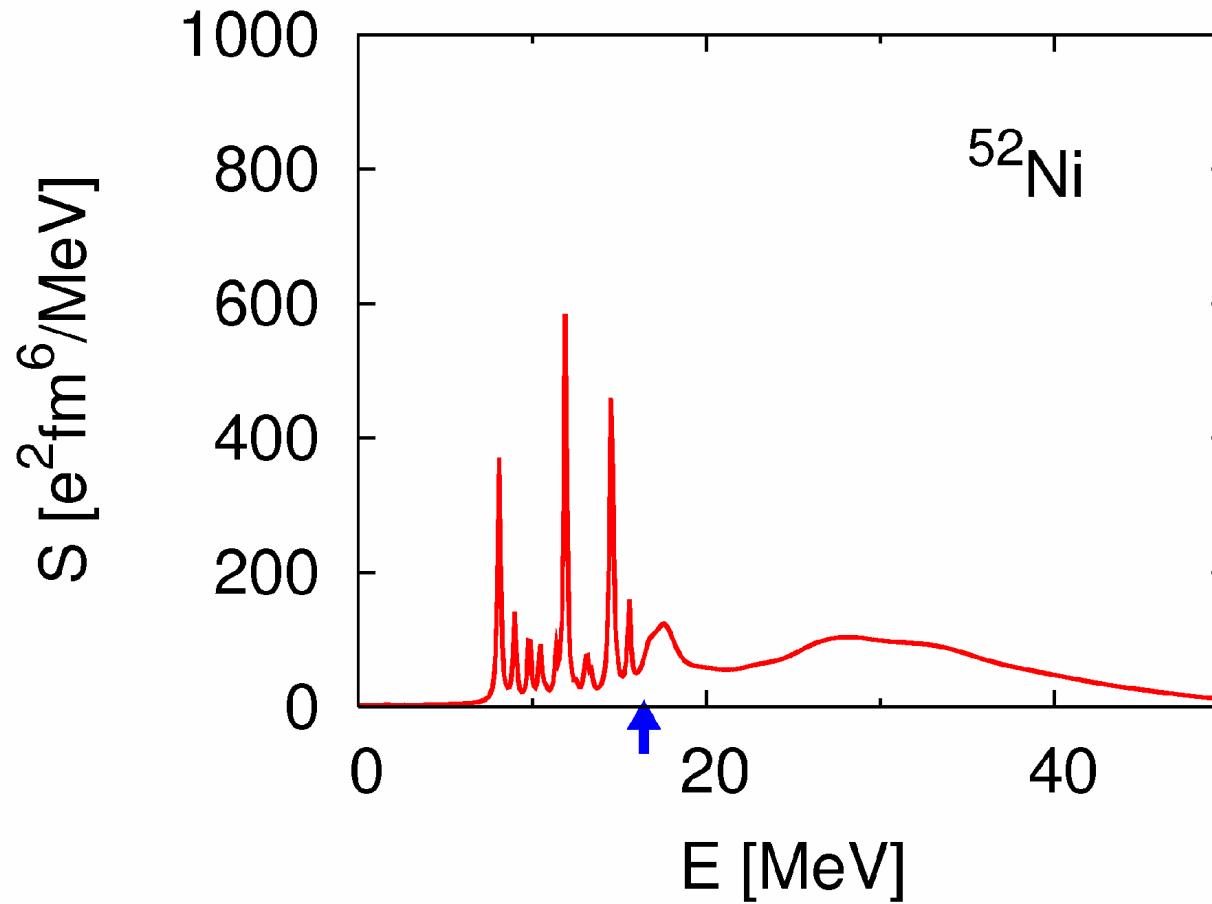
Strength function for the neutron number operator



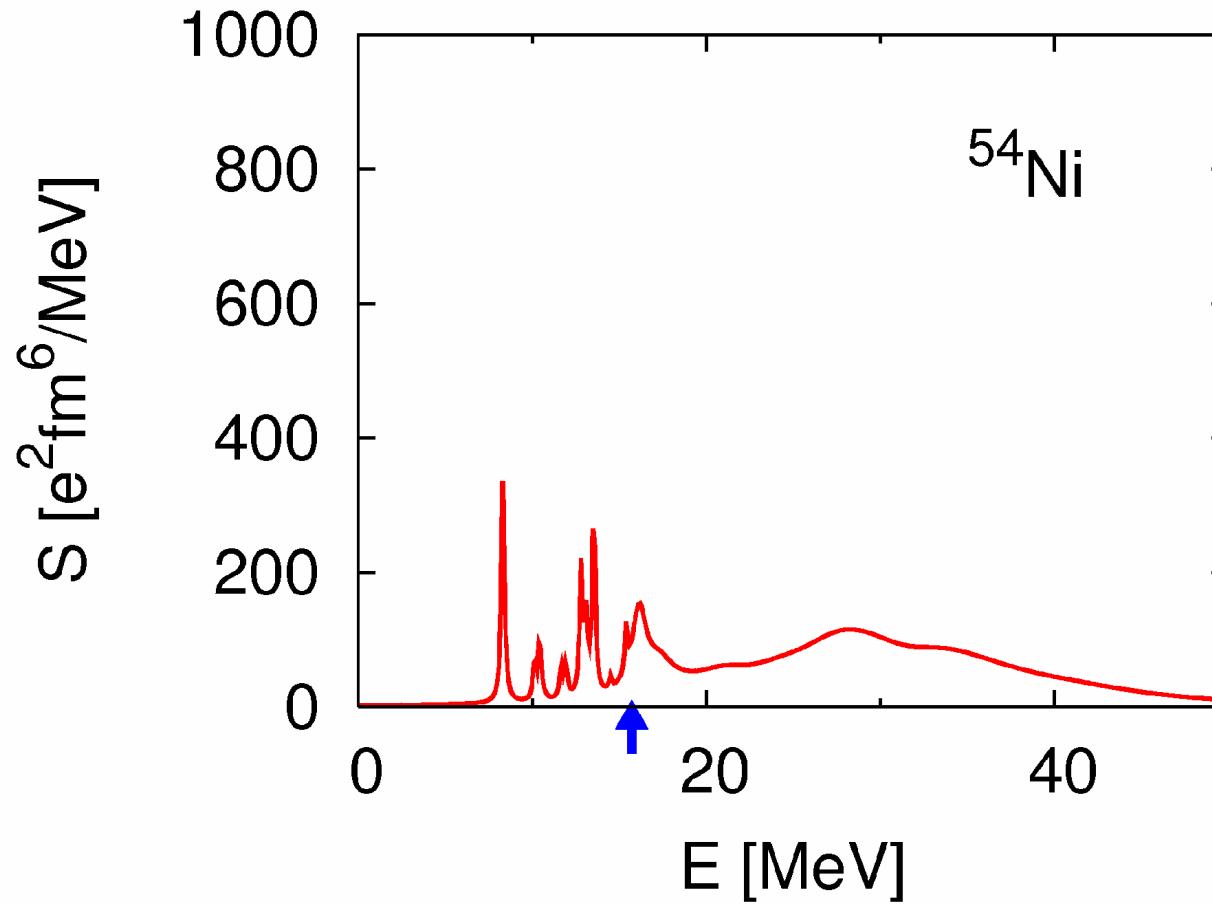
## Isoscalar $1^-$ strength functions



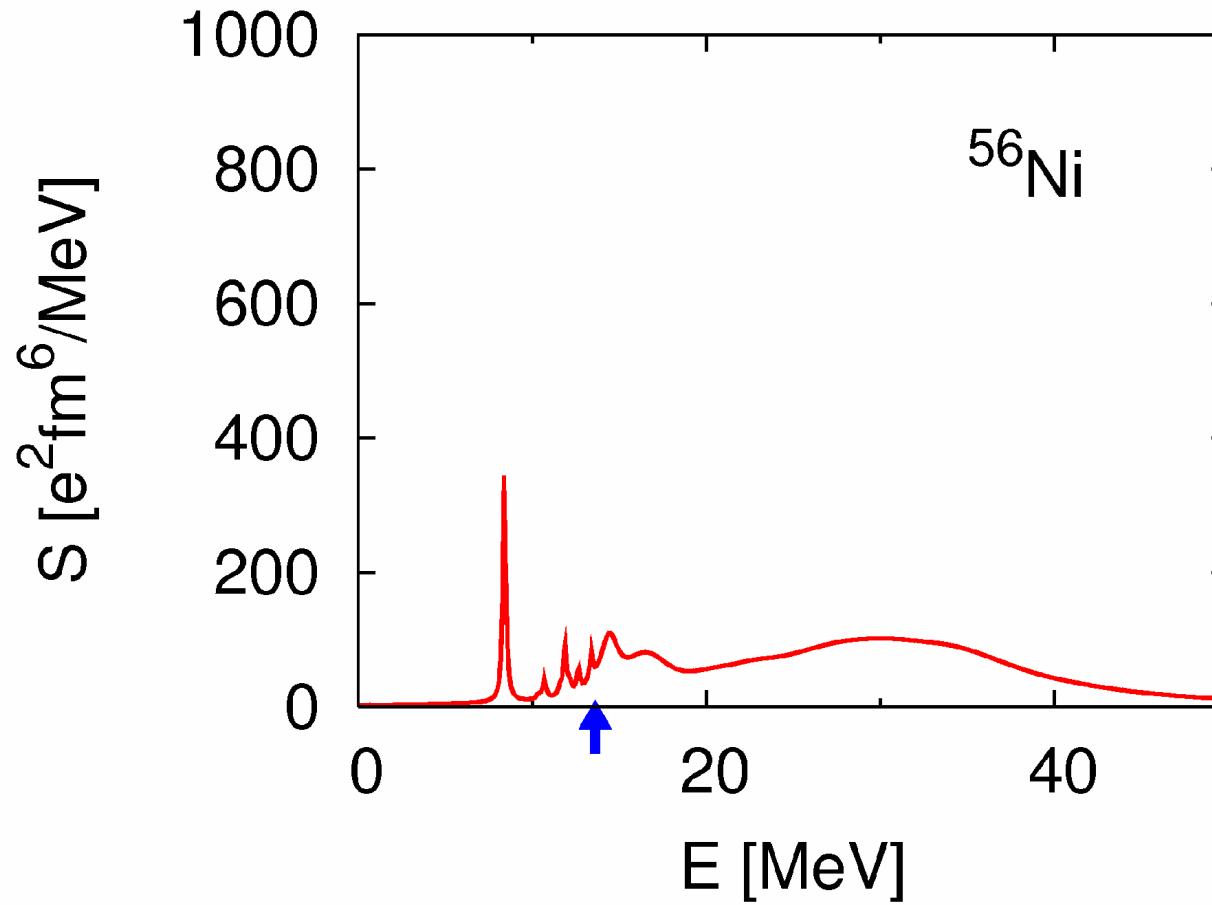
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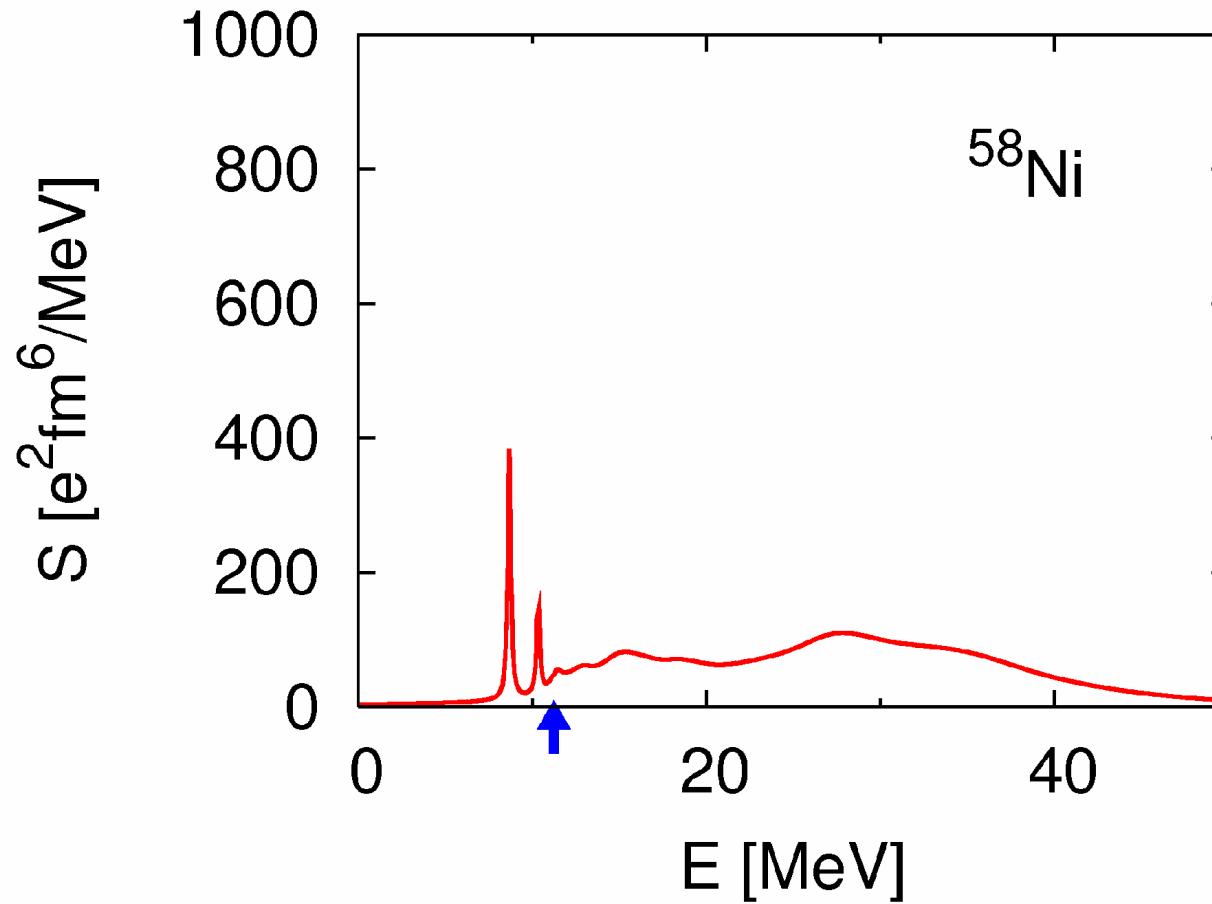
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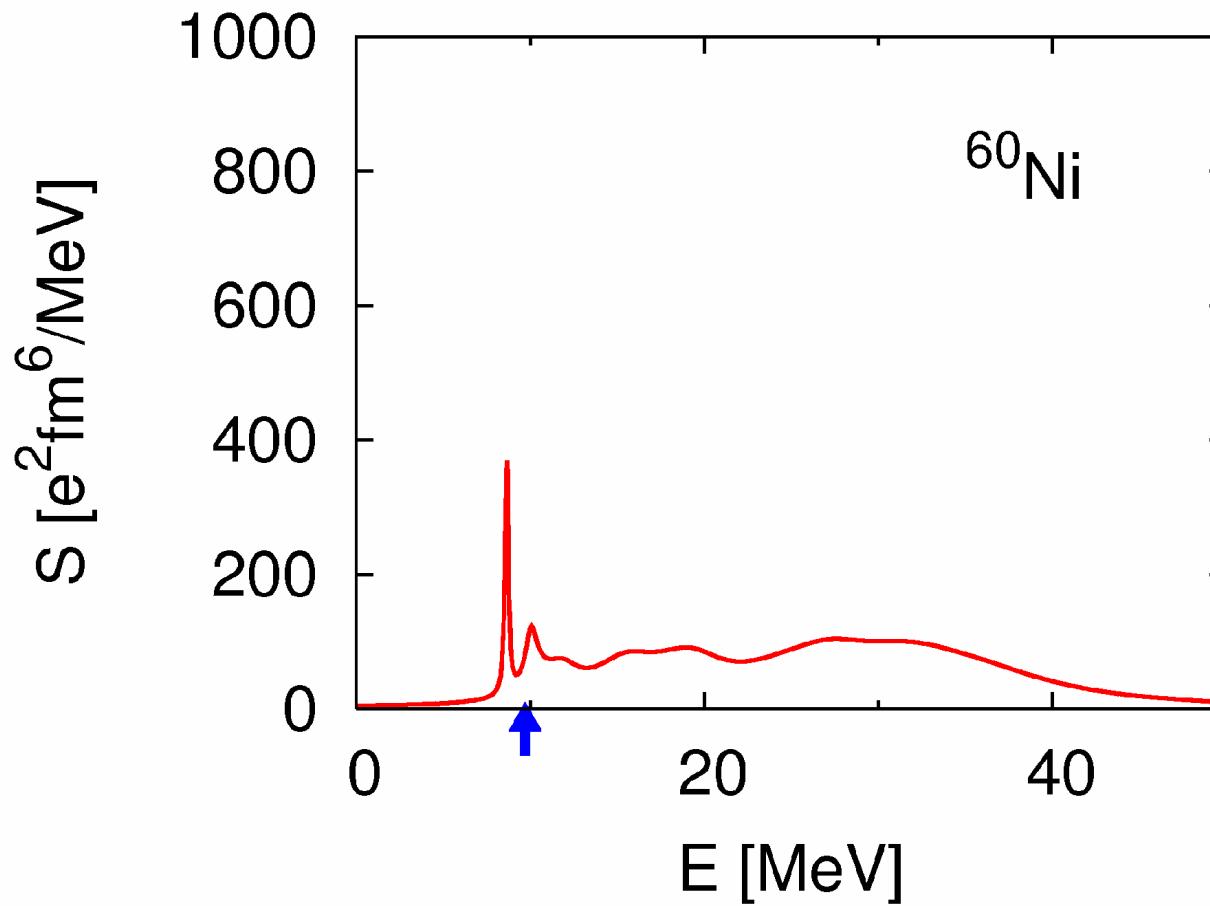
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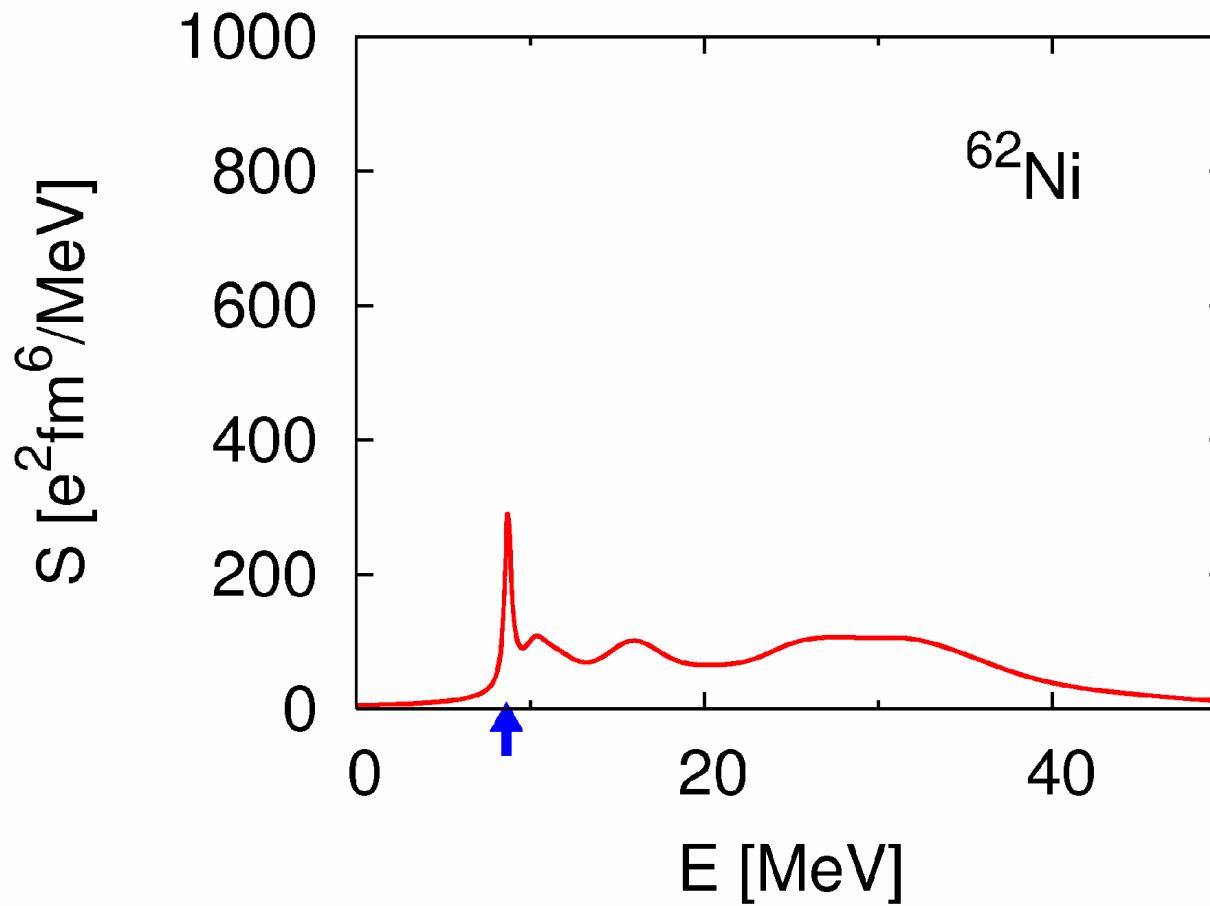
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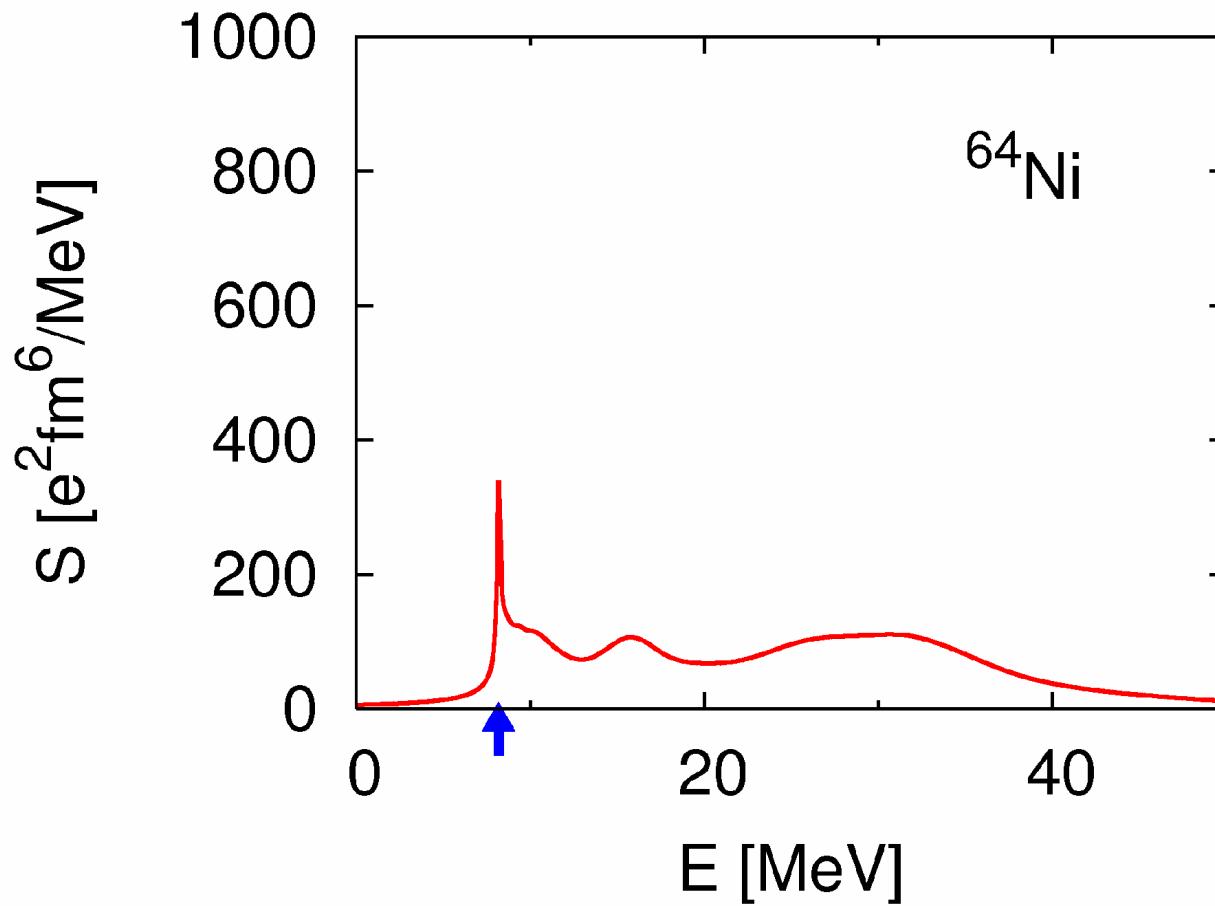
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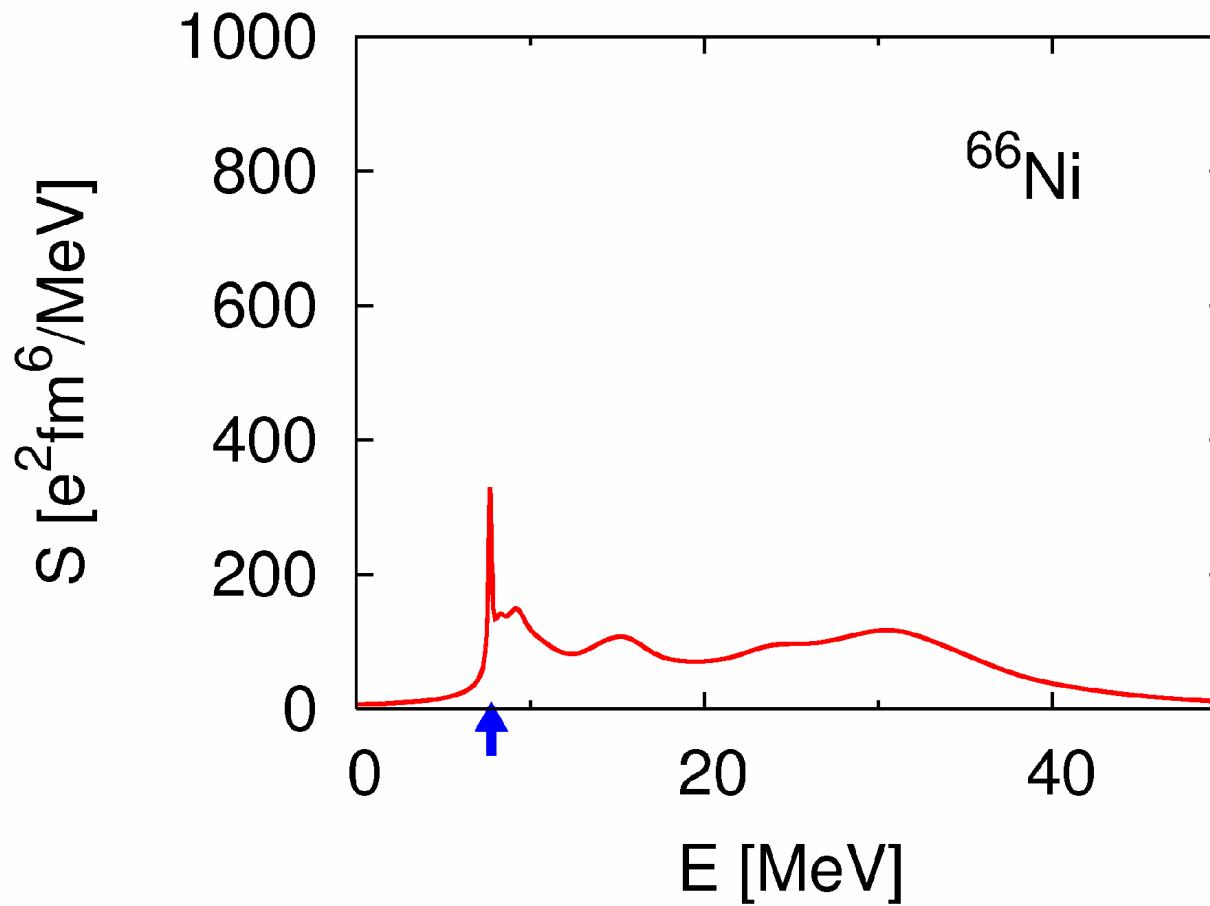
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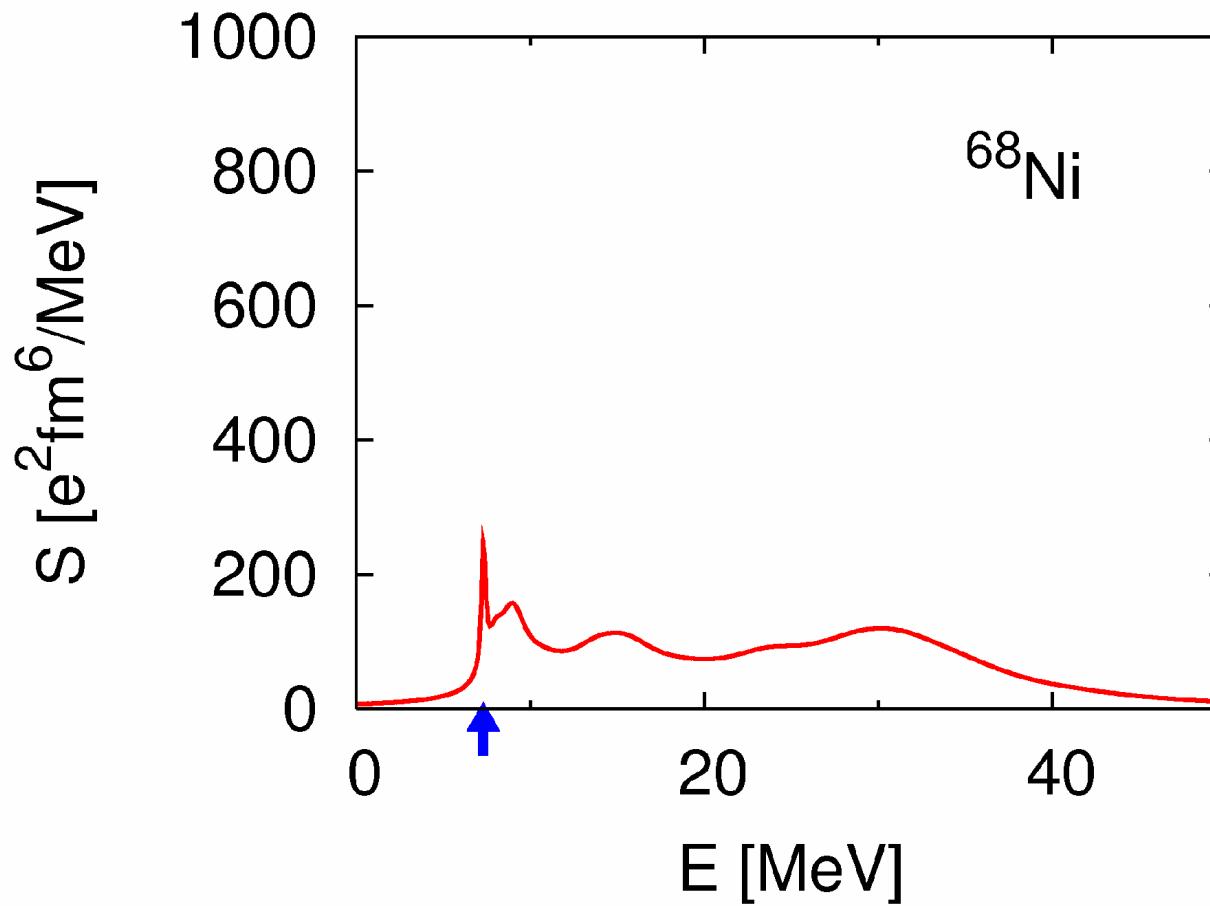
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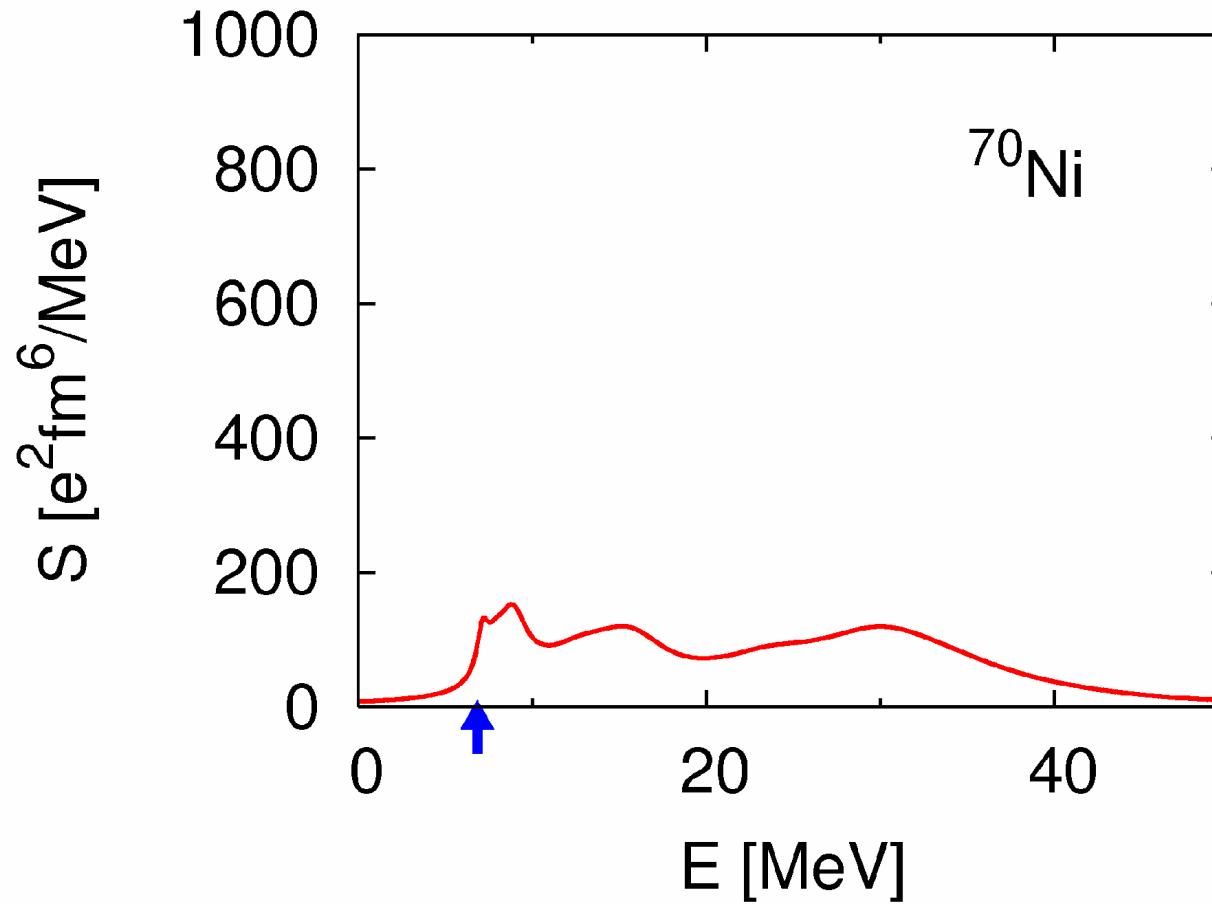
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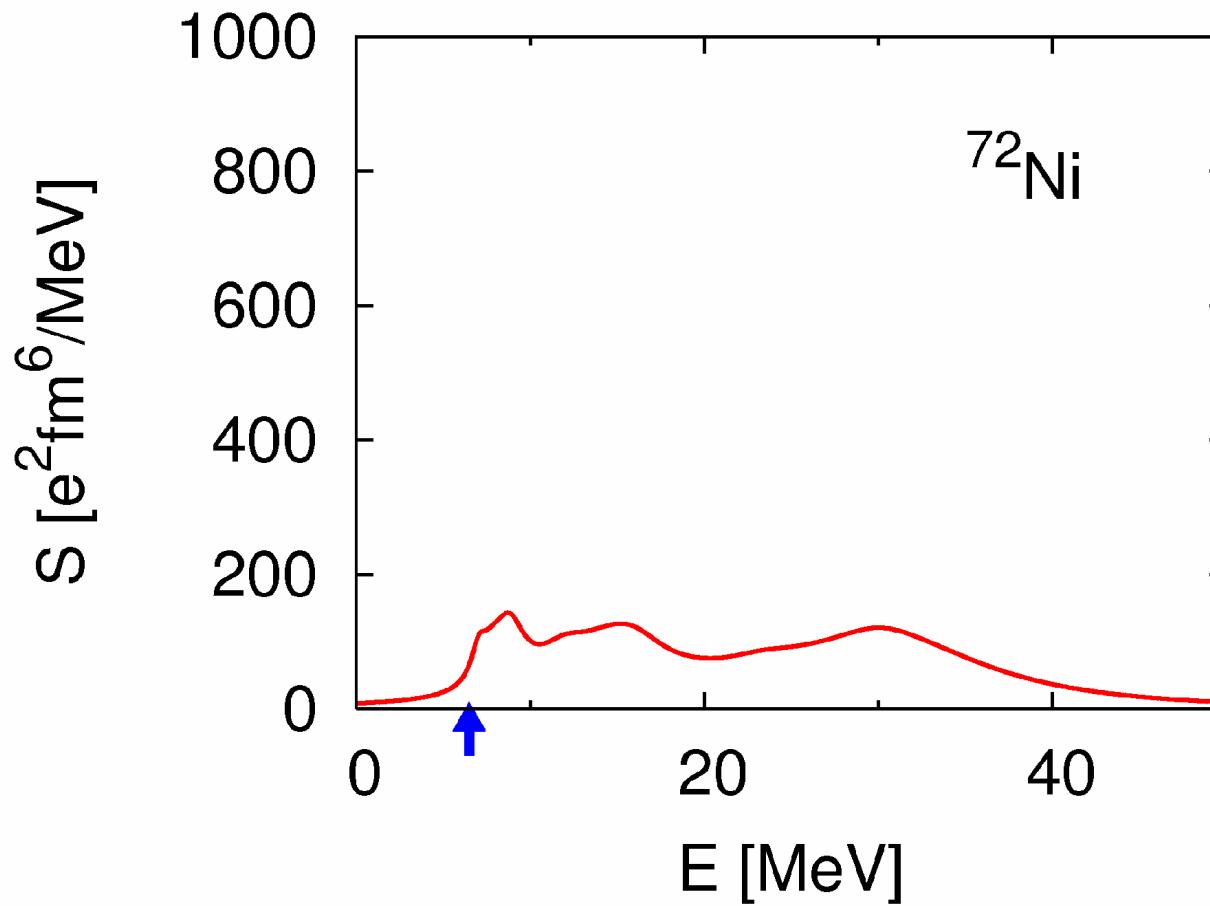
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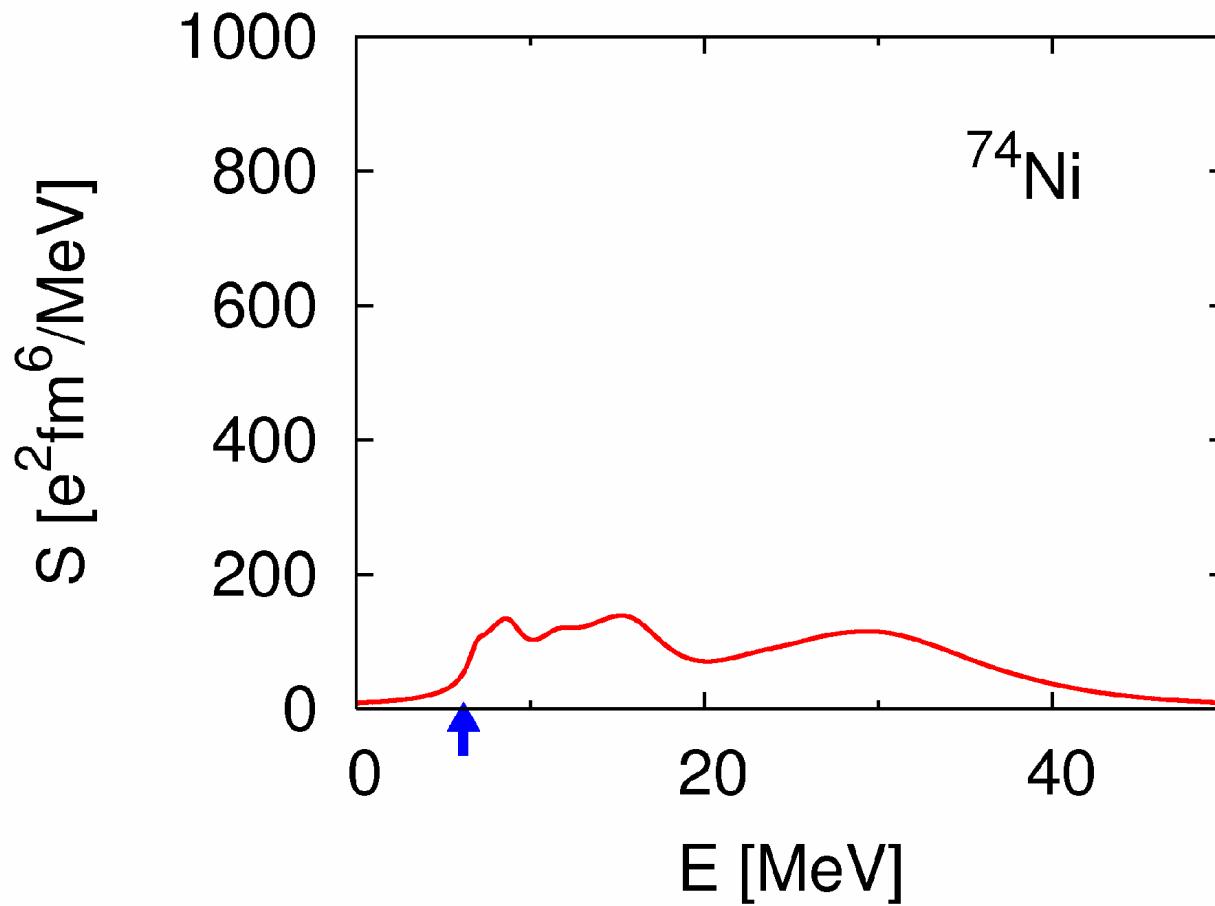
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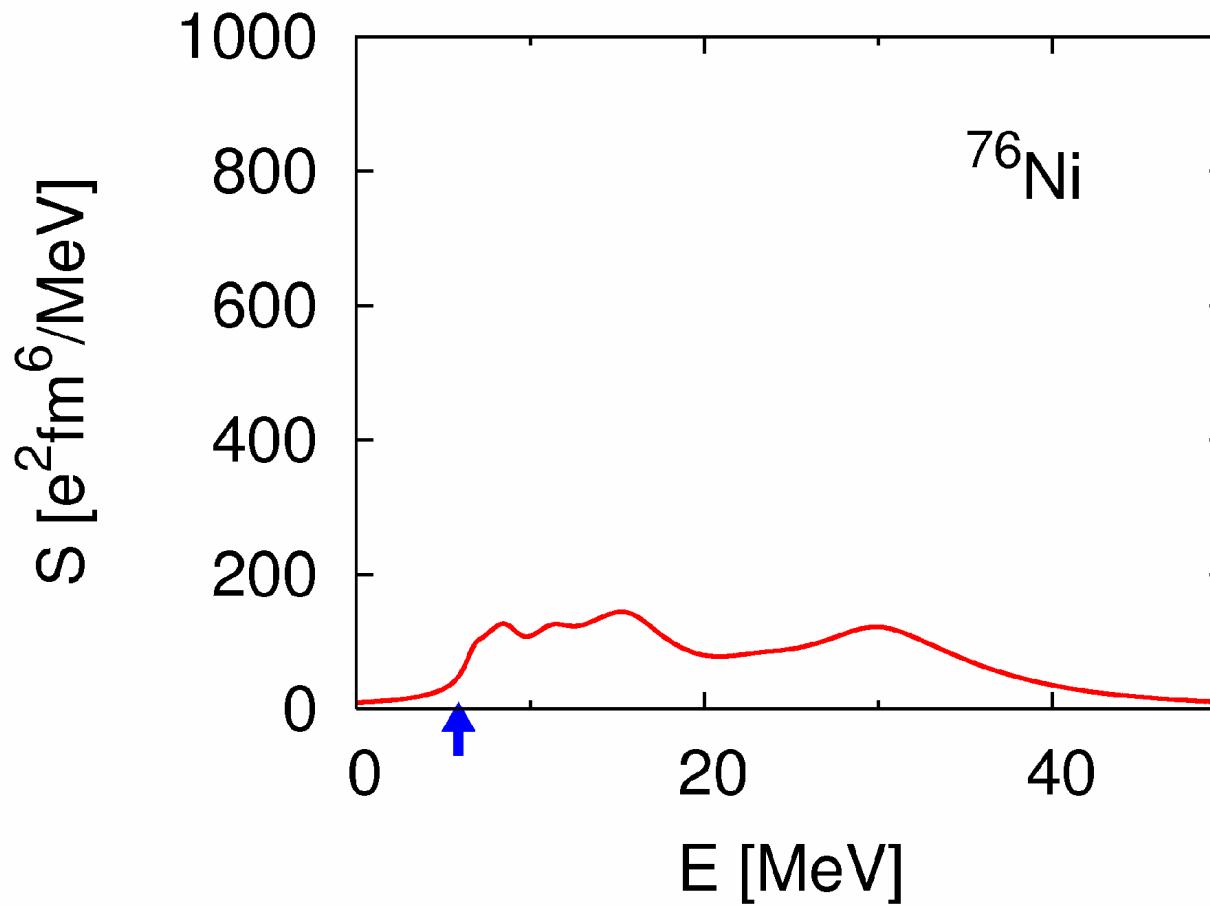
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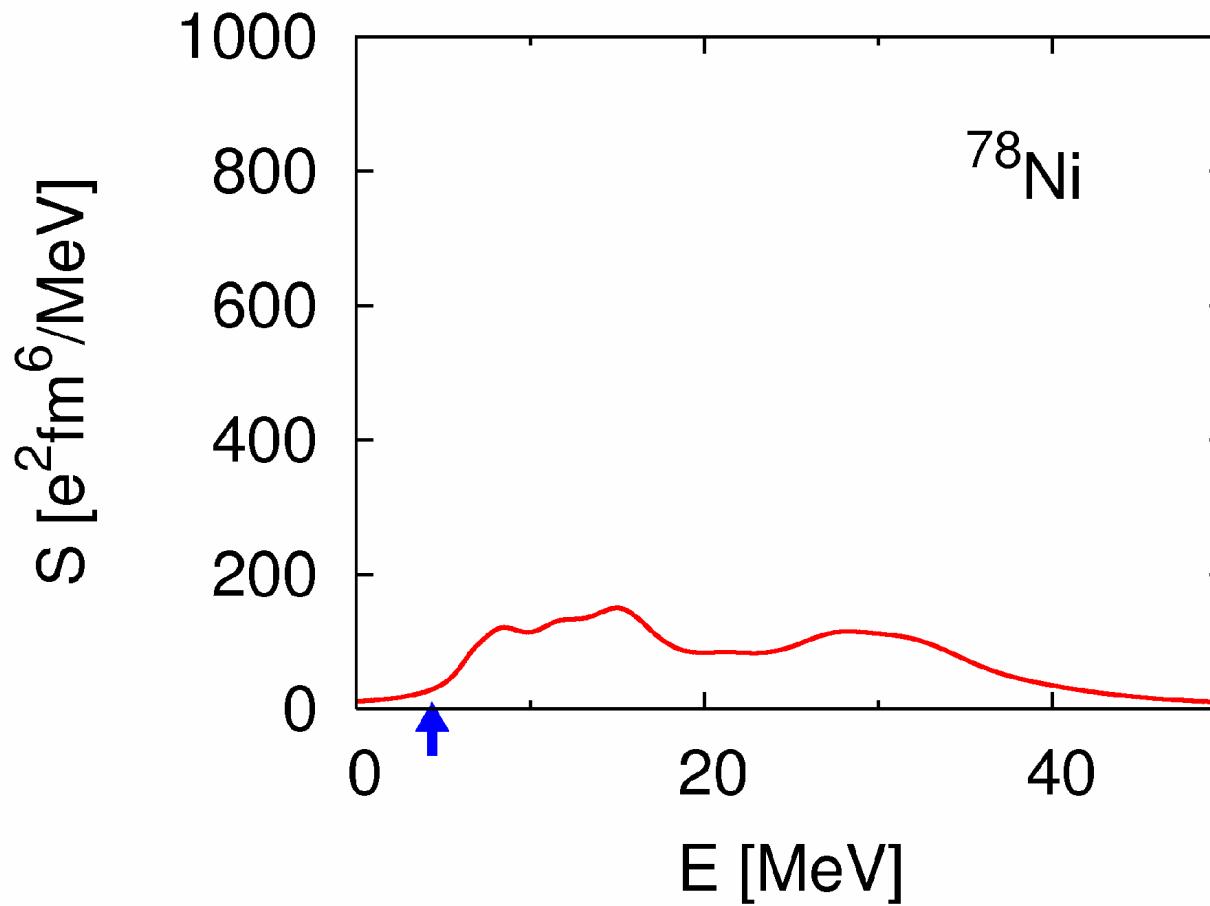
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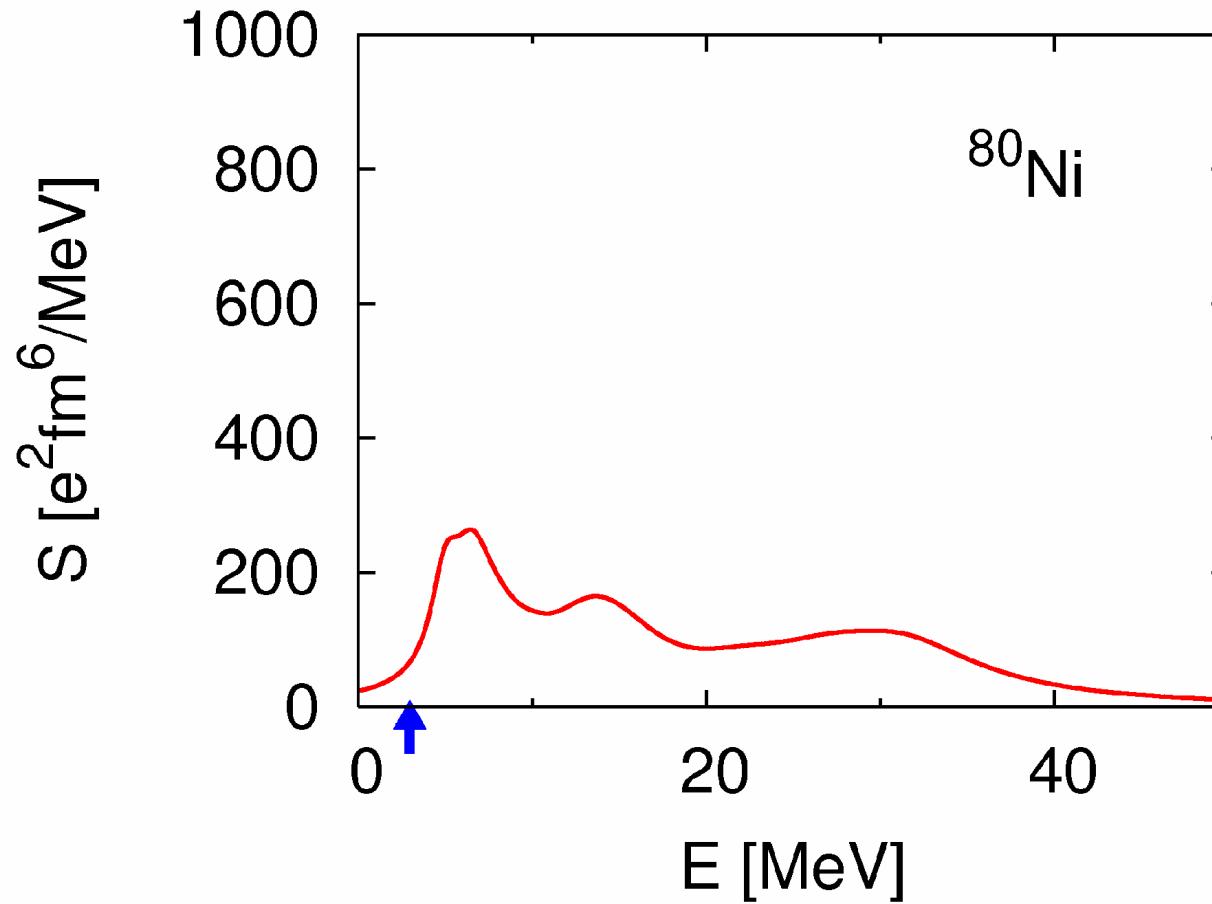
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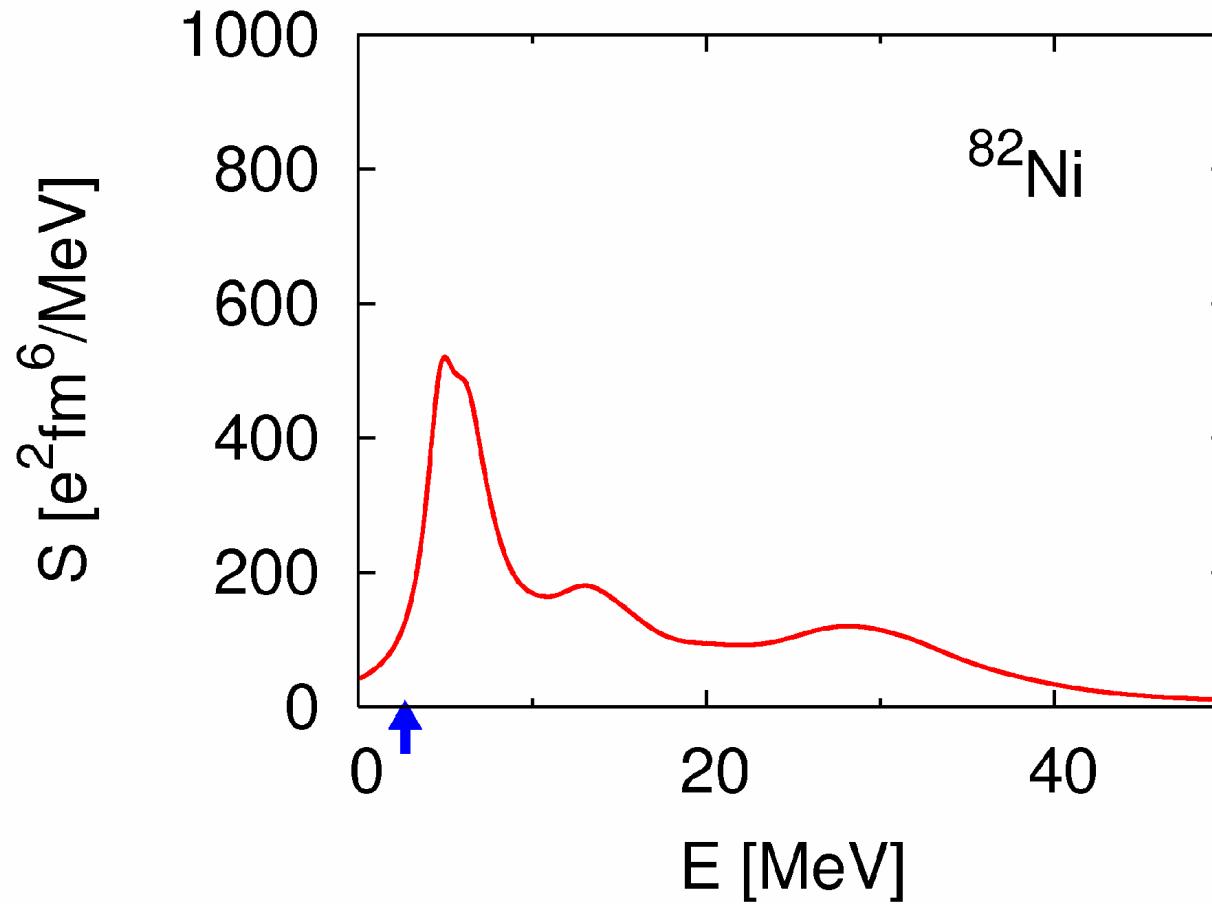
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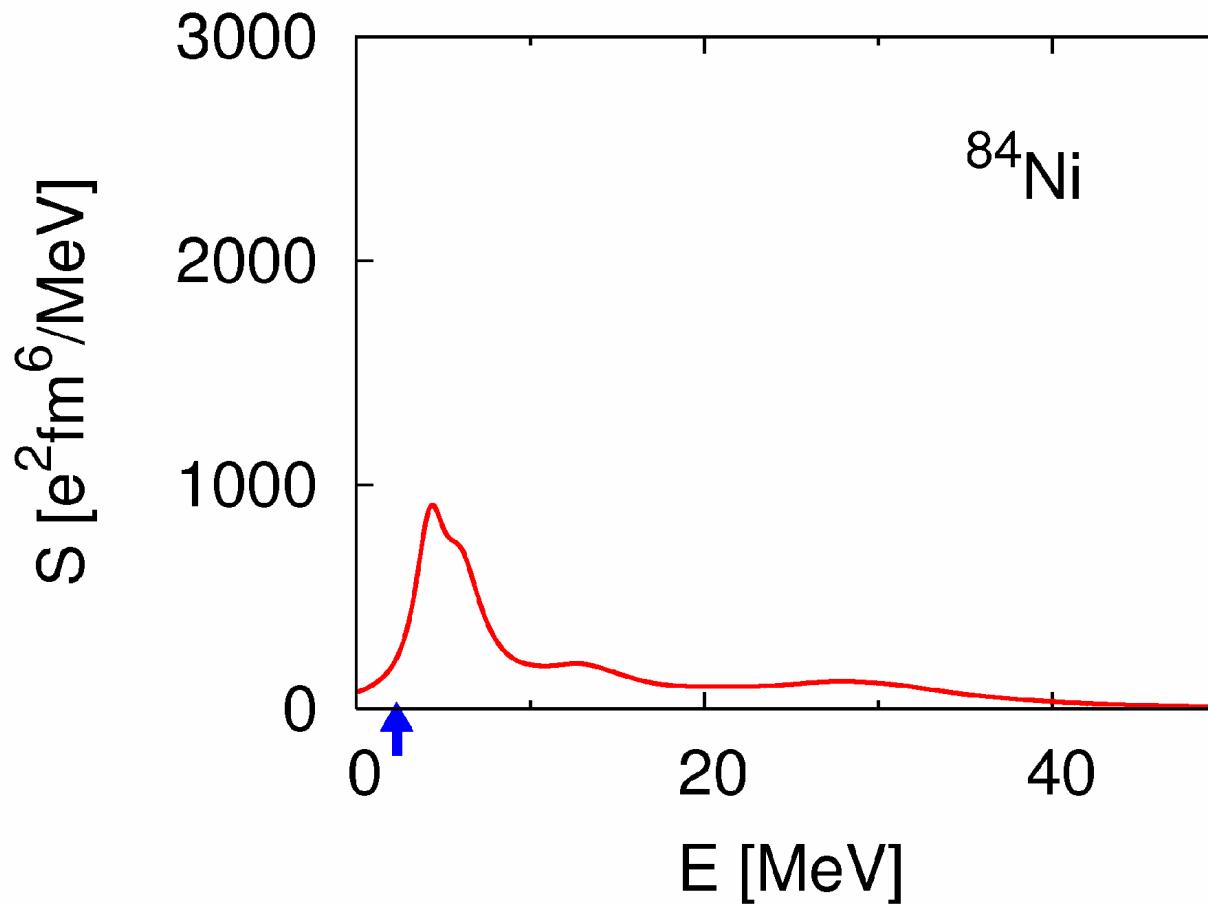
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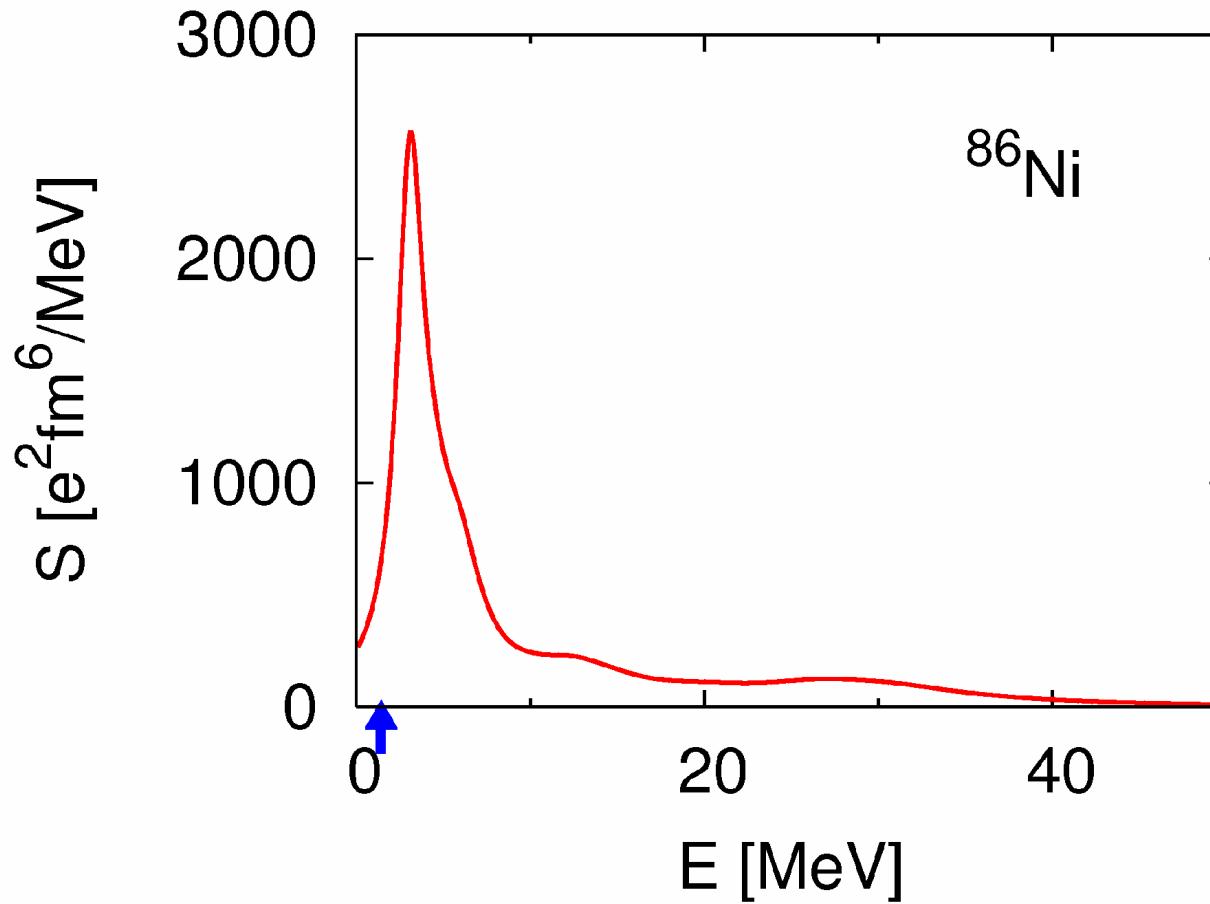
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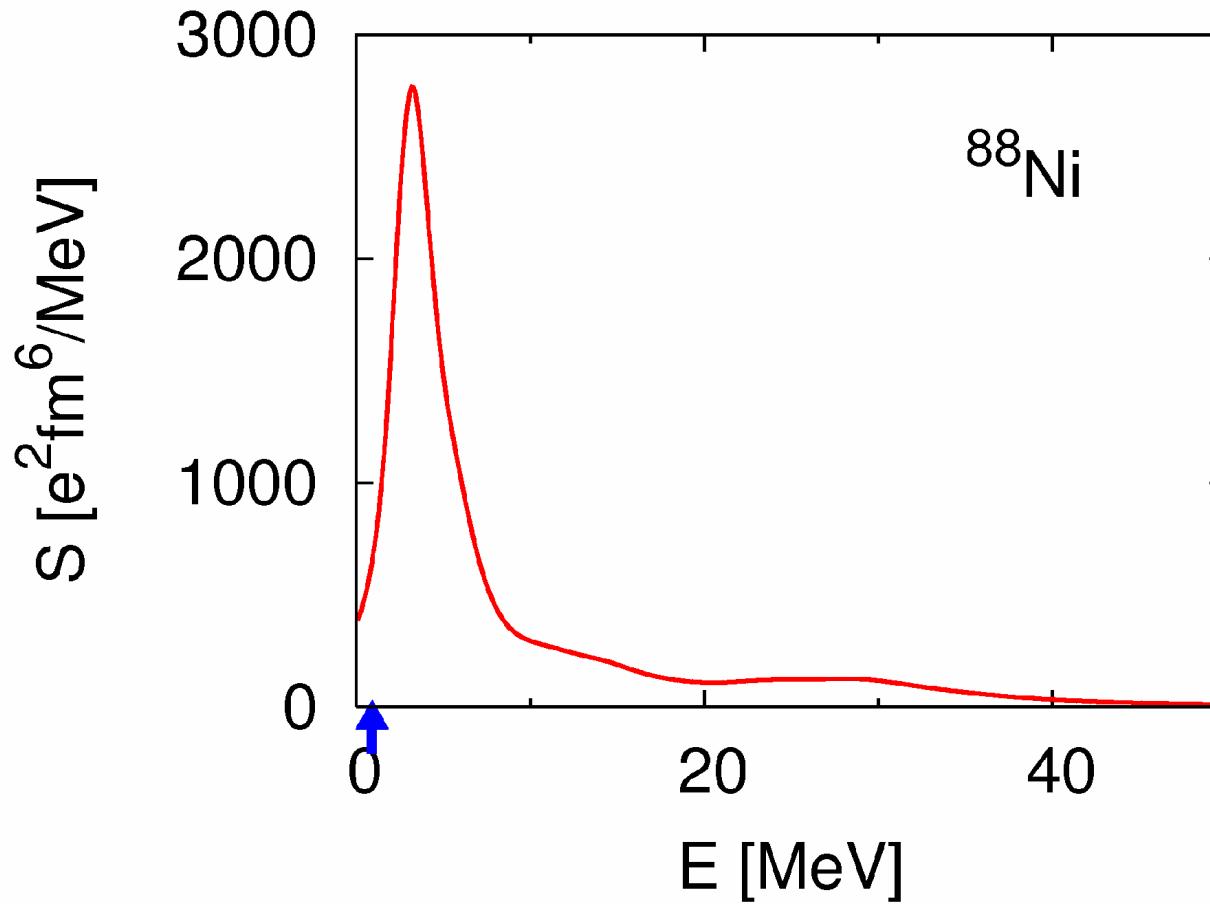
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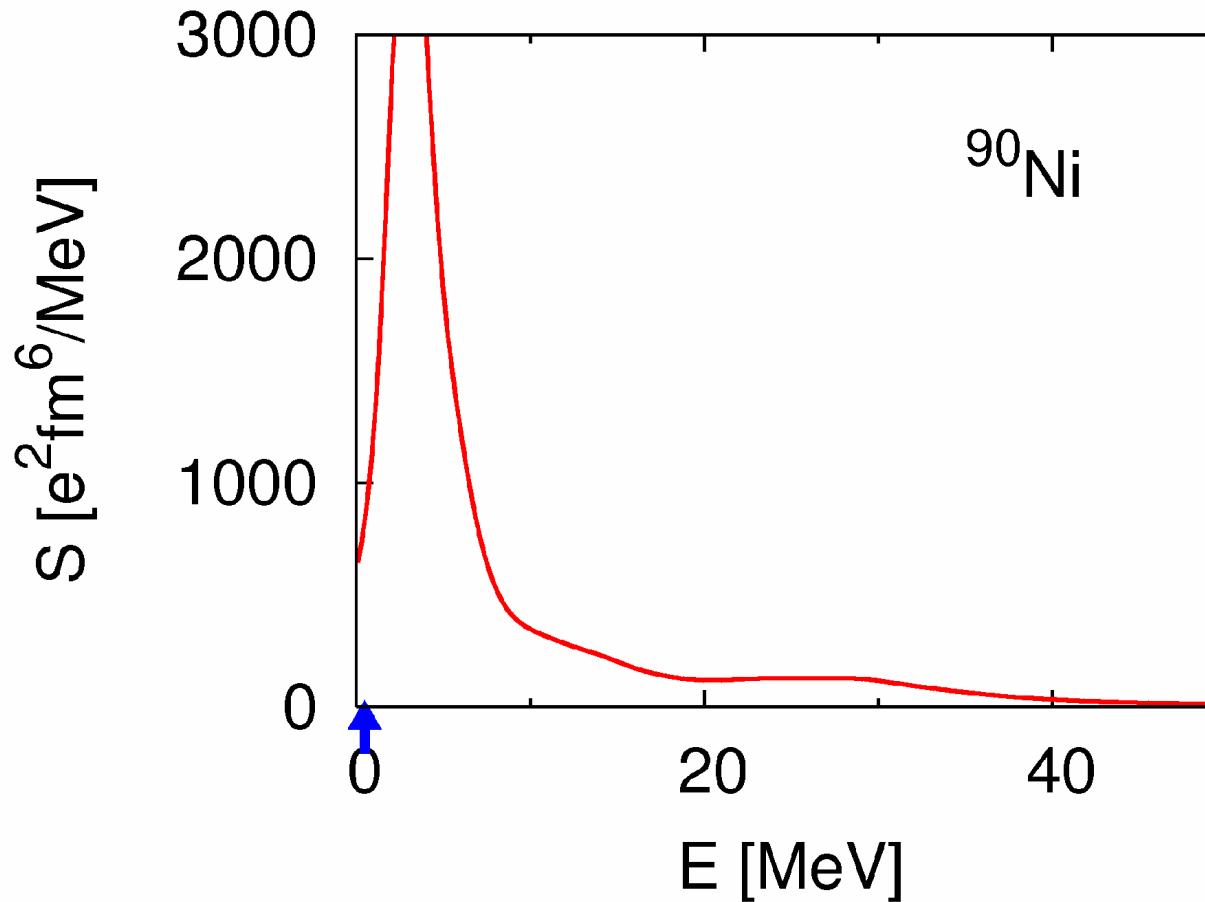
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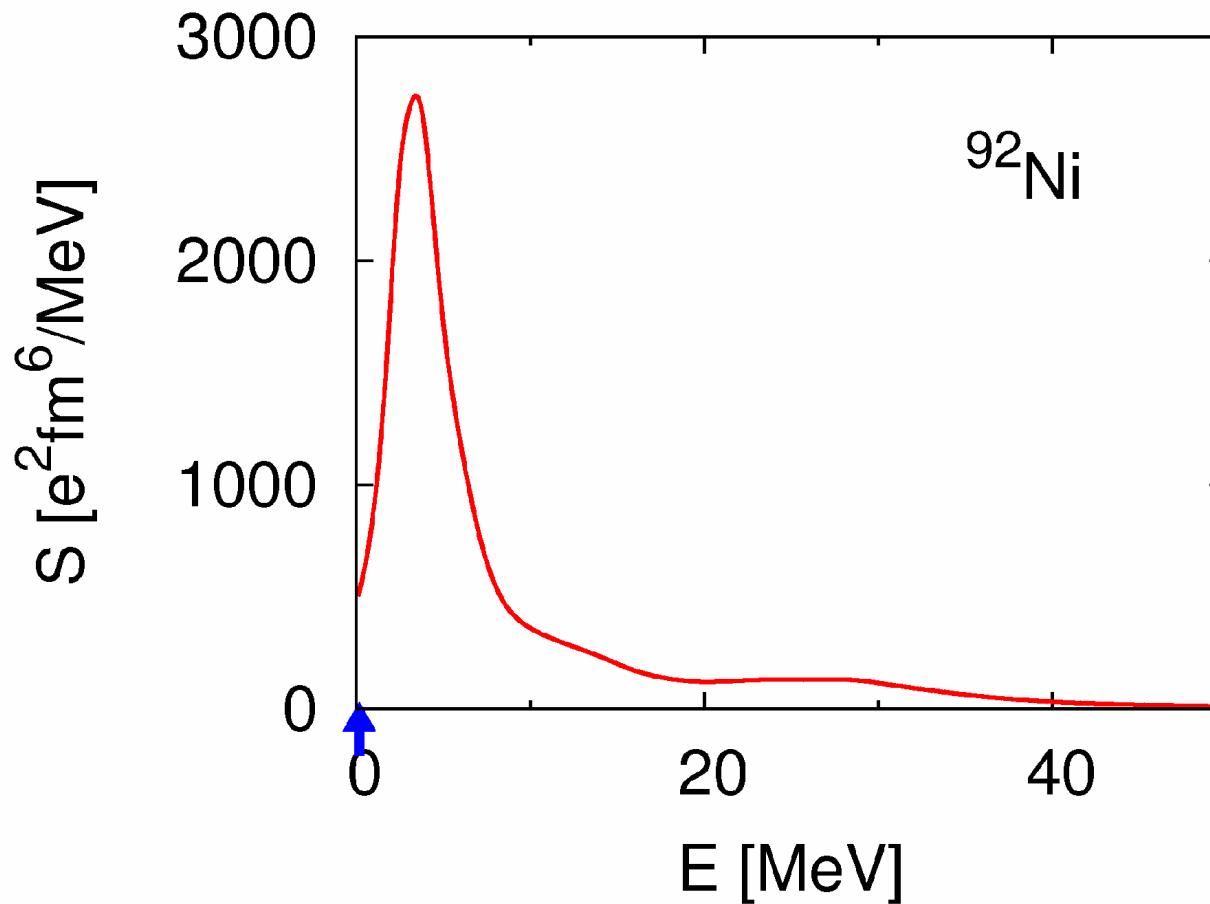
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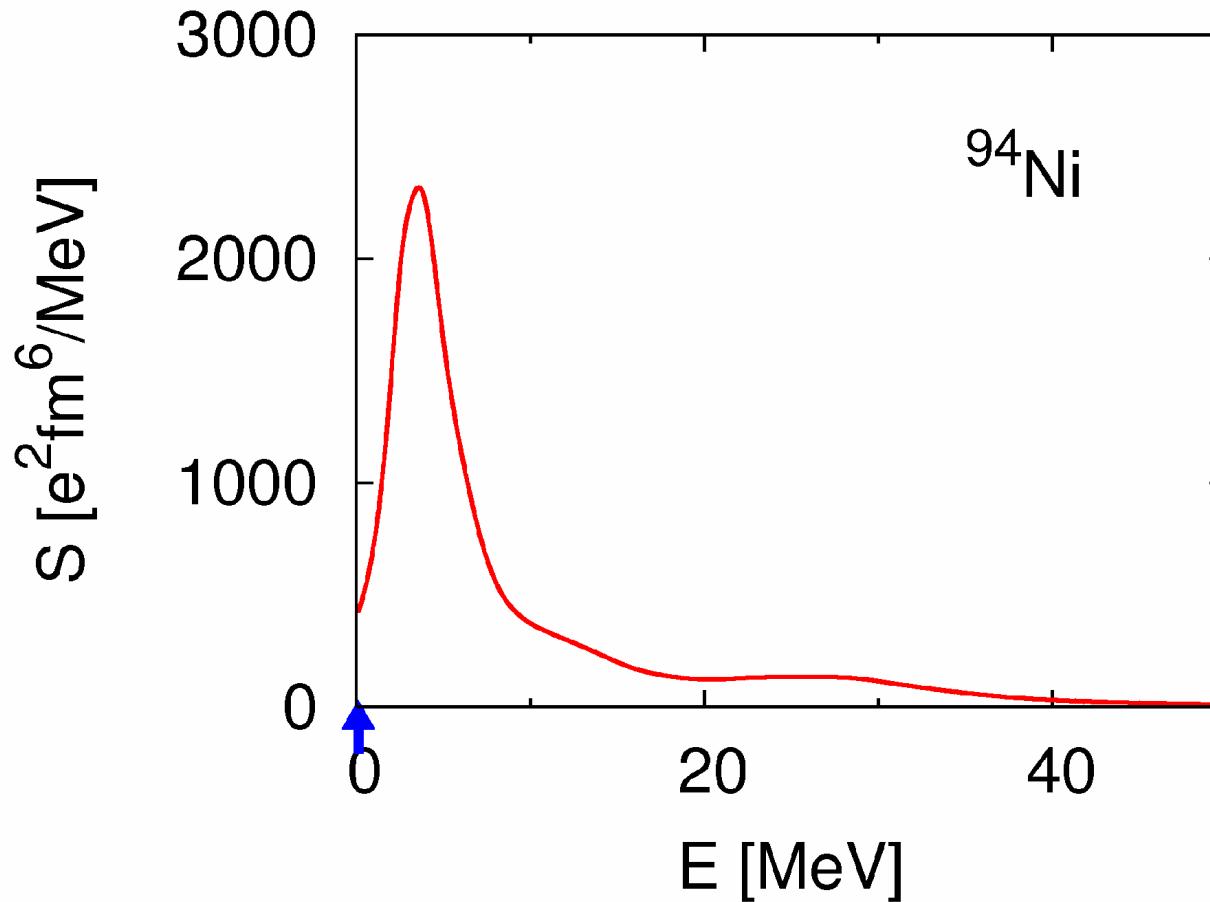
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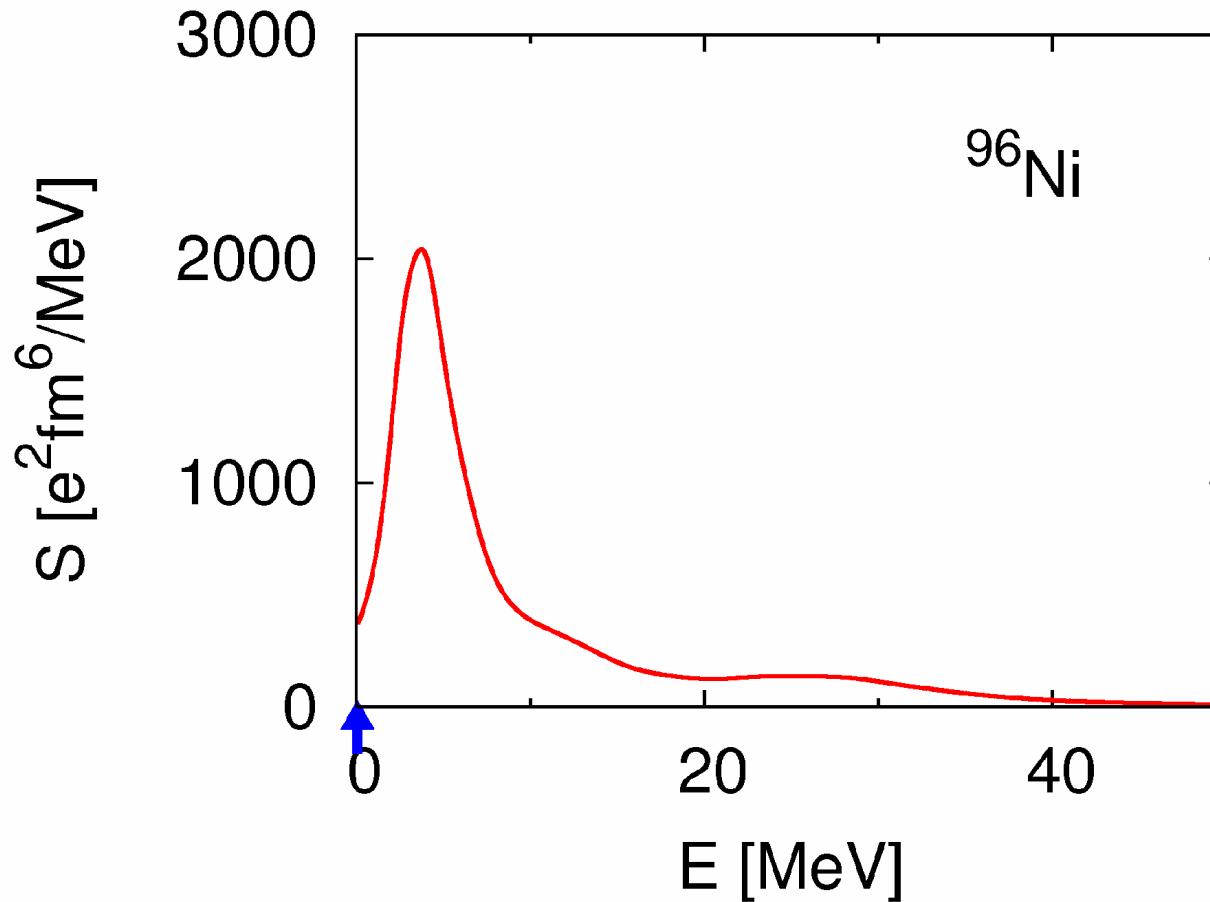
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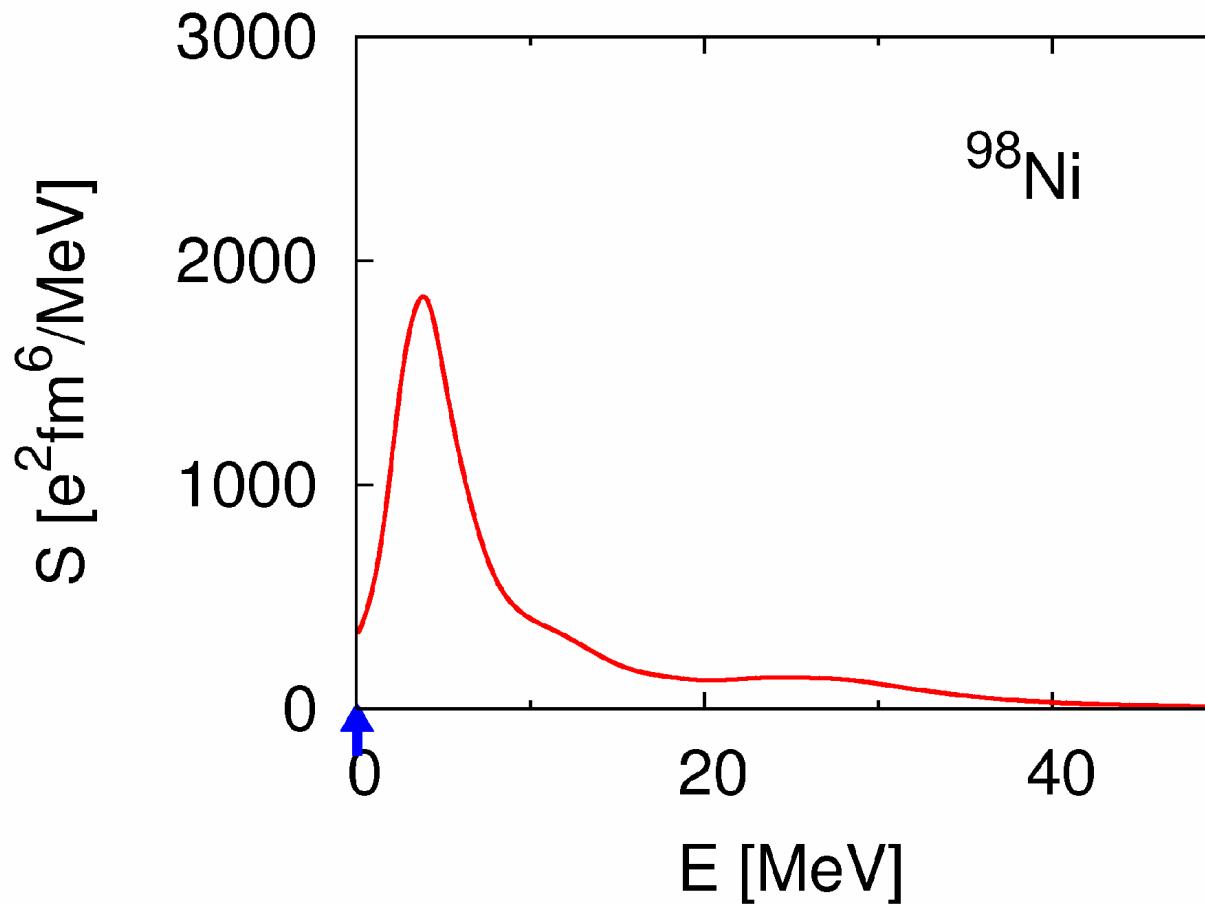
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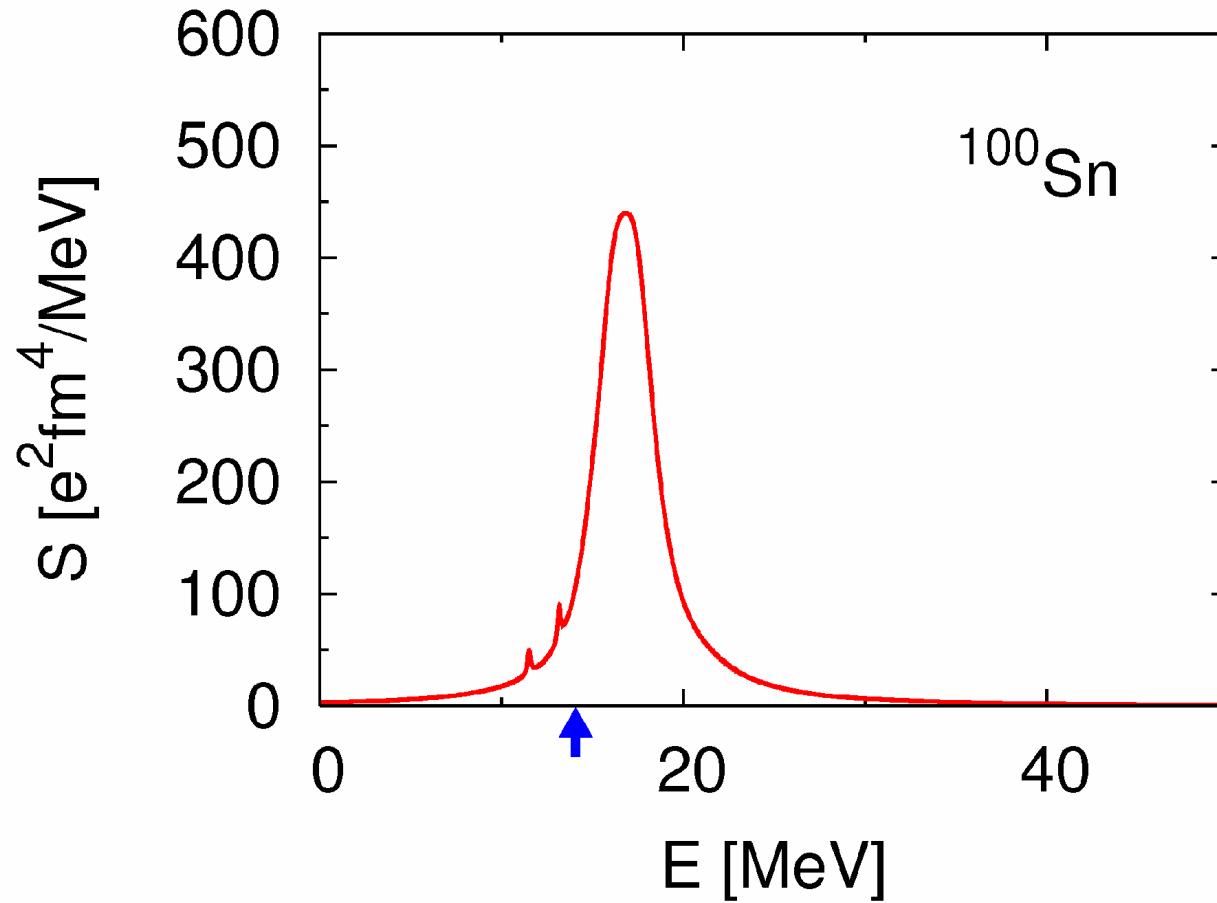
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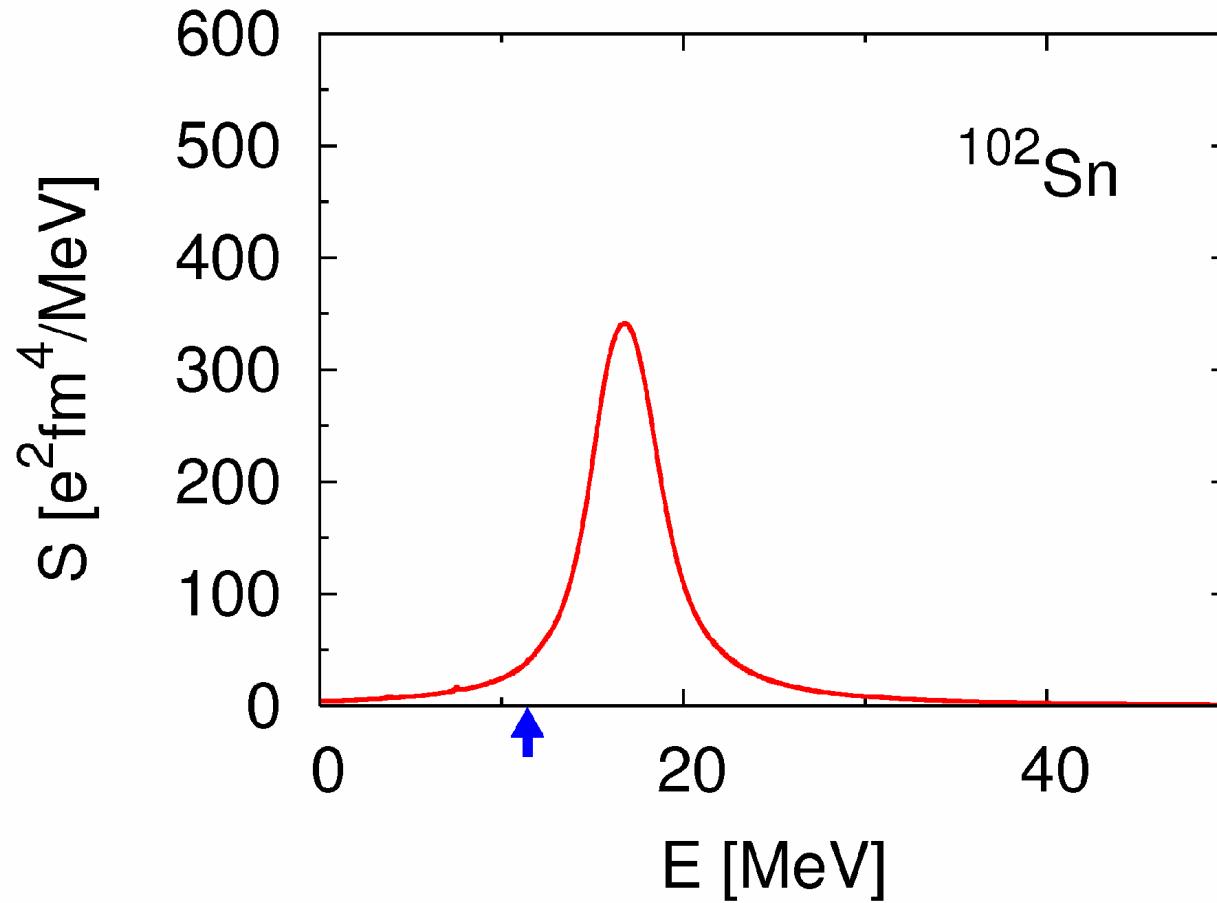
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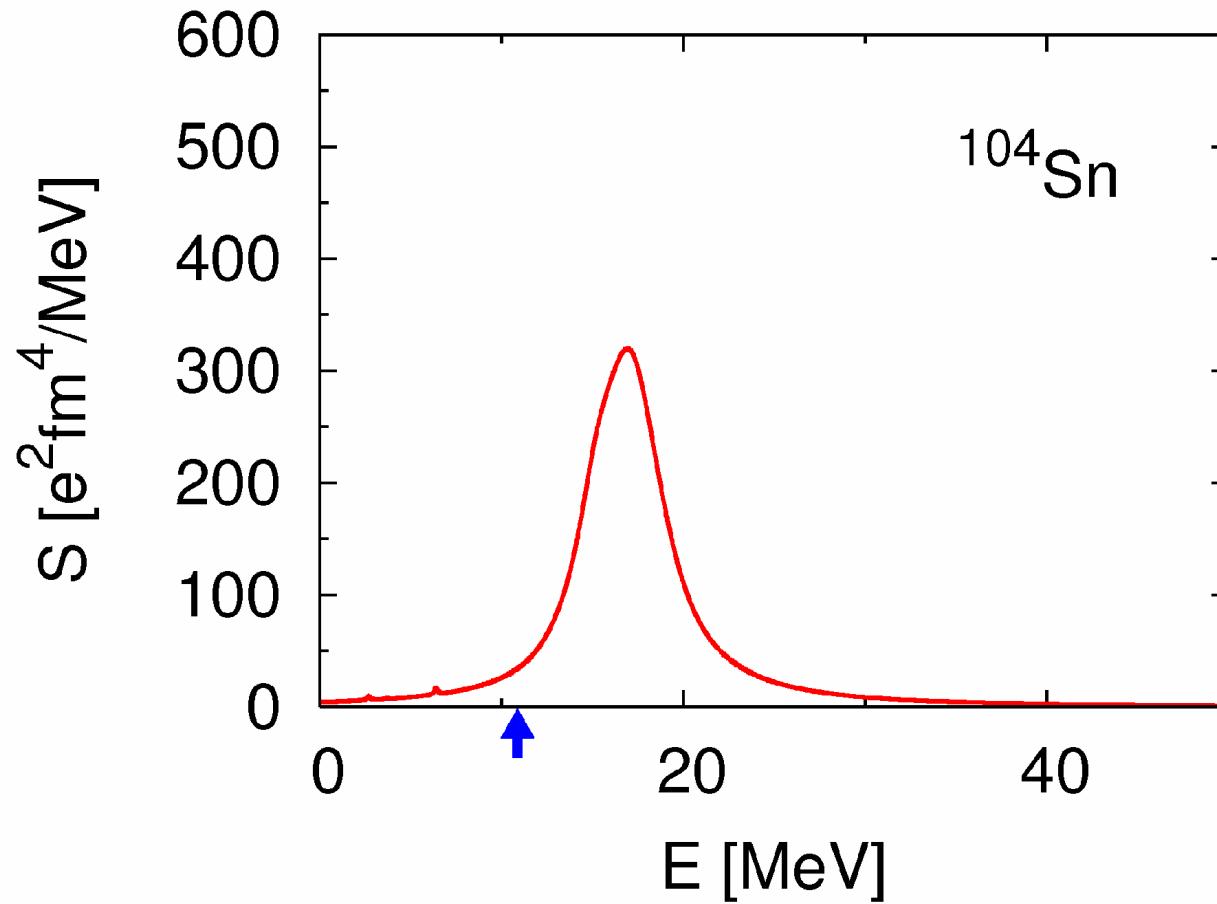
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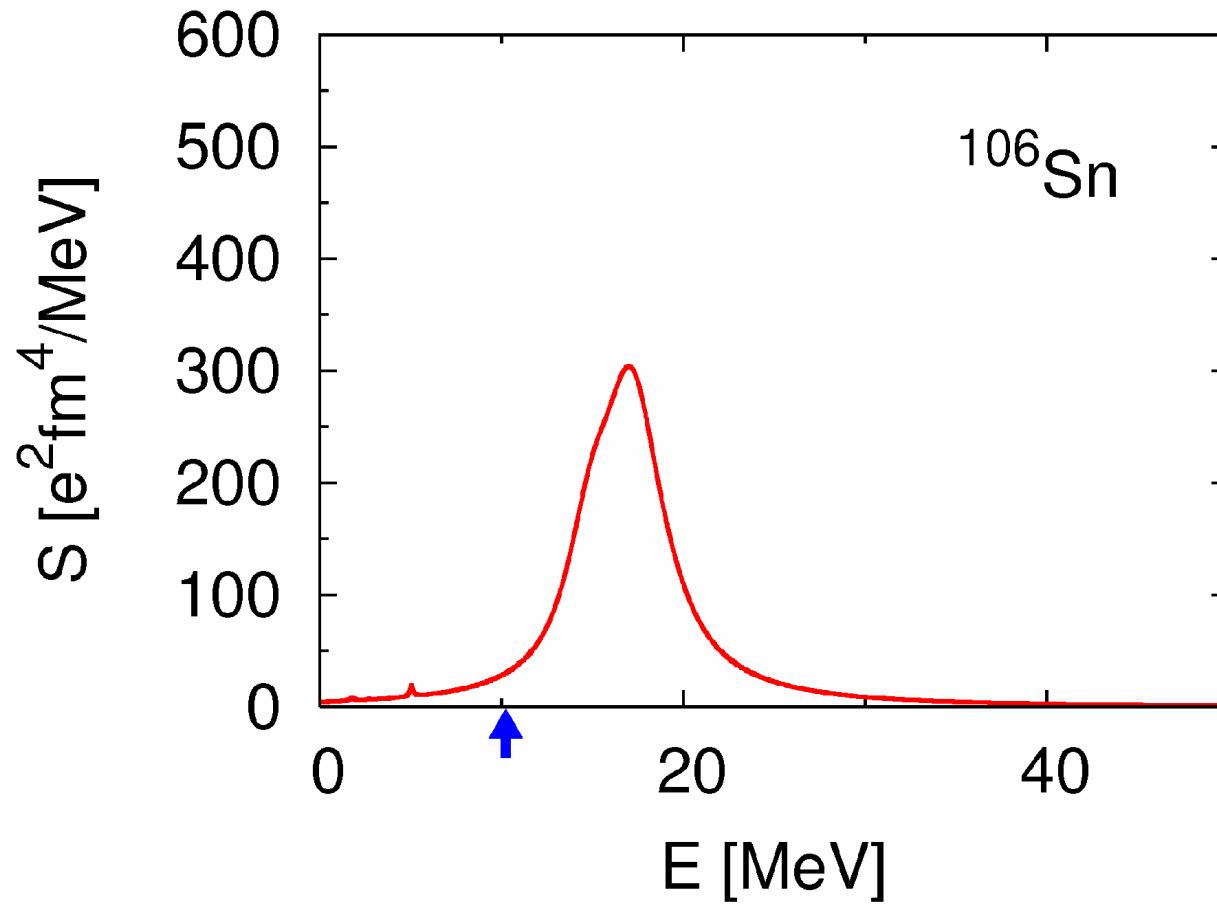
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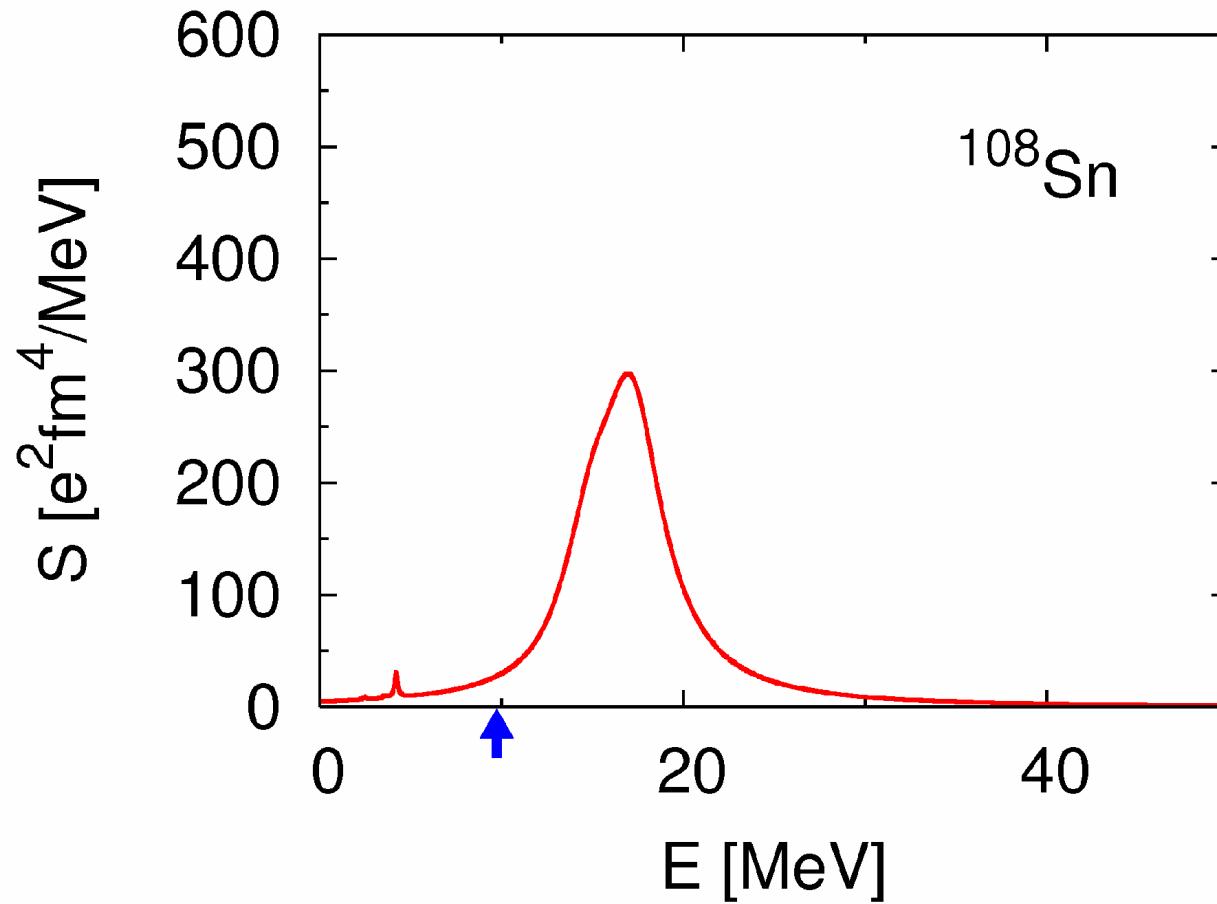
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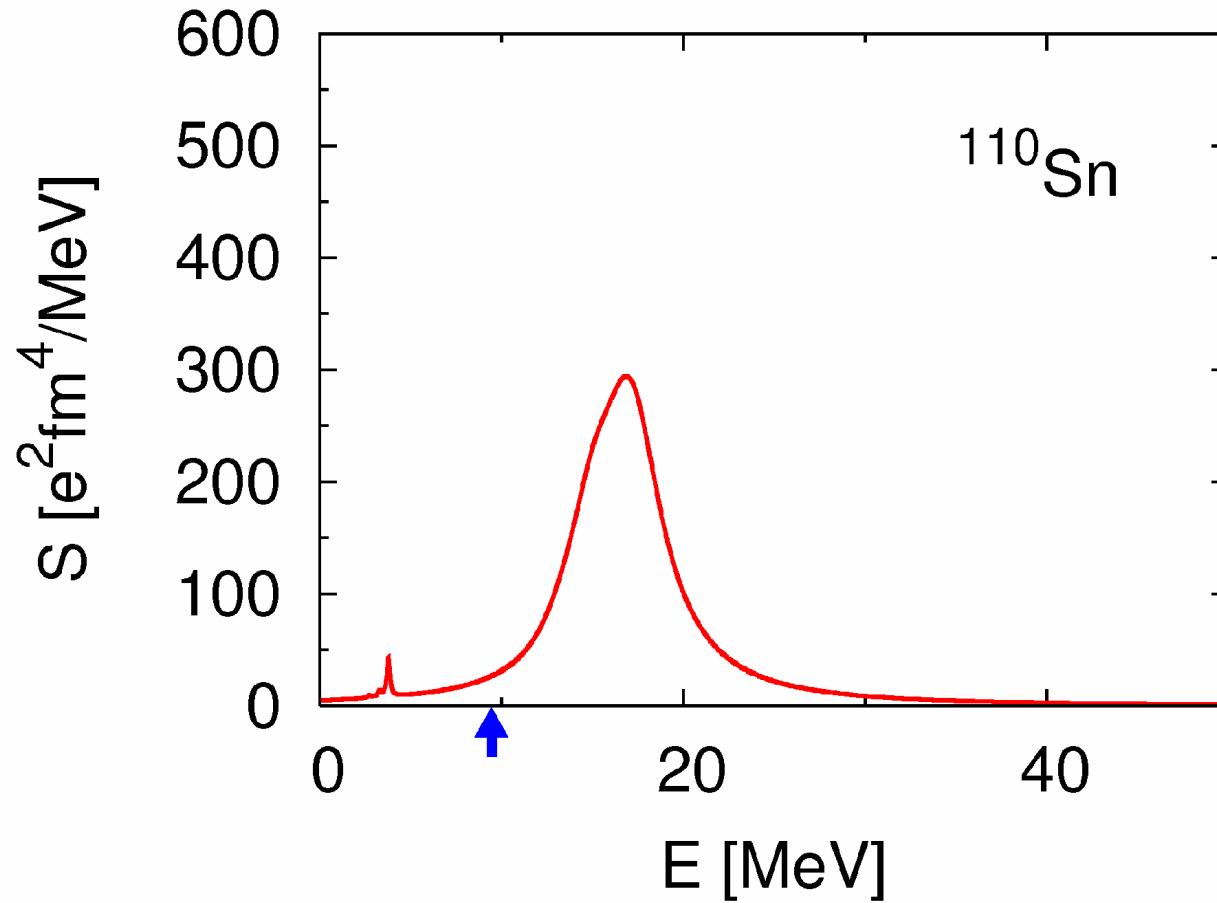
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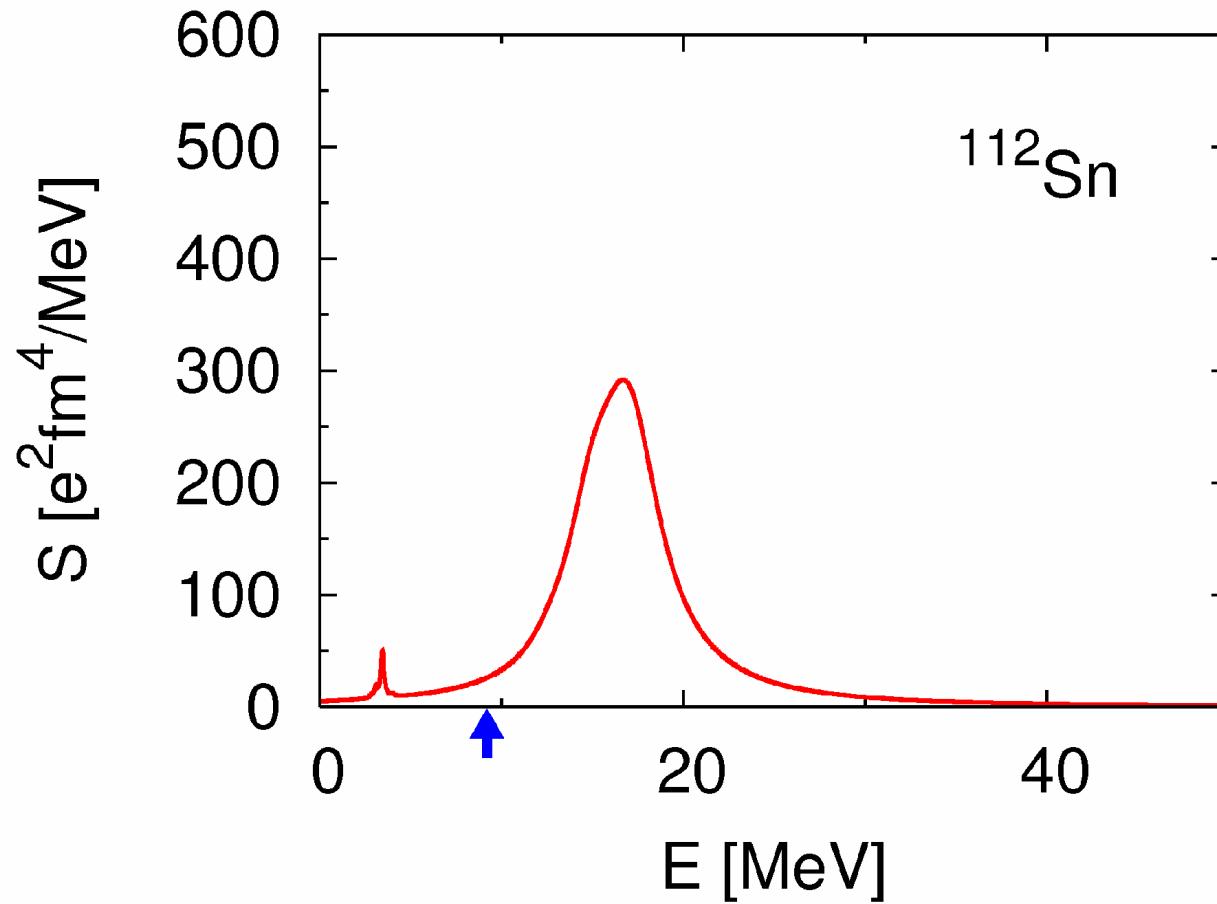
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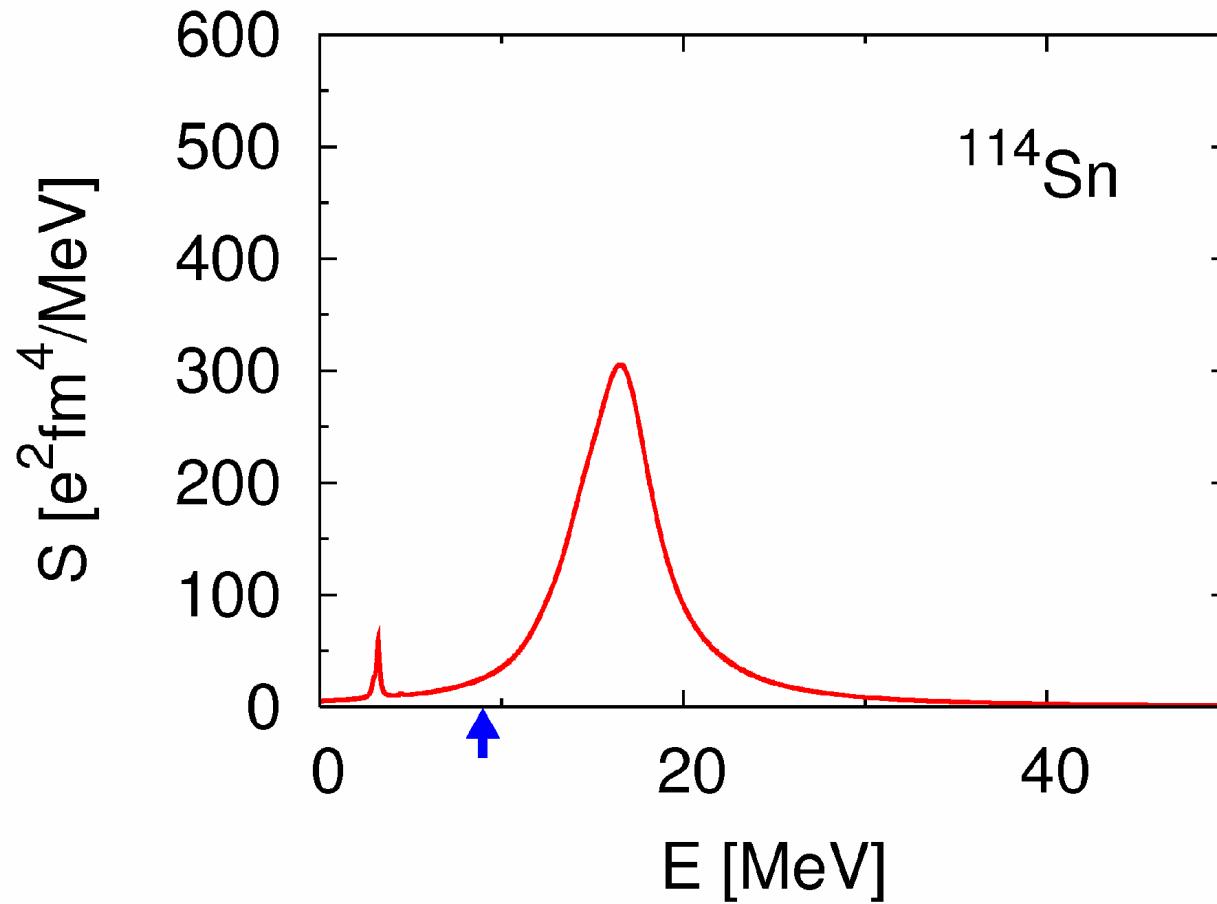
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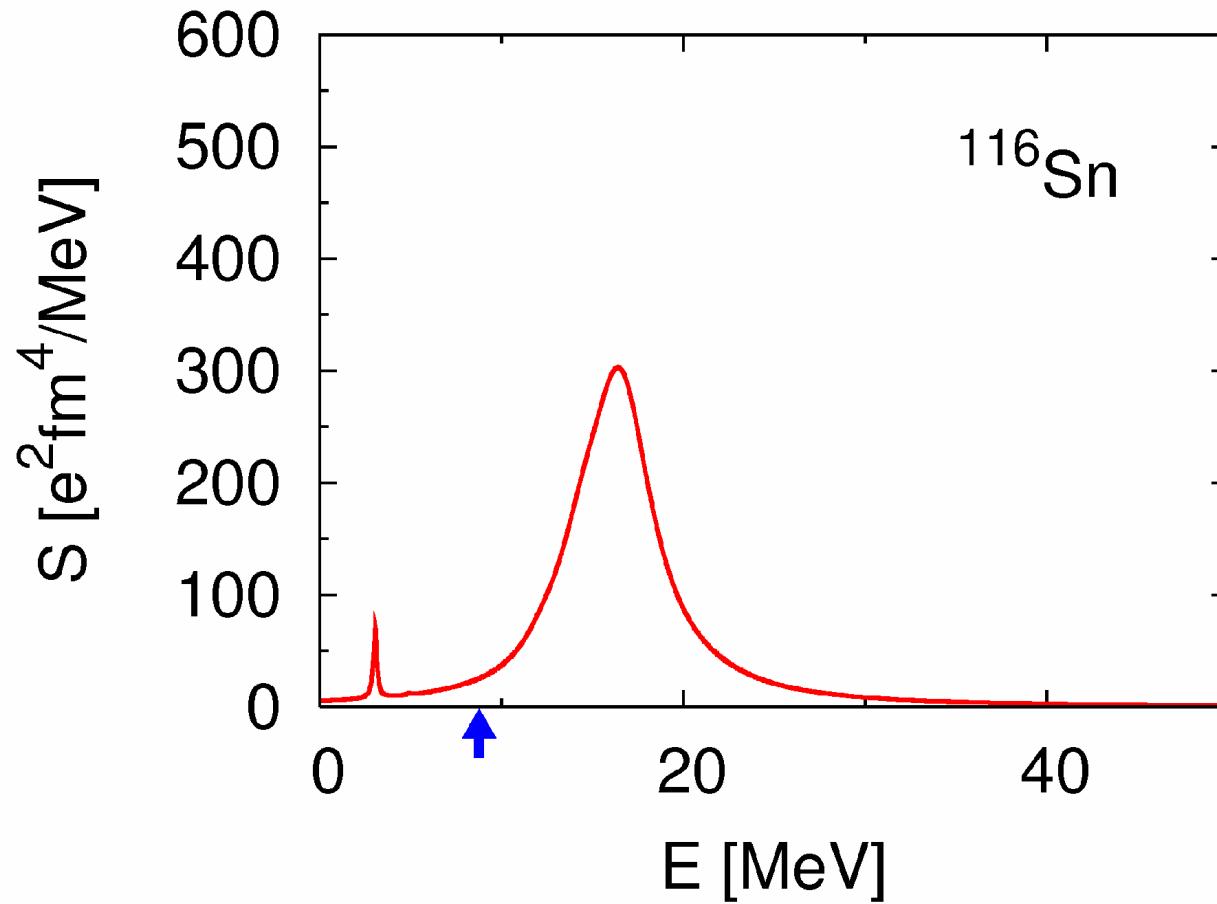
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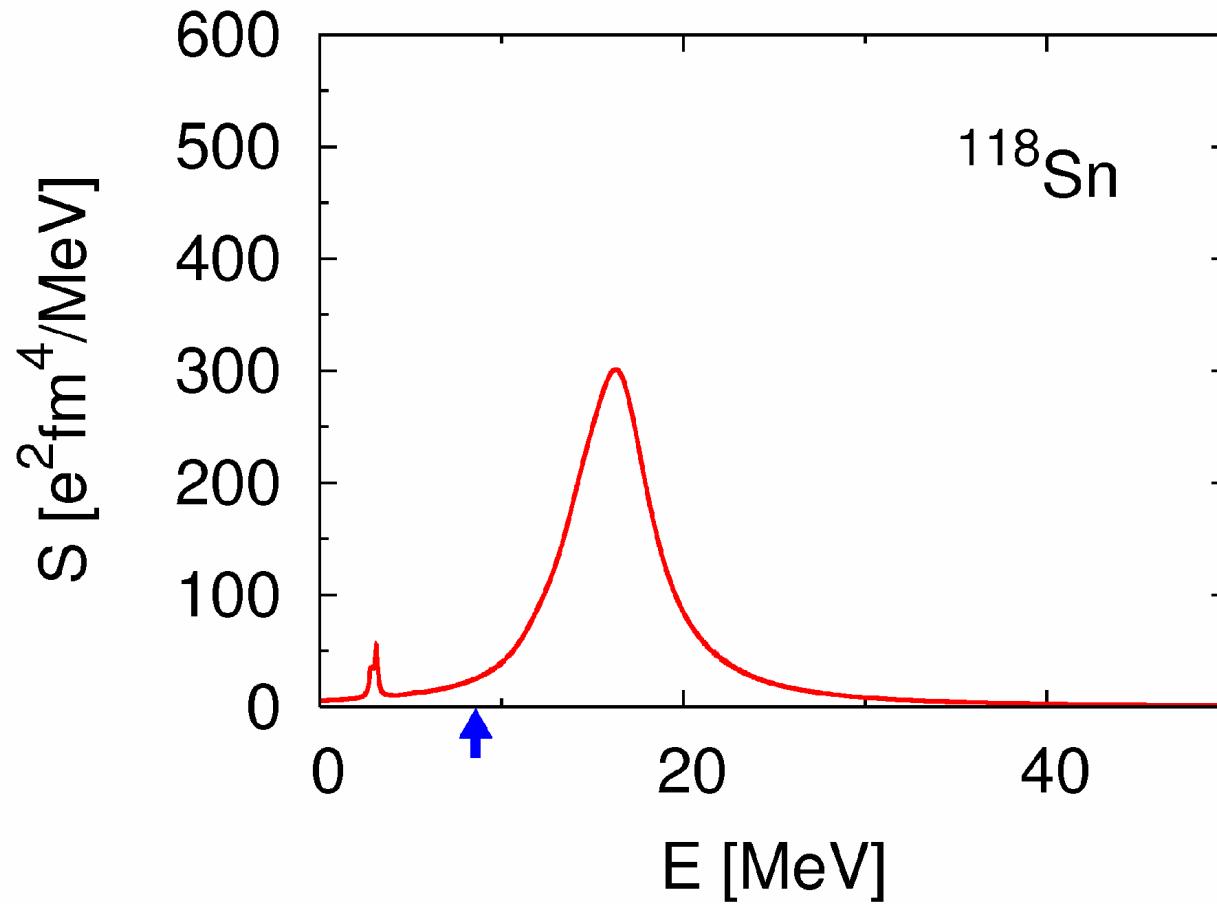
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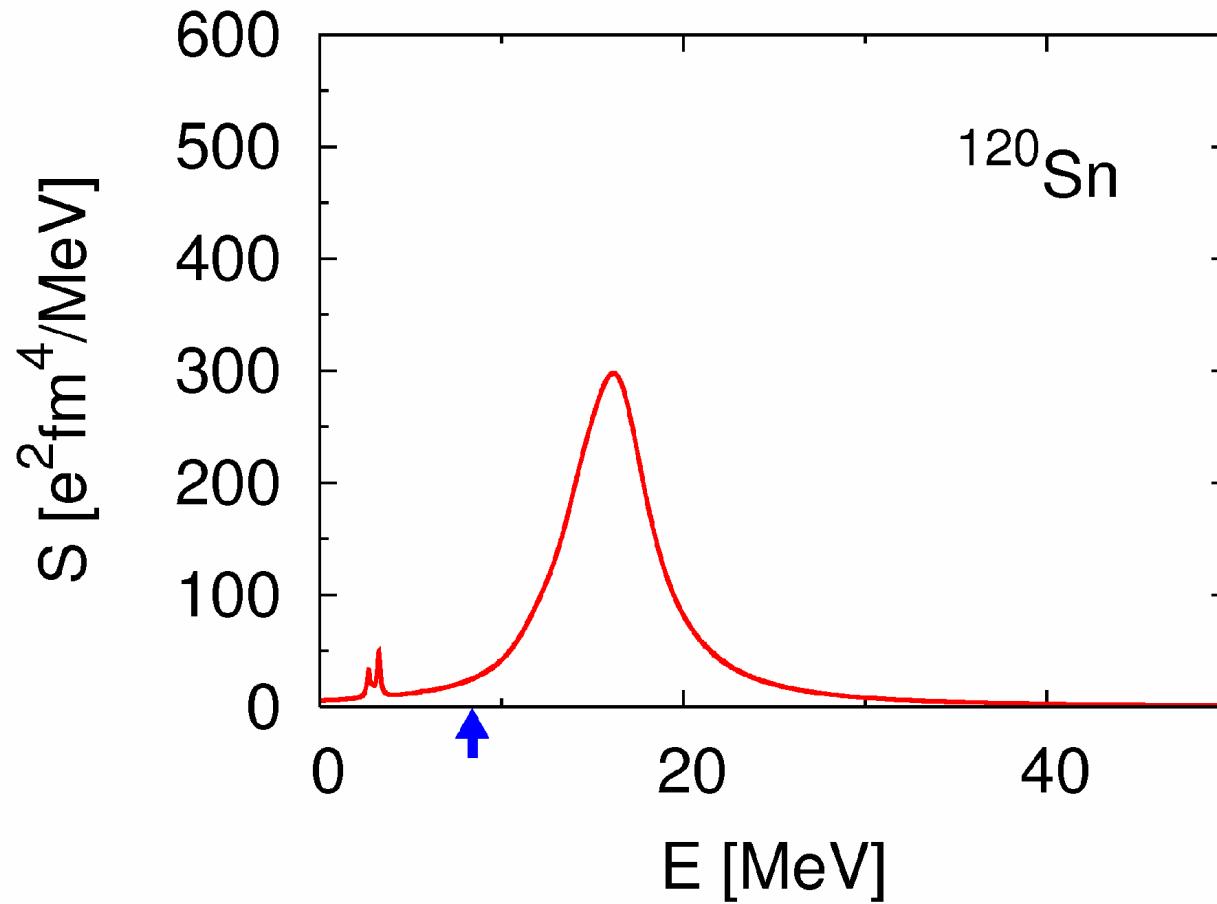
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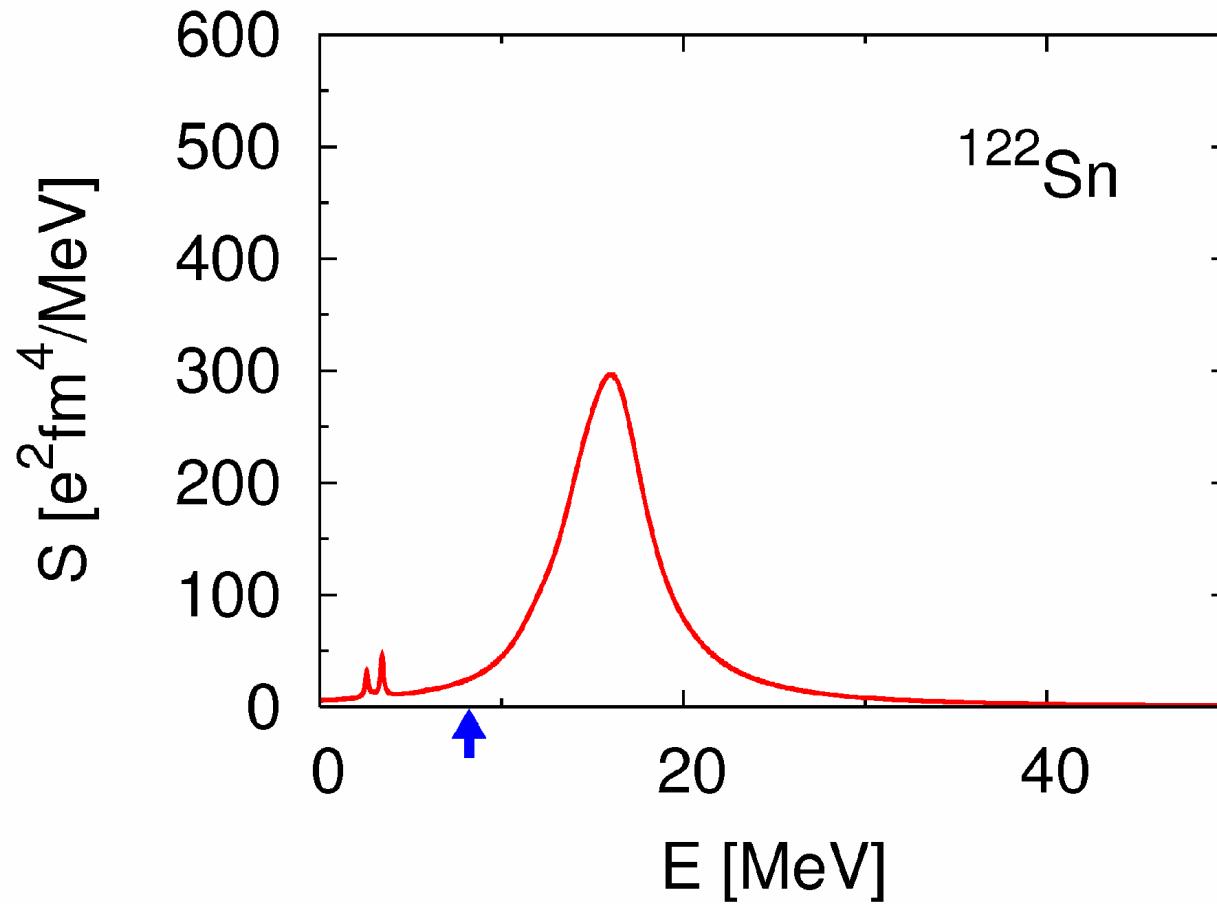
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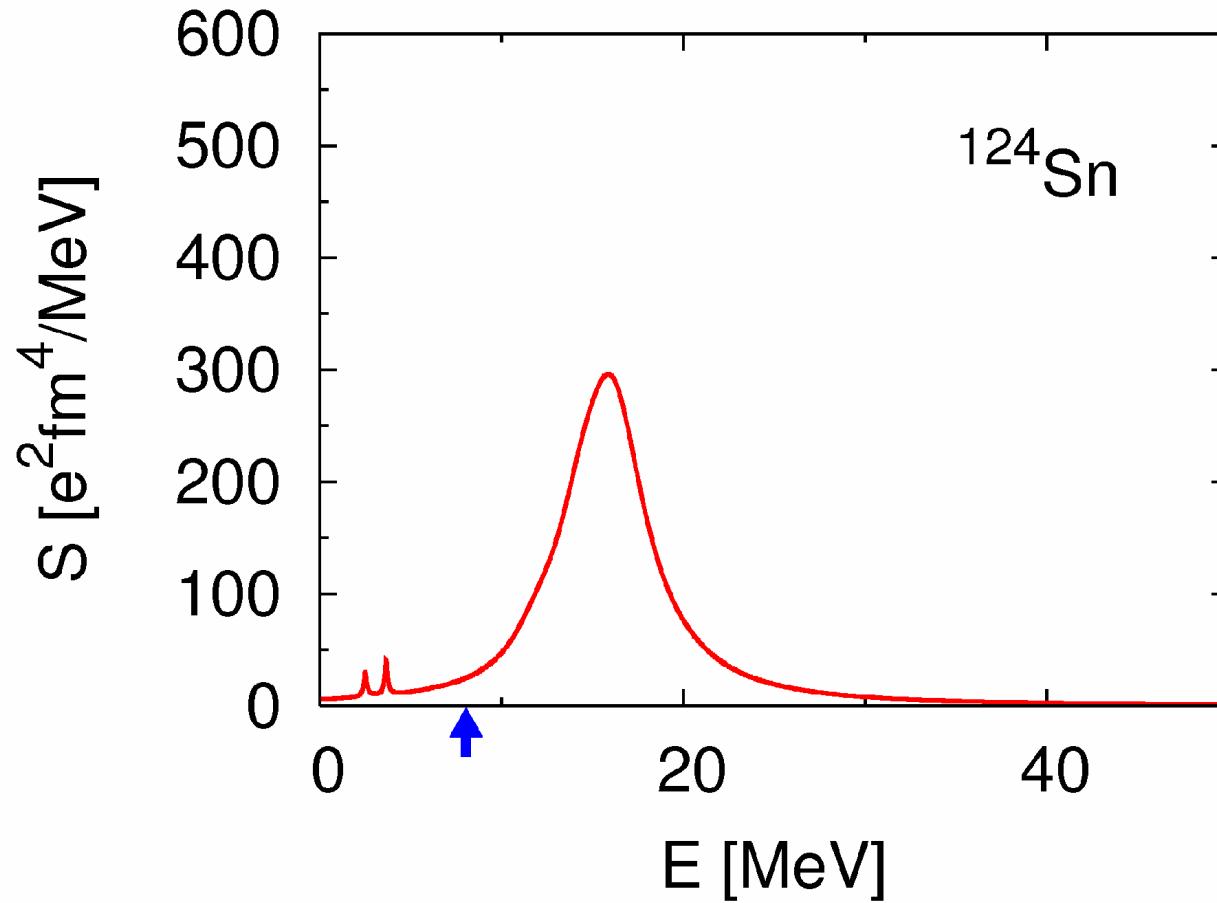
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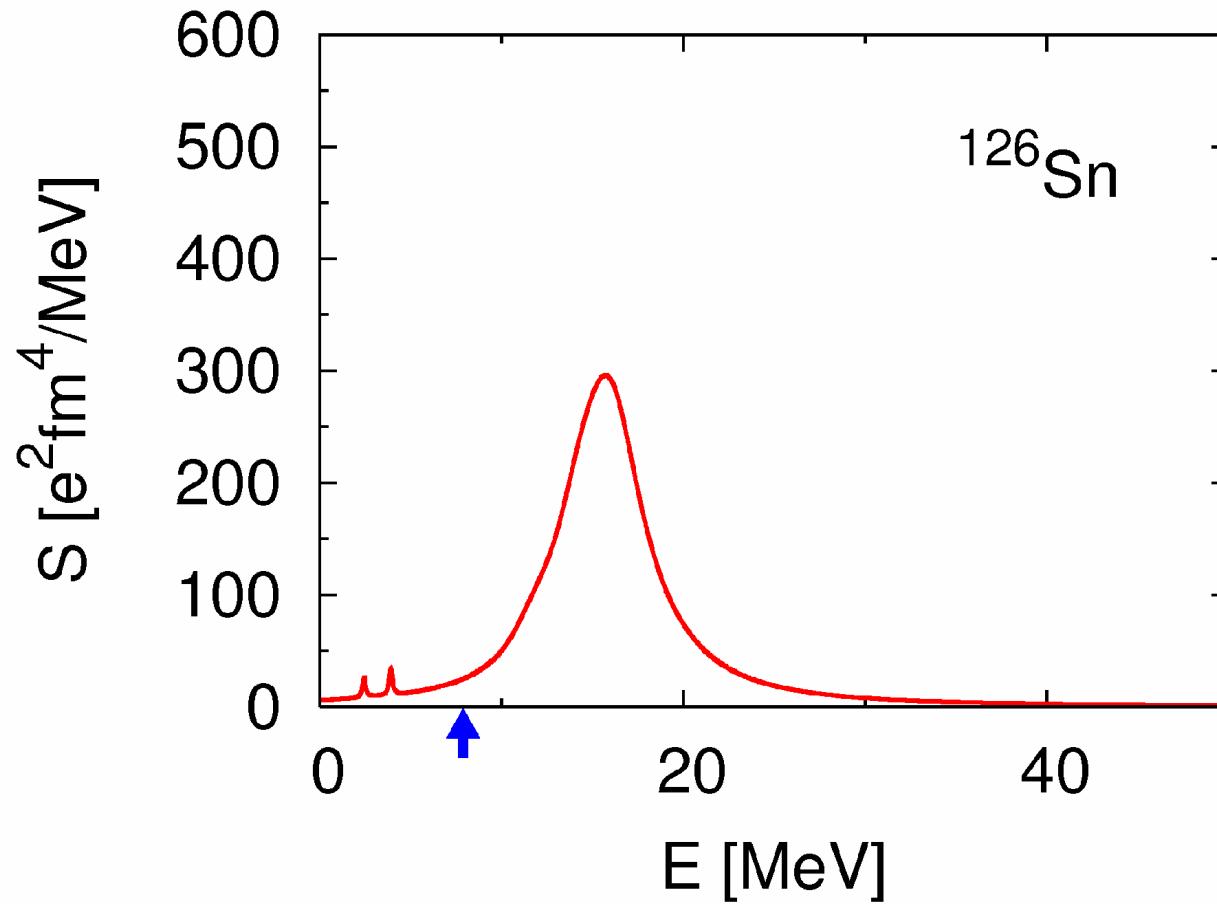
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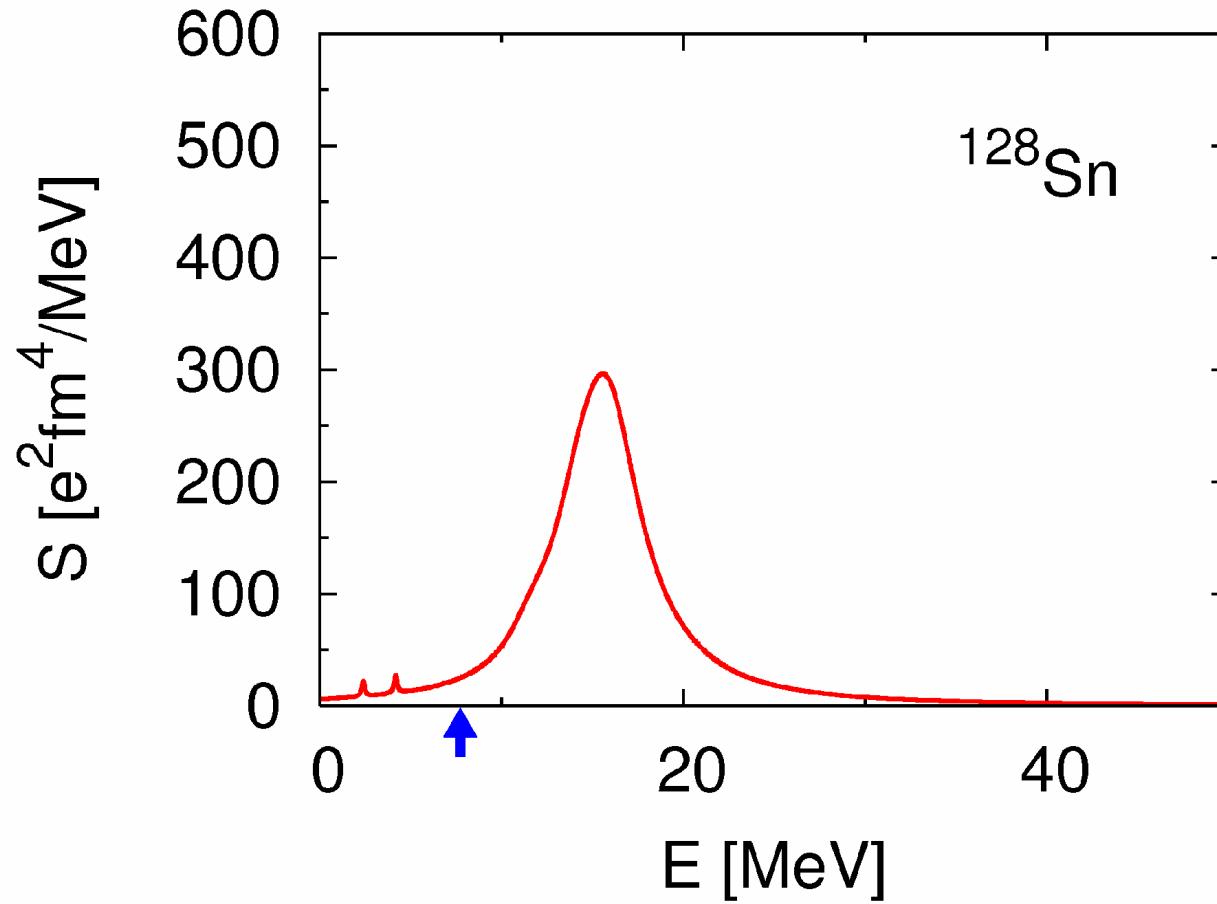
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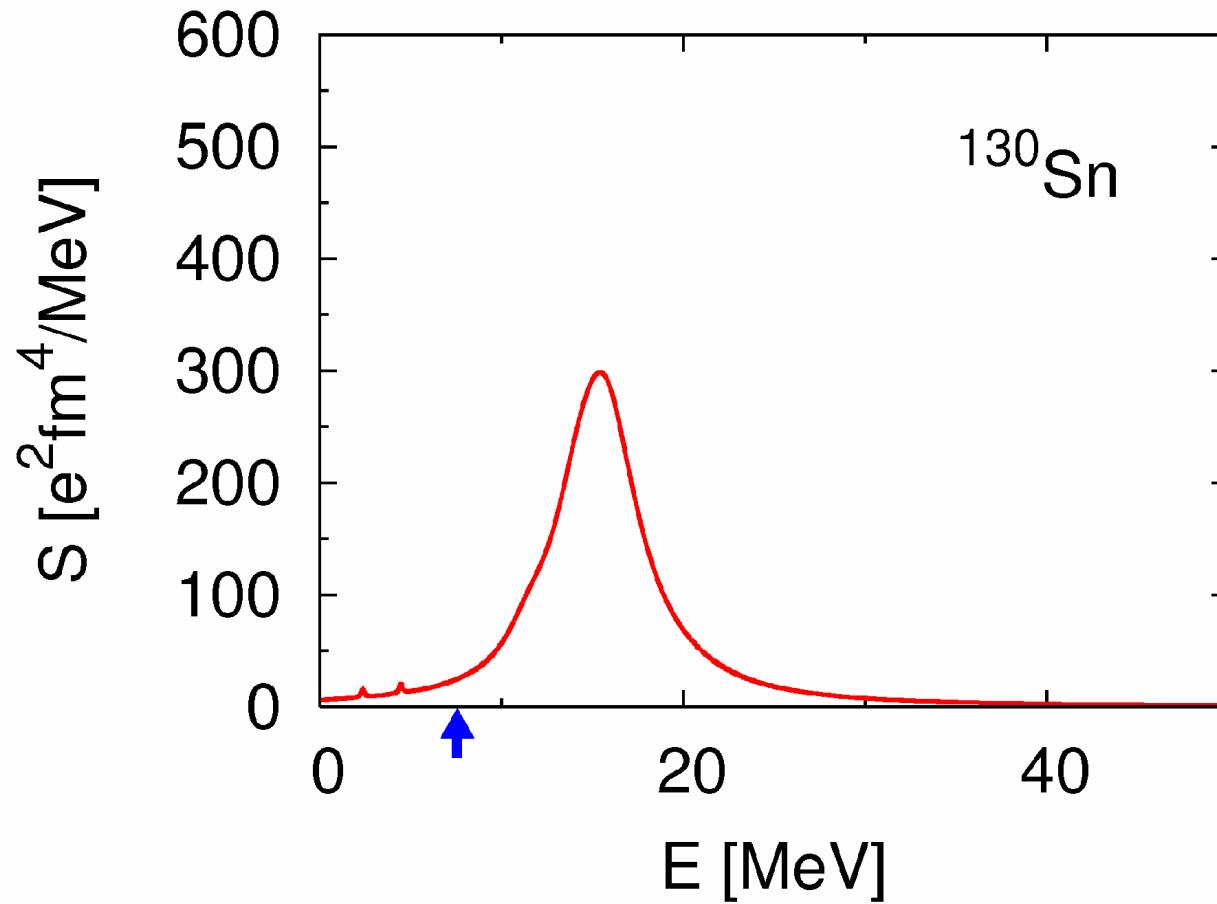
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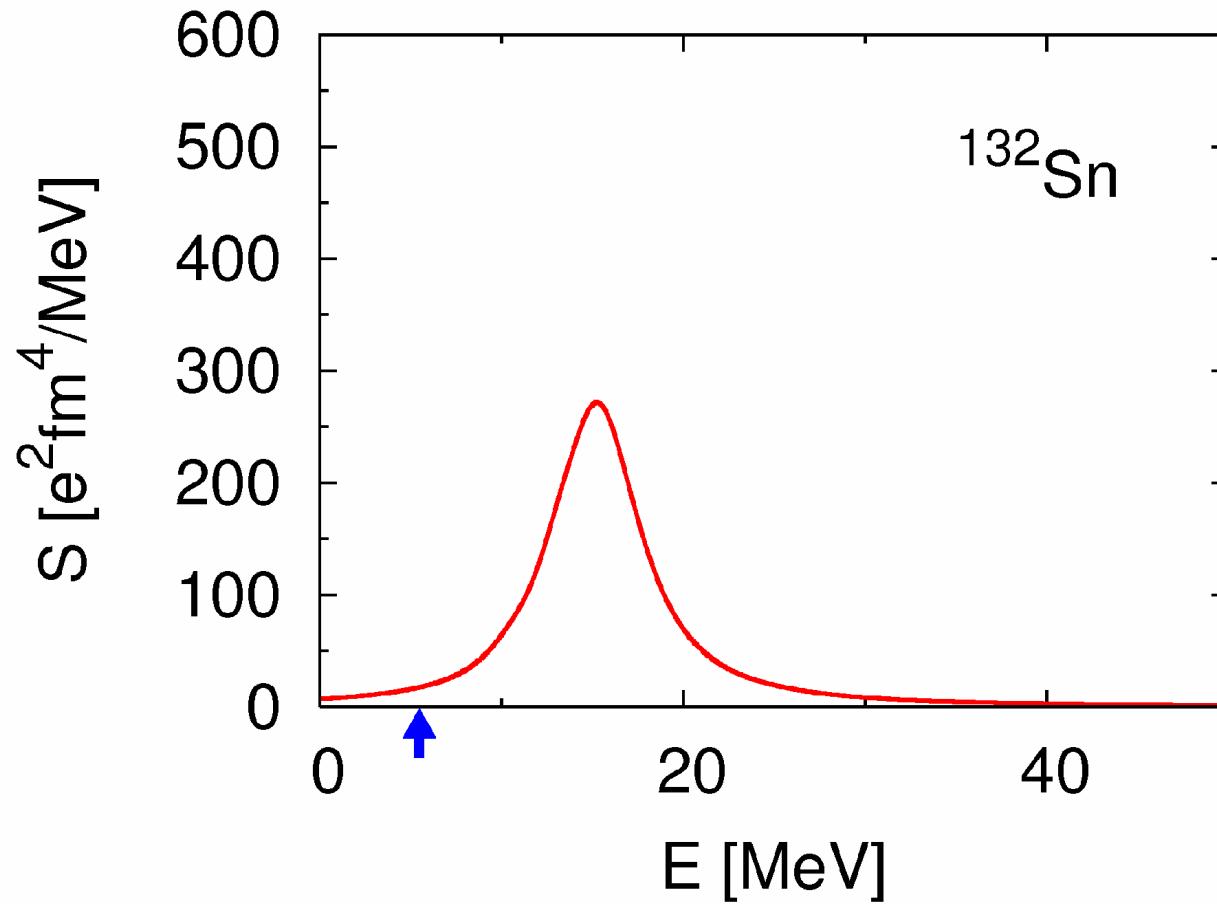
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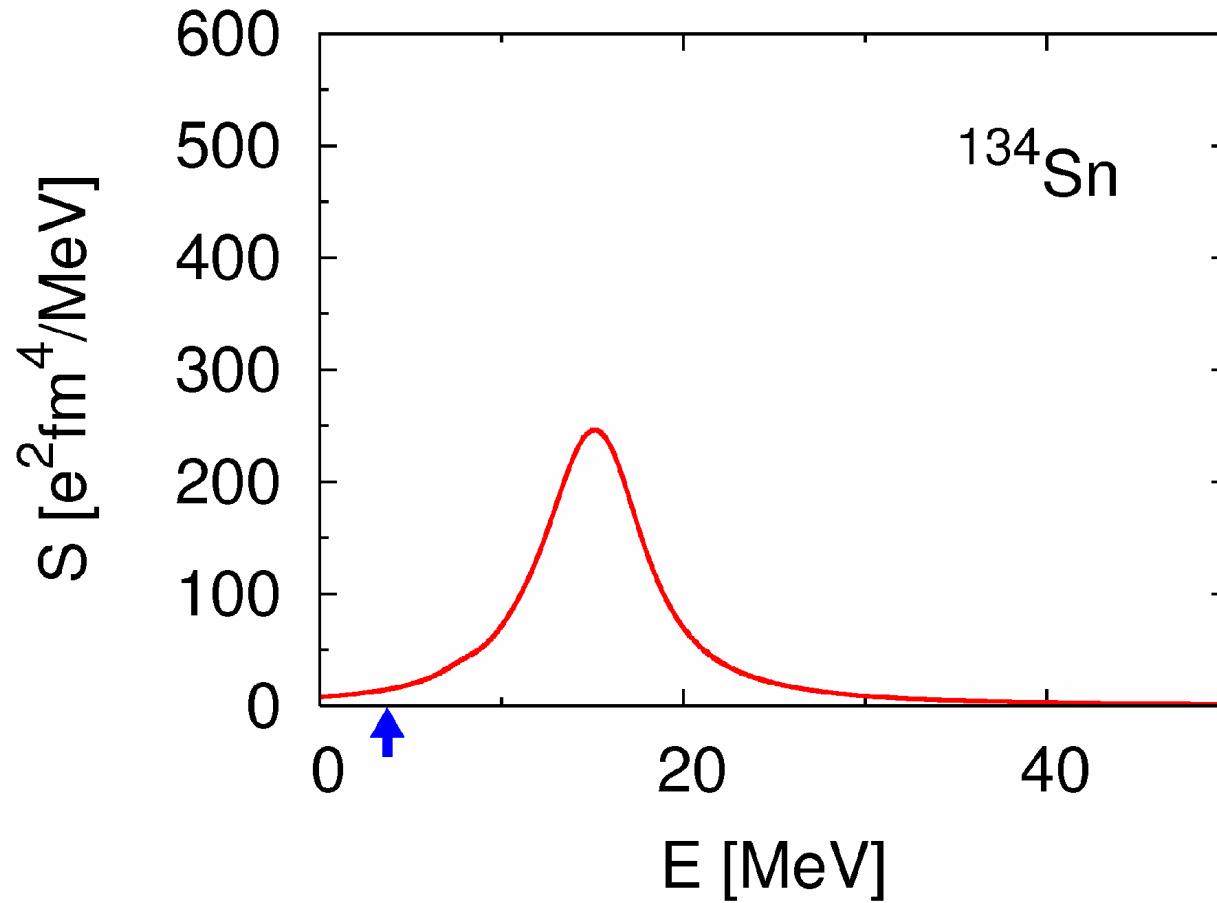
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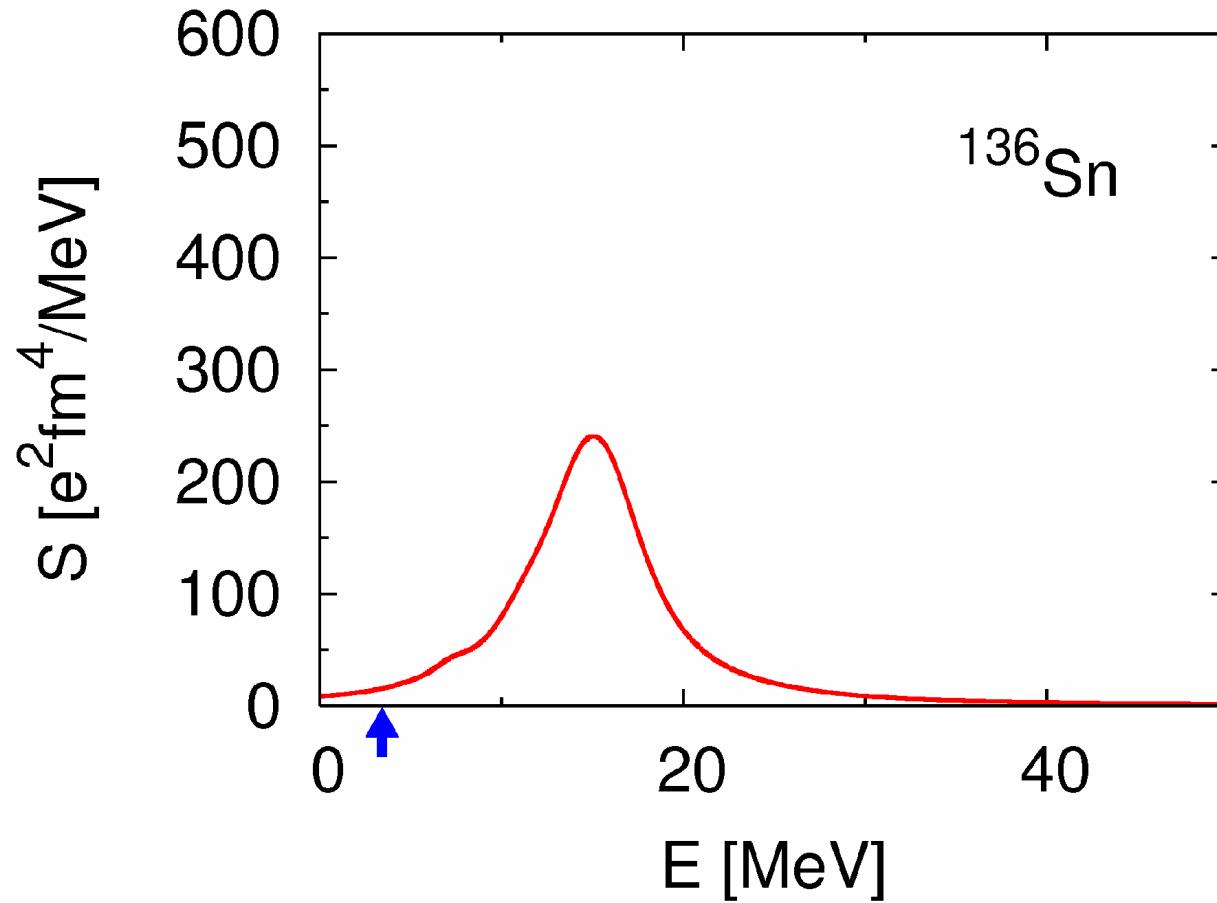
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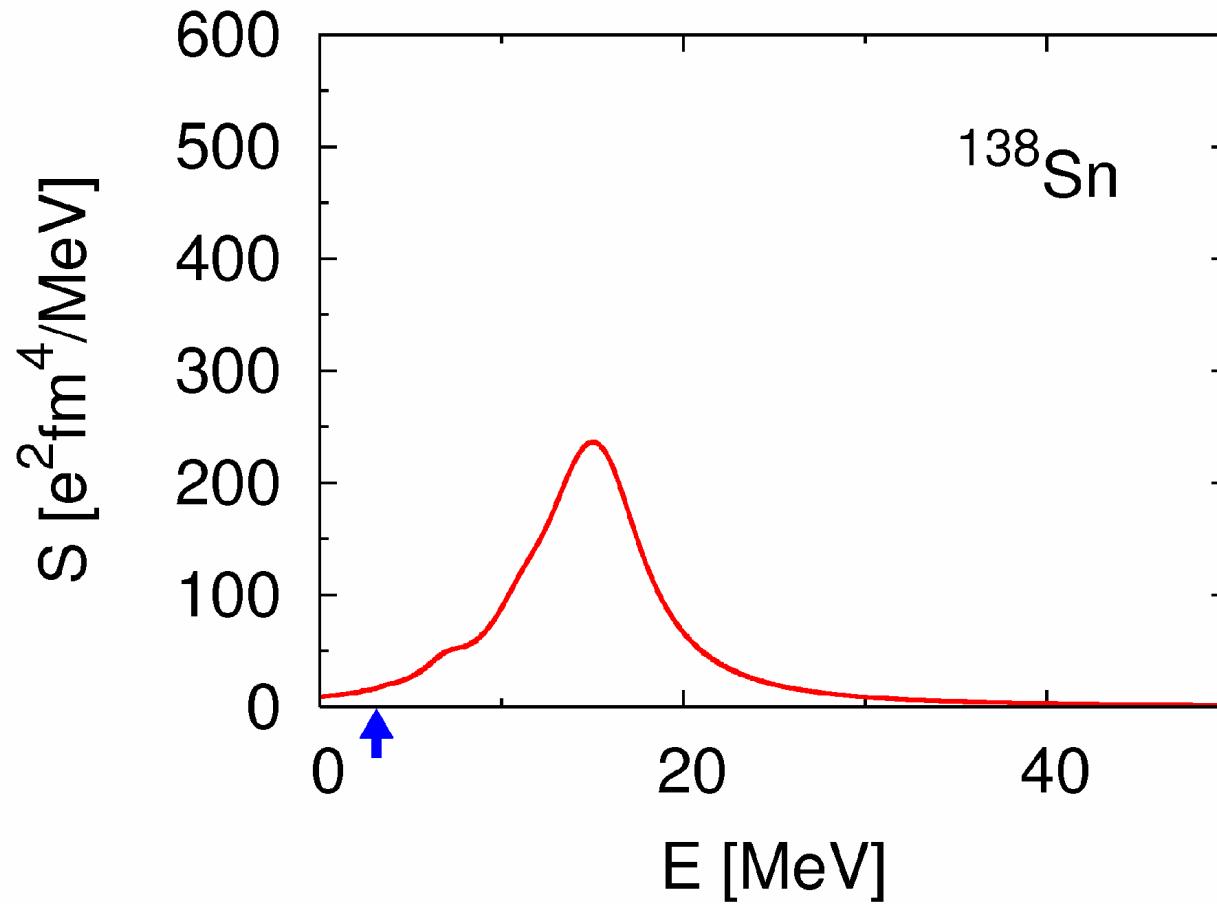
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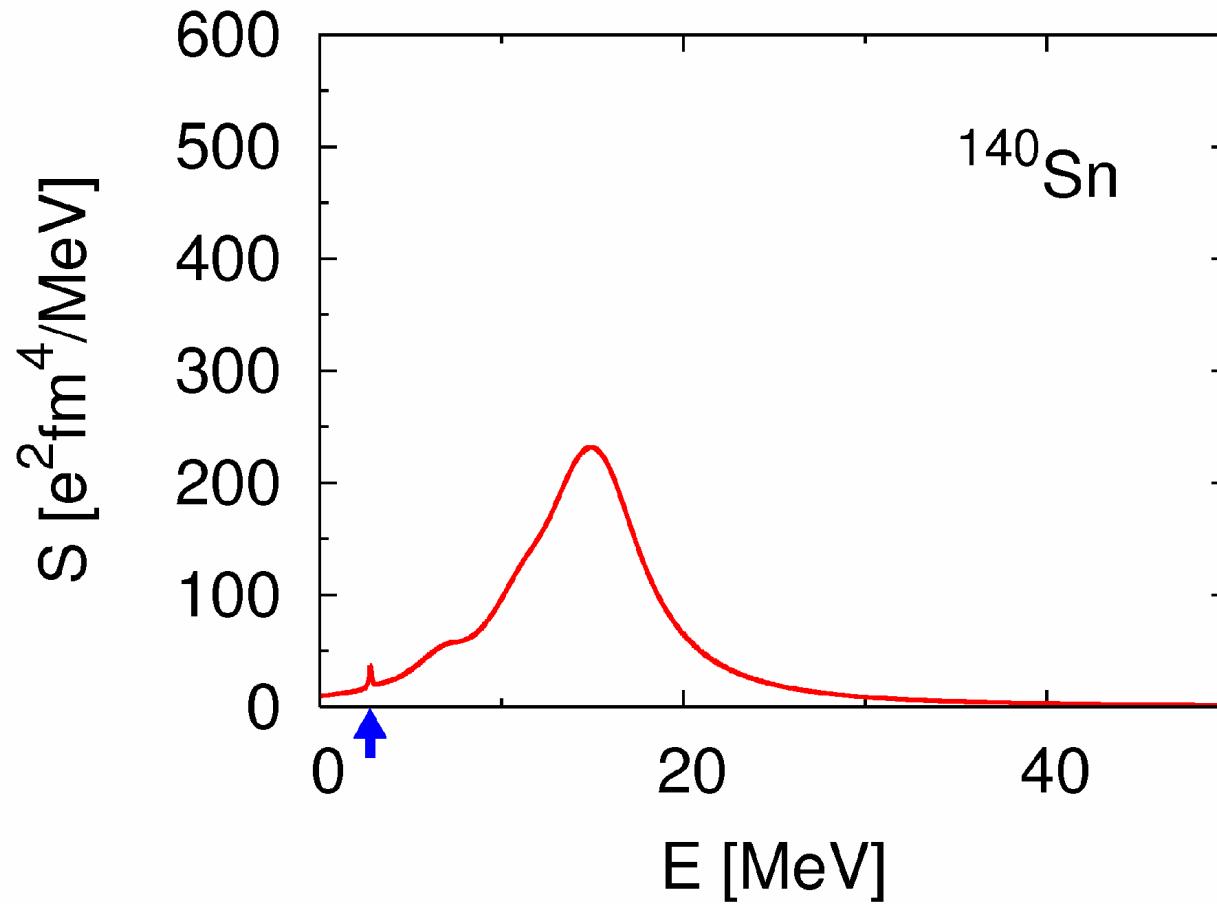
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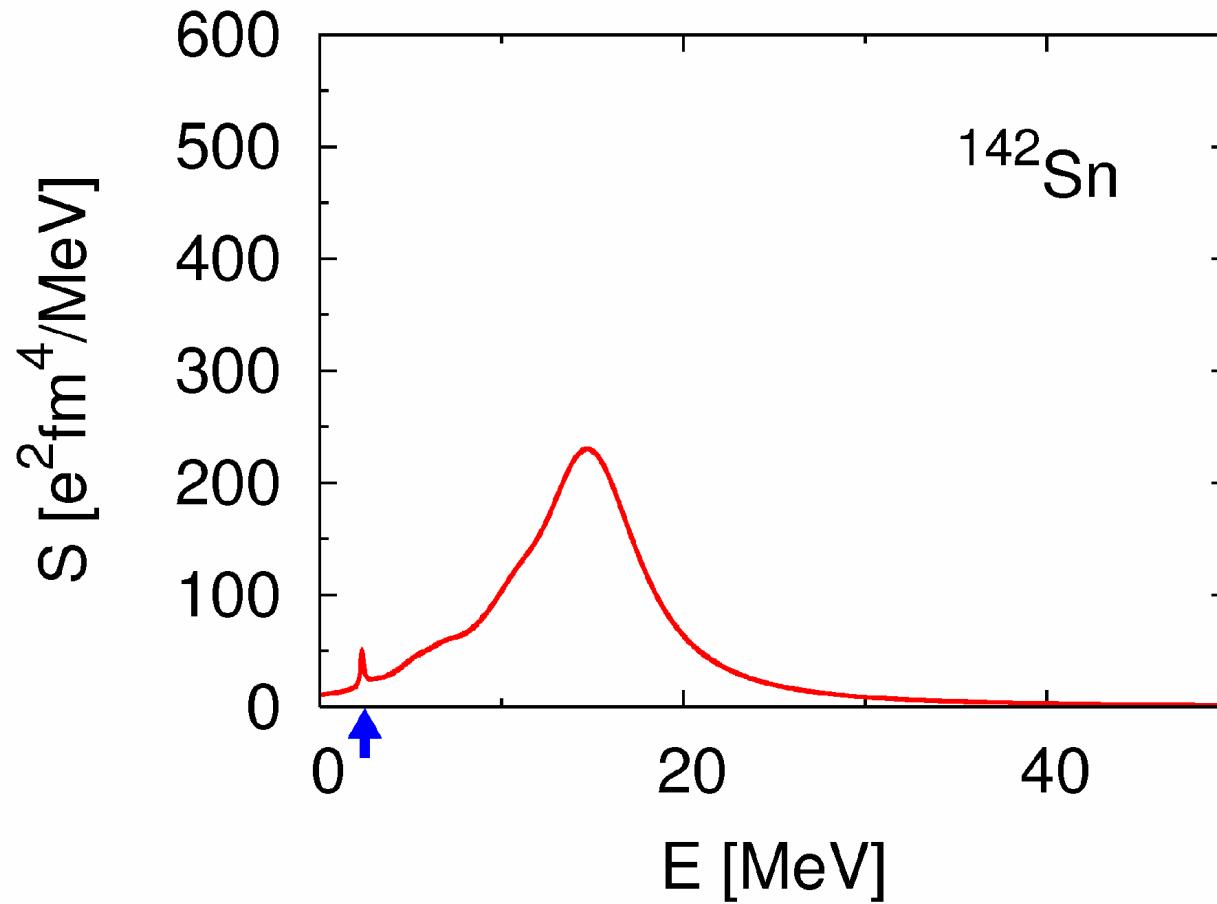
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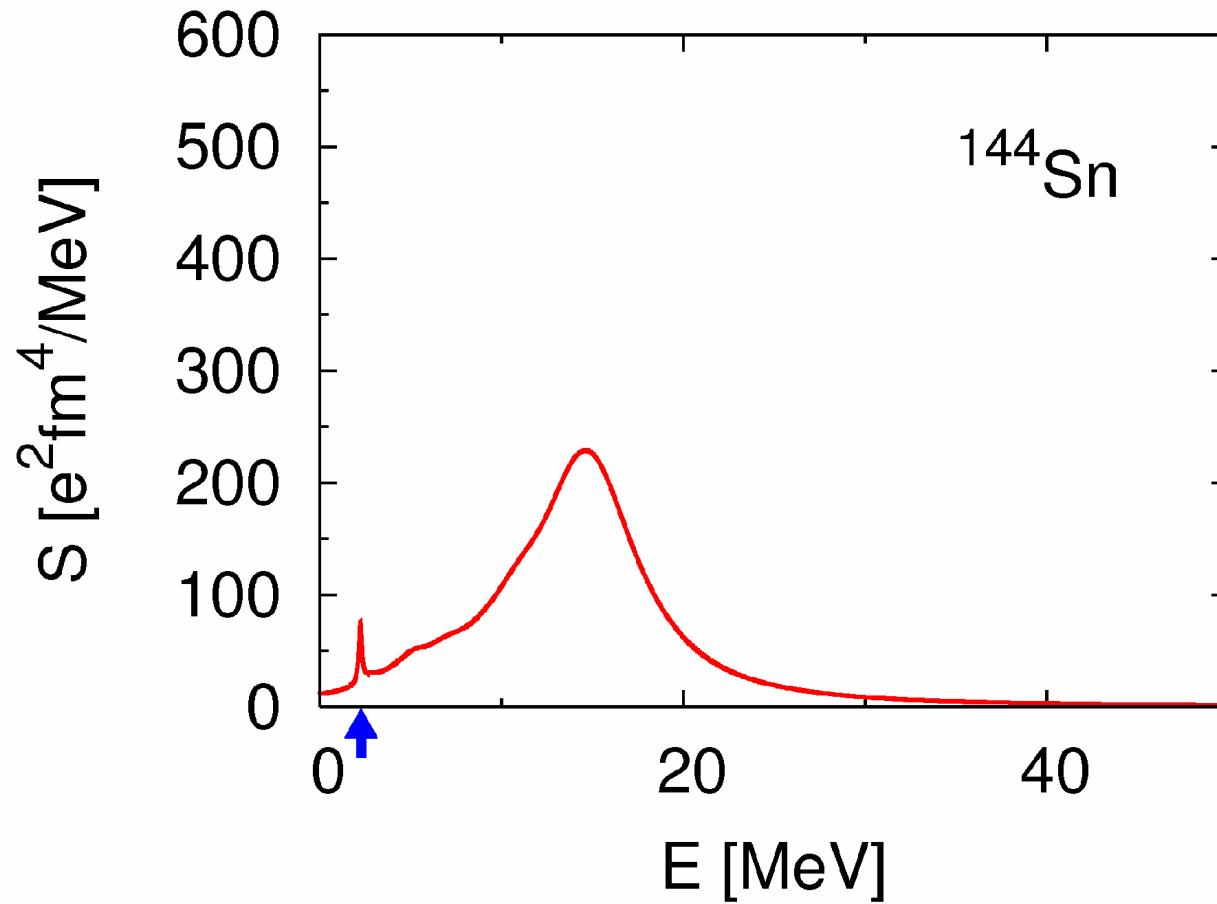
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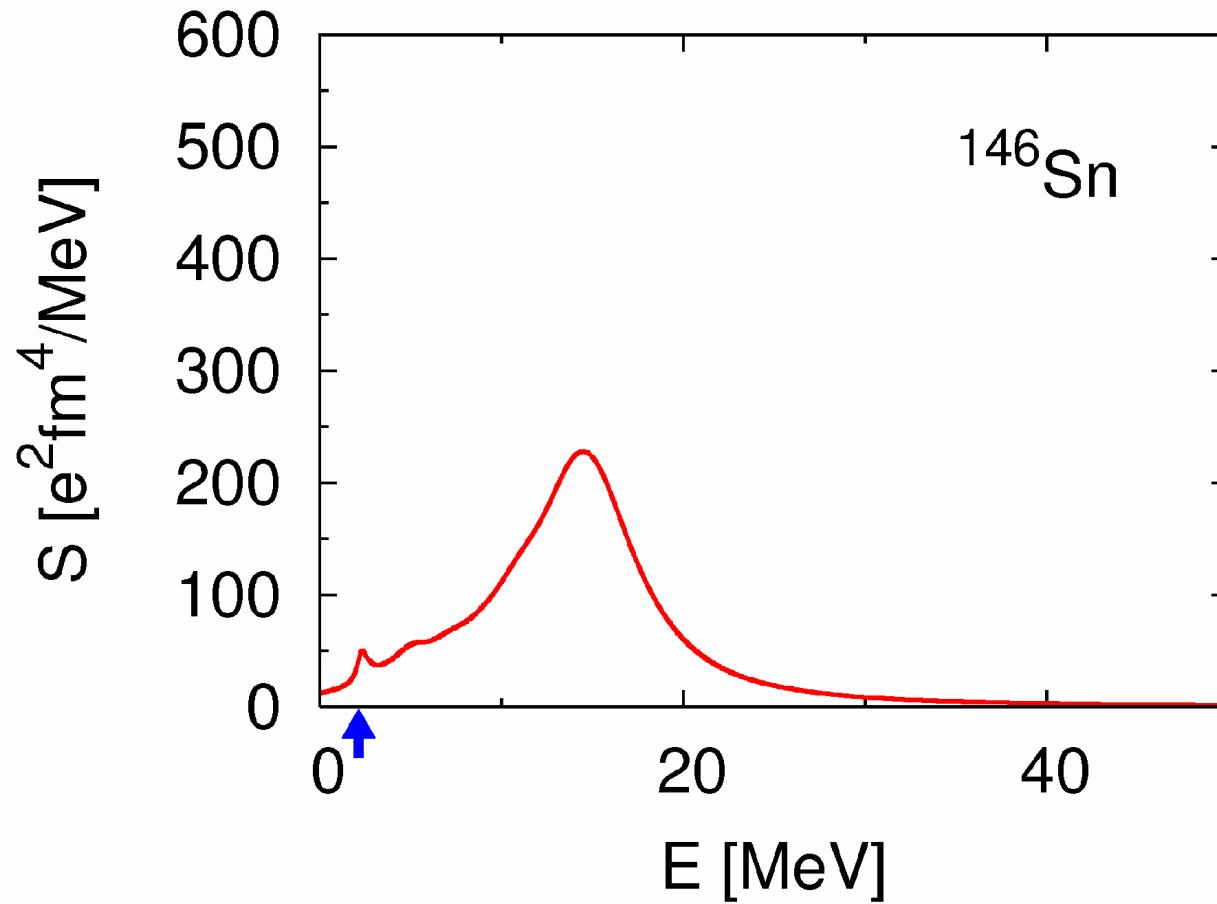
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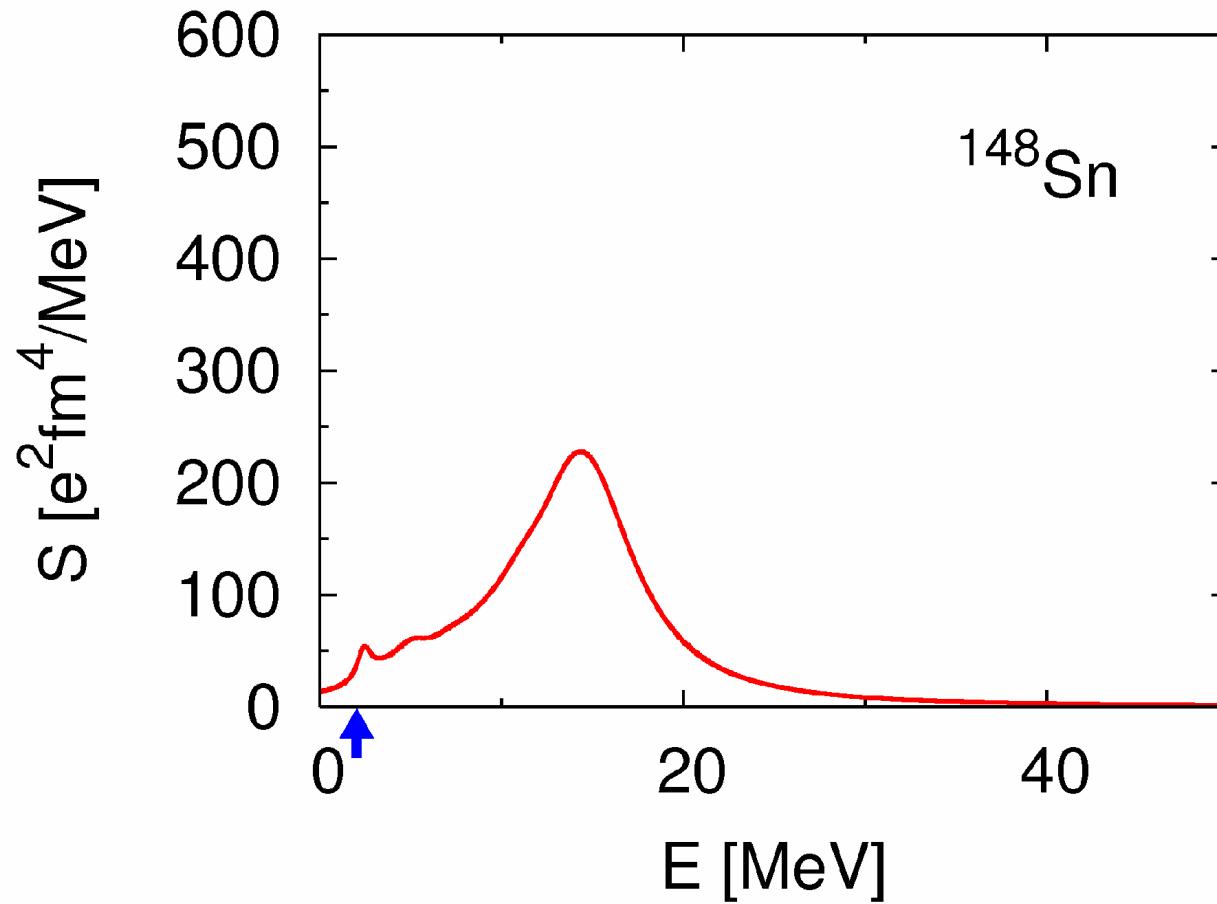
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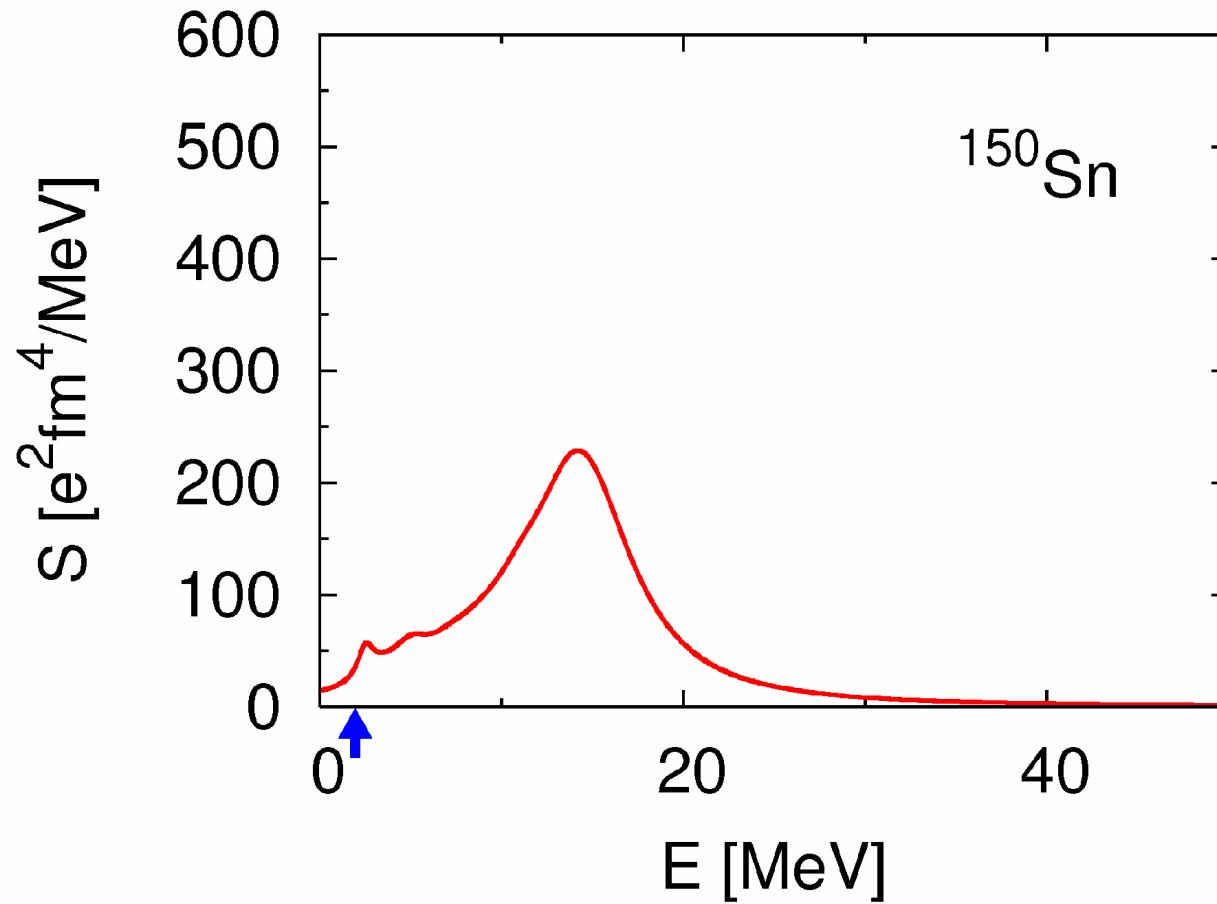
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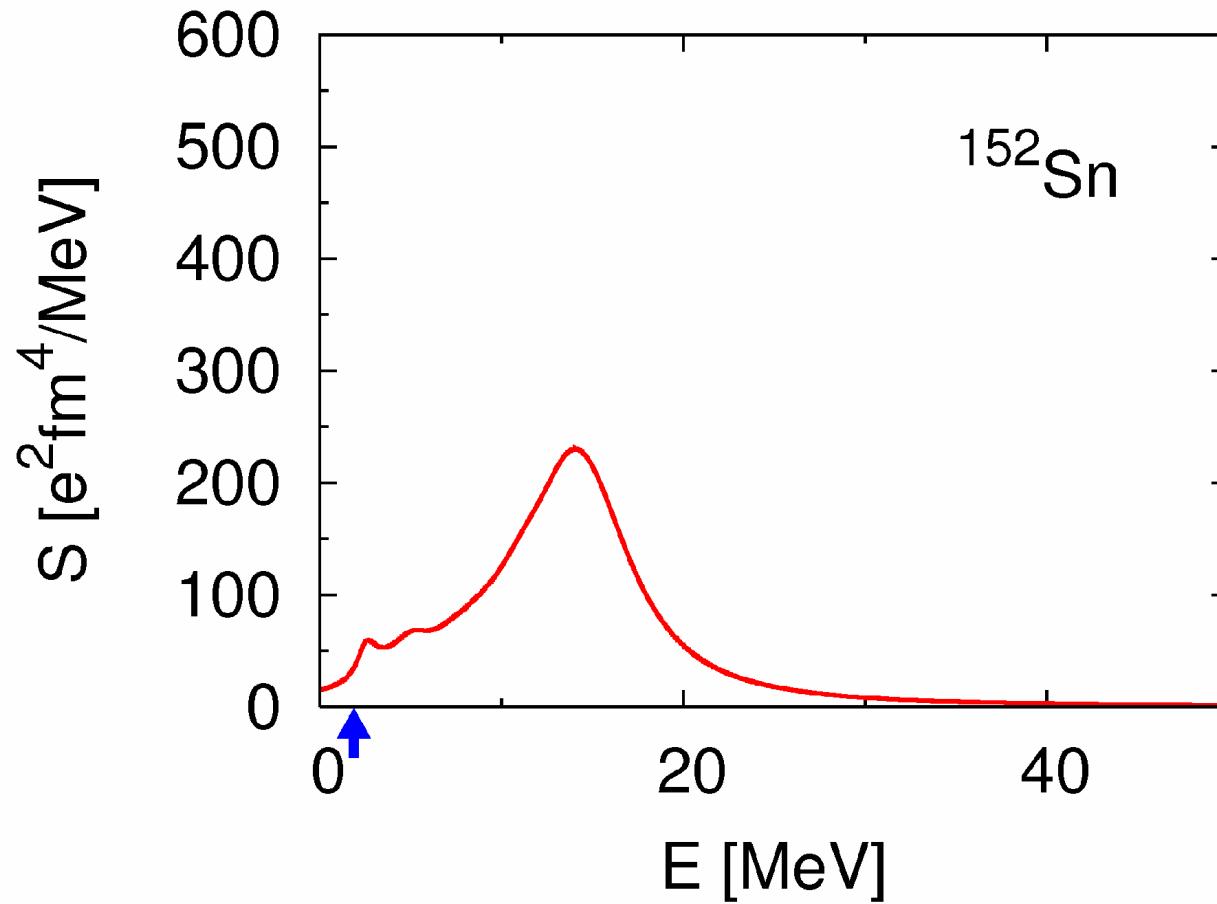
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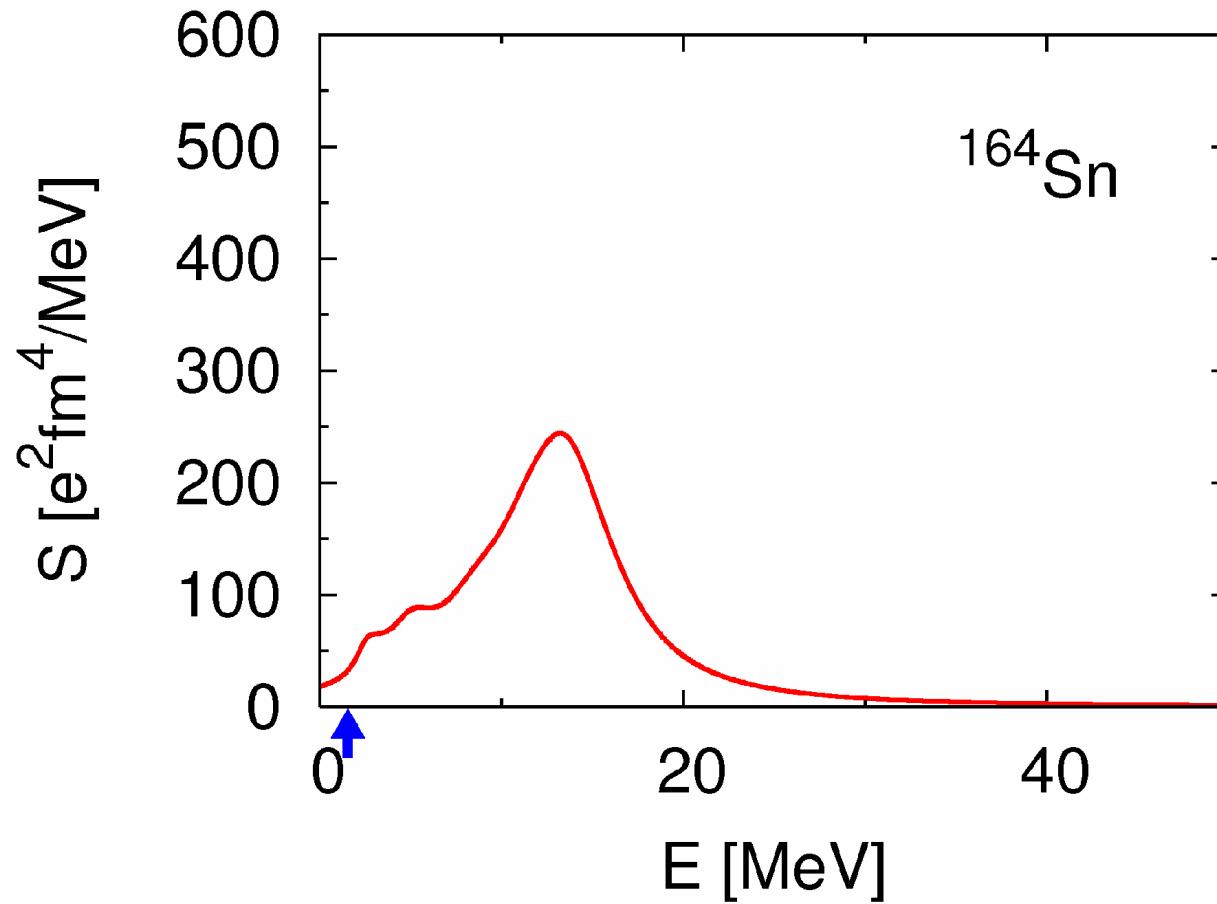


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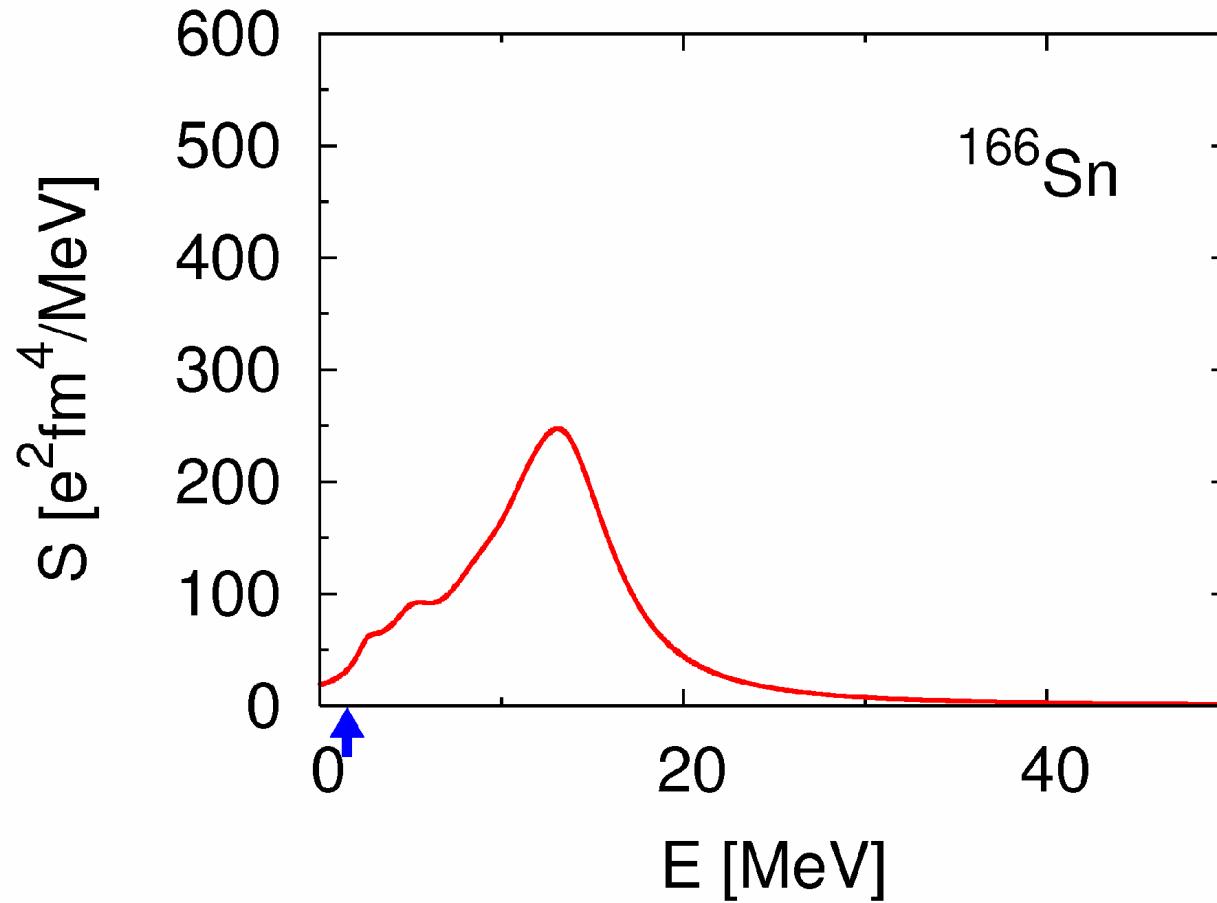


At  $A = 154 - 162$  ( $N = 104 - 112$ )  
ground states : deformed

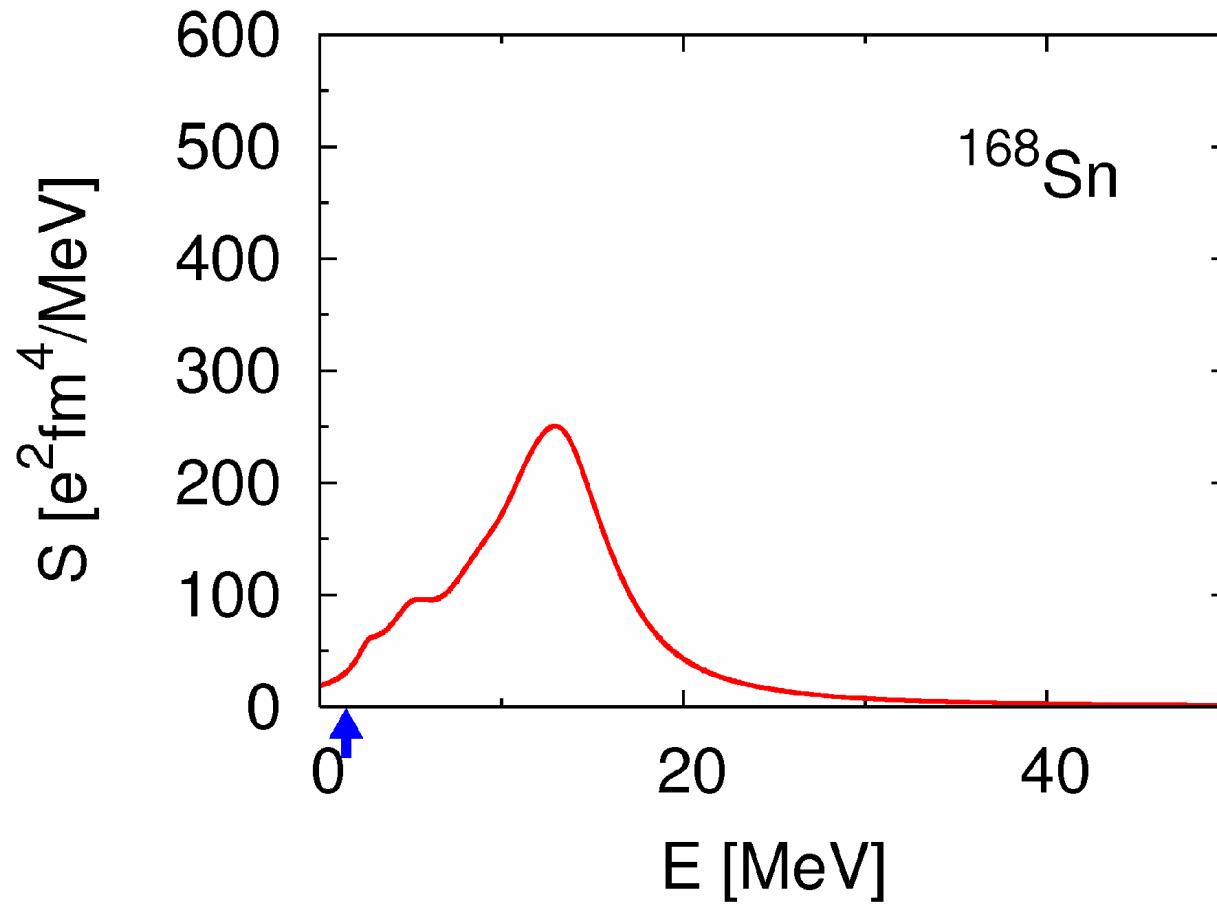
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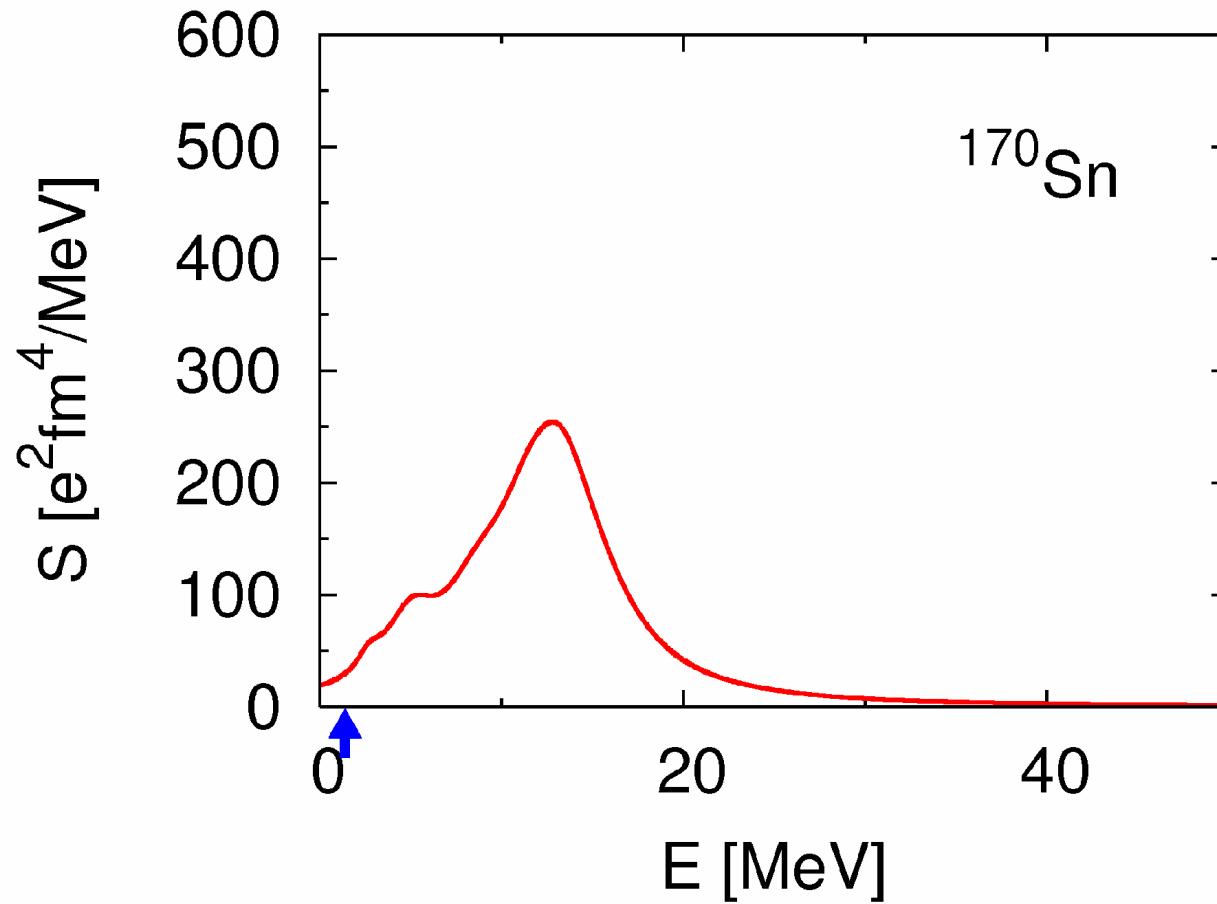
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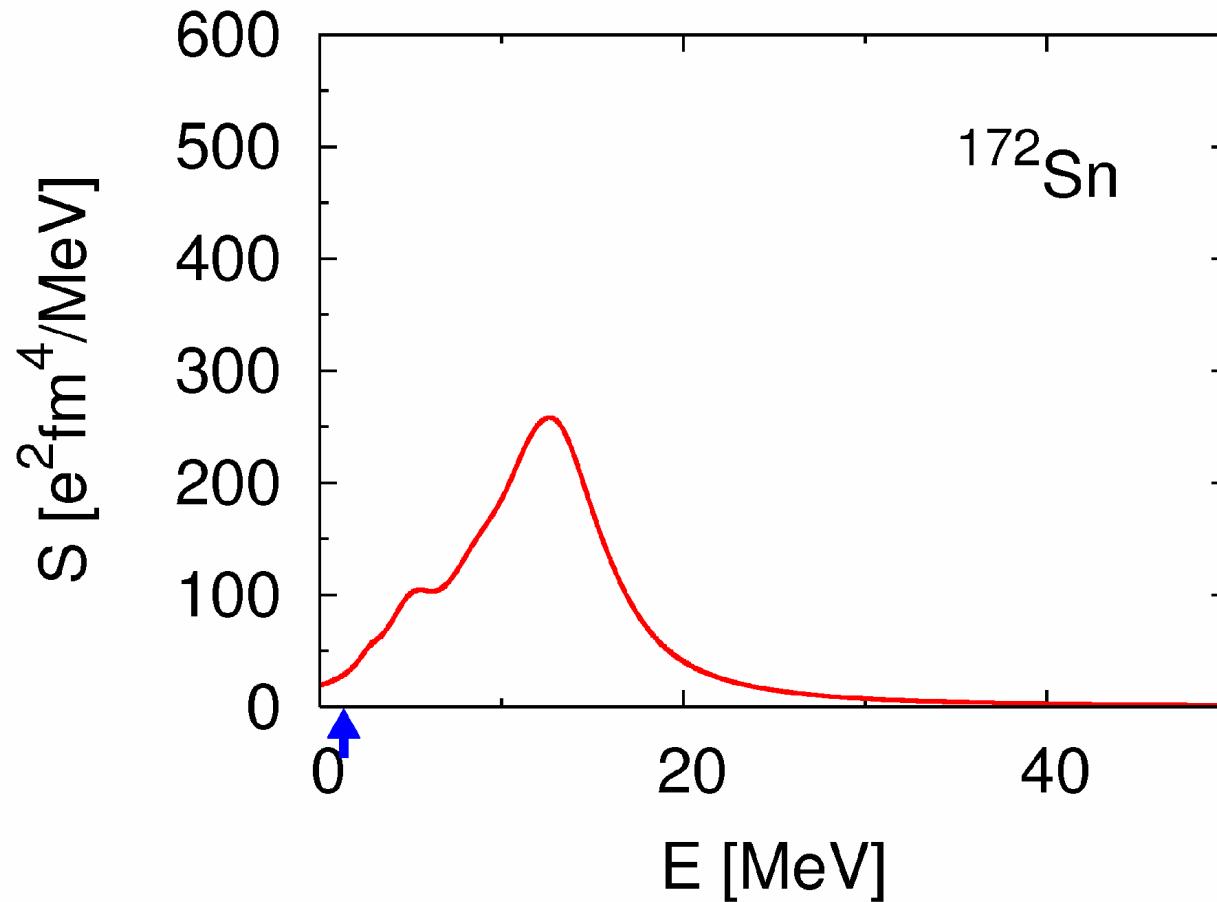
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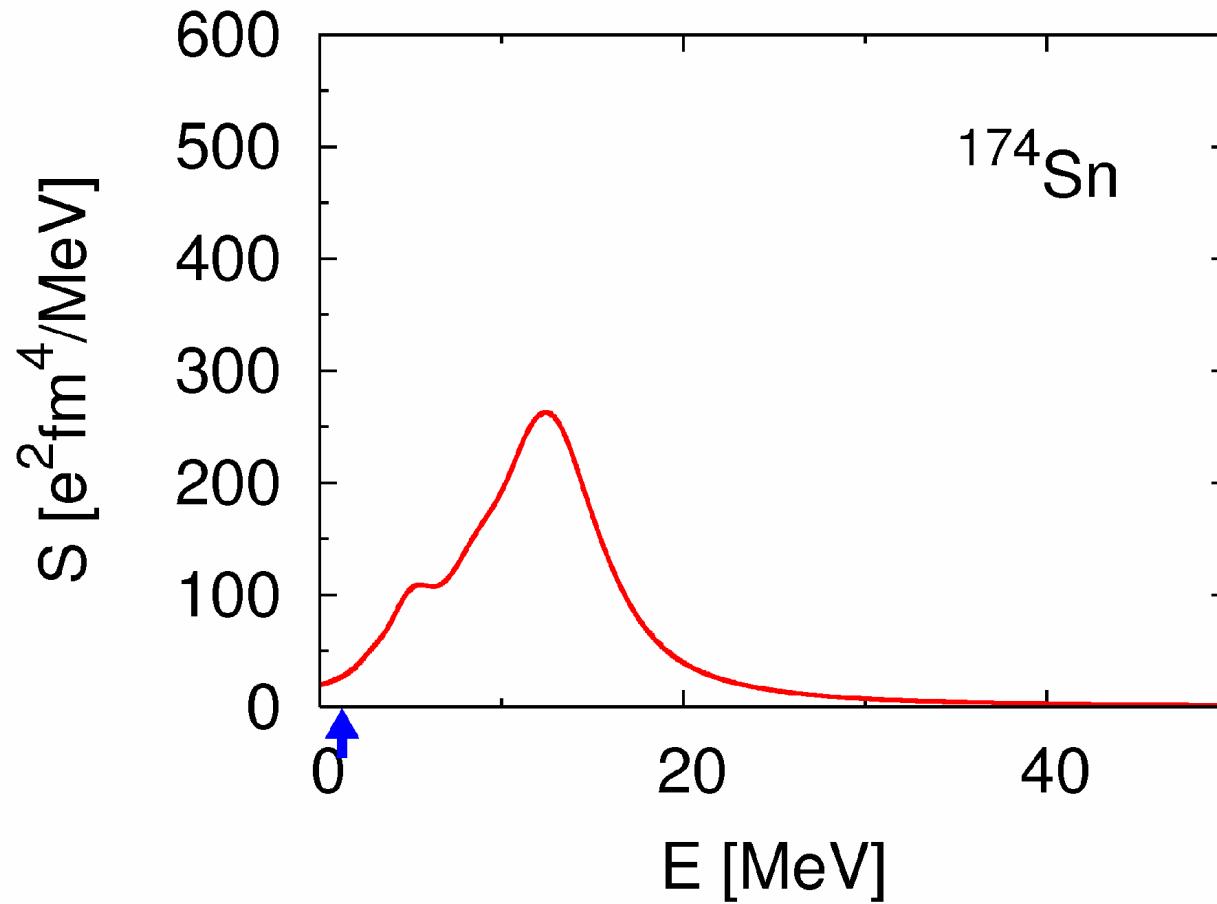
## Isoscalar $0^+$ strength functions



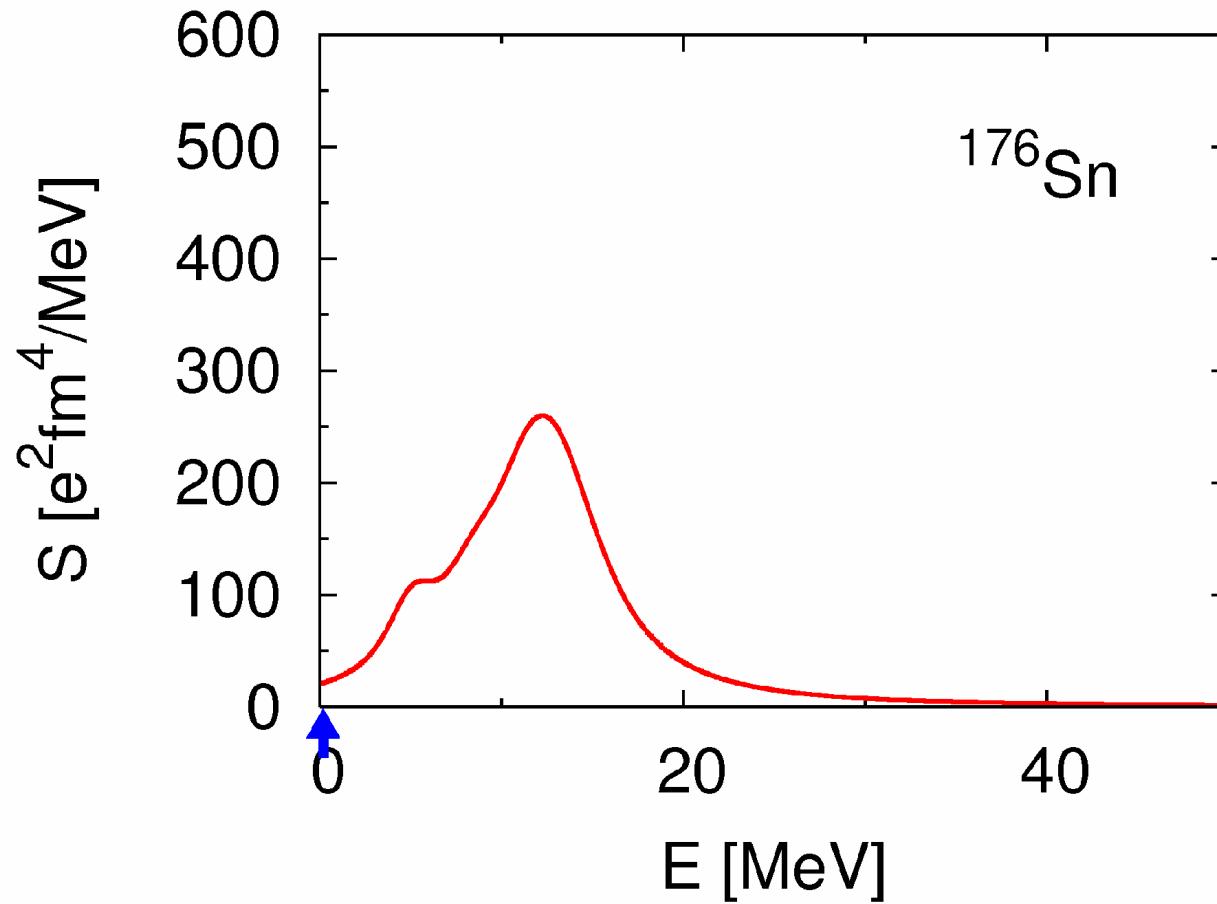
## Isoscalar $0^+$ strength functions



## Isoscalar $0^+$ strength functions



## Isoscalar $0^+$ strength functions



## Exp. data

S.Shlomo and D.H. Youngblood, Phys.Rev.C, **47**, 529 (1993) Tab.III

However,

D.H. Youngblood et al., Phys.Rev.Lett., **82**, 691 (1999)

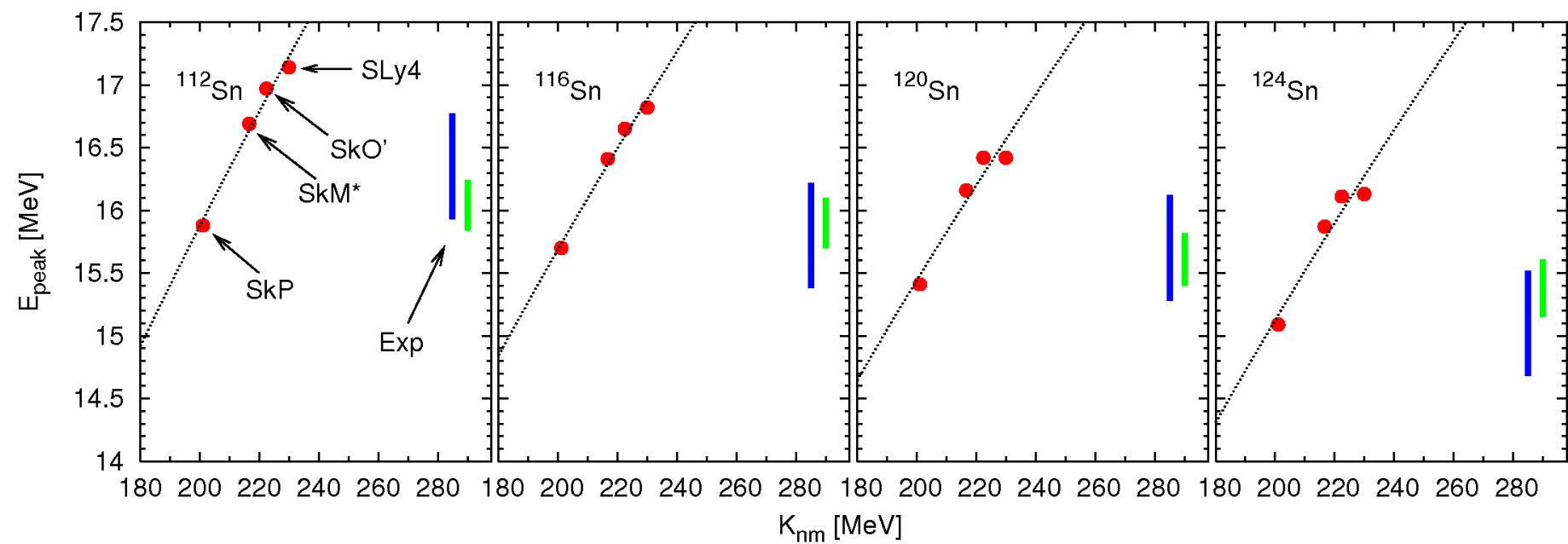
TABLE I. GMR energies and errors in MeV.

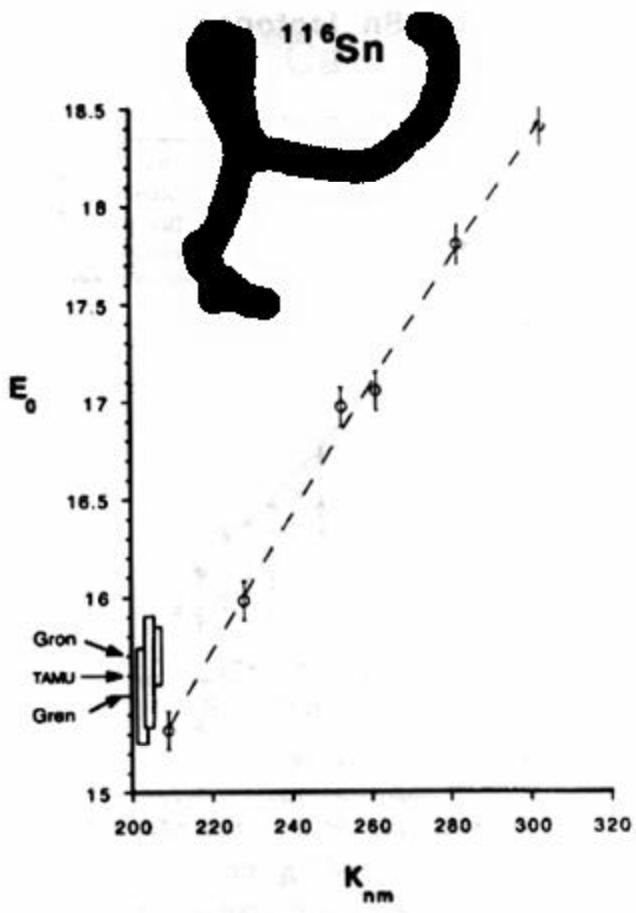
	TAMU 1998 Gaussian Cross Section		Previous Work Gaussian Cross Section		TAMU 1998 $E0$ Strength $m_1/m_0$ MeV		TAMU 1998 Slice Analysis $E0$ Strength $m_1/m_0$ MeV	
	Centroid MeV	error MeV	Centroid MeV	error MeV			Centroid MeV	error MeV
<sup>90</sup> Zr	16.44	0.07	16.10	0.28	16.80		17.89	0.20
<sup>116</sup> Sn	15.77	0.07	15.60	0.16	16.00		16.07	0.12
<sup>144</sup> Sm	15.16	0.11	15.10	0.14	15.31		15.39	0.28
<sup>208</sup> Pb	13.91	0.11	13.90	0.30	14.24		14.17	0.28

15.50 0.20 in 1993

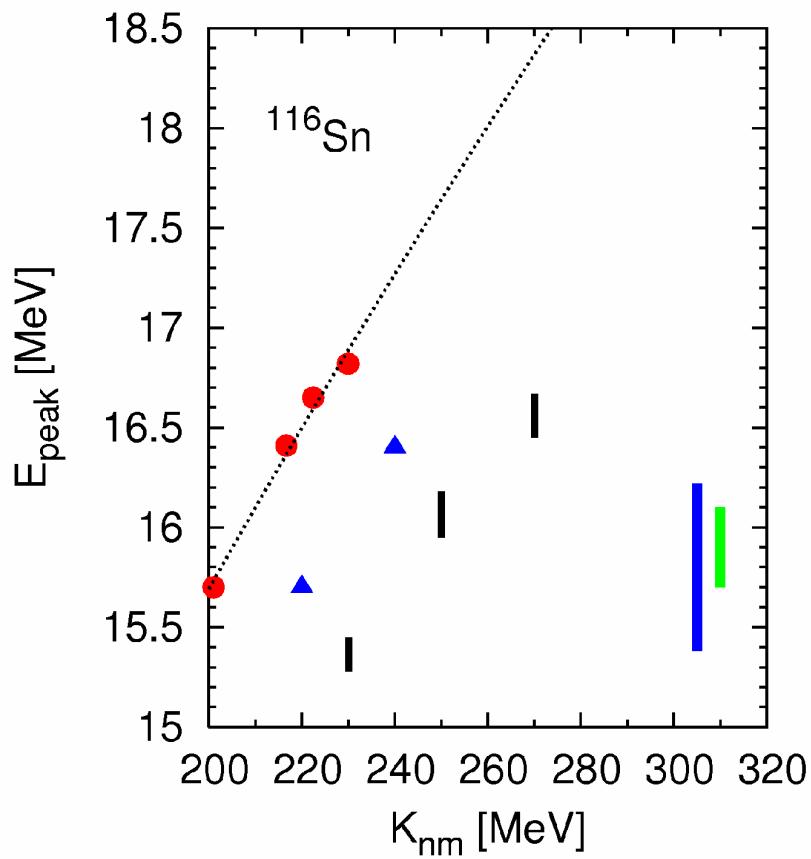
TAMU 1998 Adopted Energies $E0$ Strength $(m_1/m_{-1})^{1/2}$ MeV		error MeV
	17.81	0.35
	15.90	0.07
	15.25	0.11
	14.18	0.11

## Peak energy of GMR and $K_{nm}$





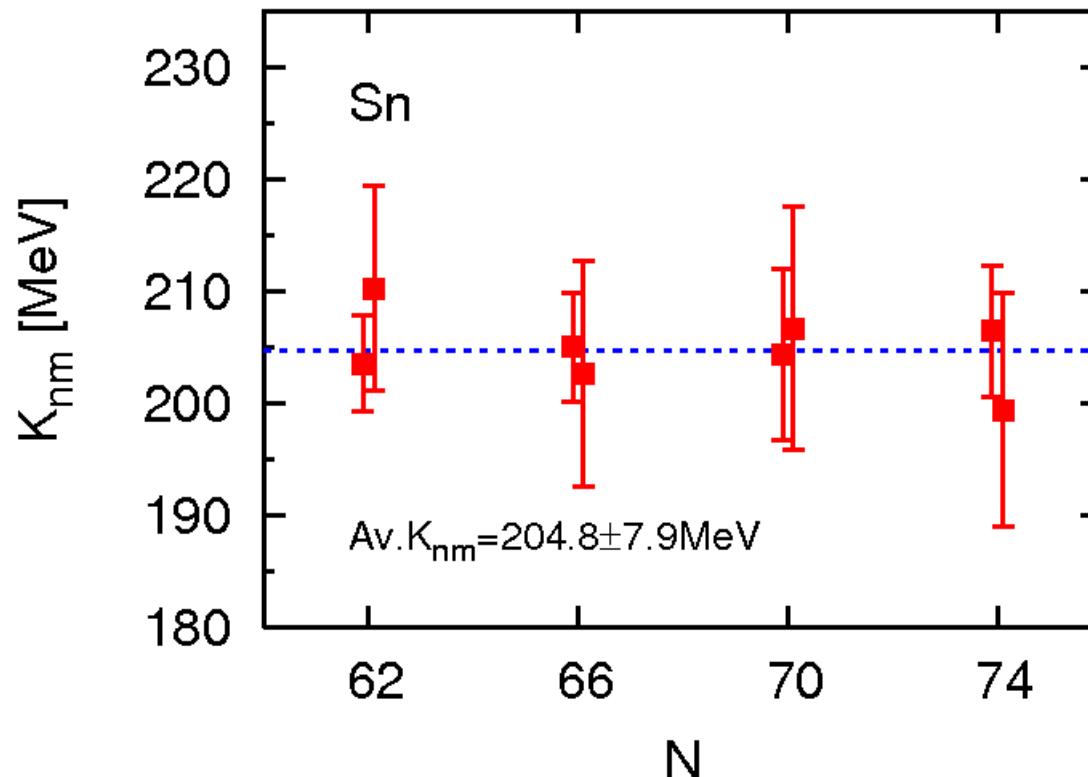
J.P. Blaizot et al., Nucl.Phys.A **591**,  
435 (1995)



▲ M.Farine et al., Nucl.Phys.A **615**,  
135 (1997)

█ D.Vretenar et al., P.R.C **68**,  
024310 (2003)

## $K_{nm}$ obtained



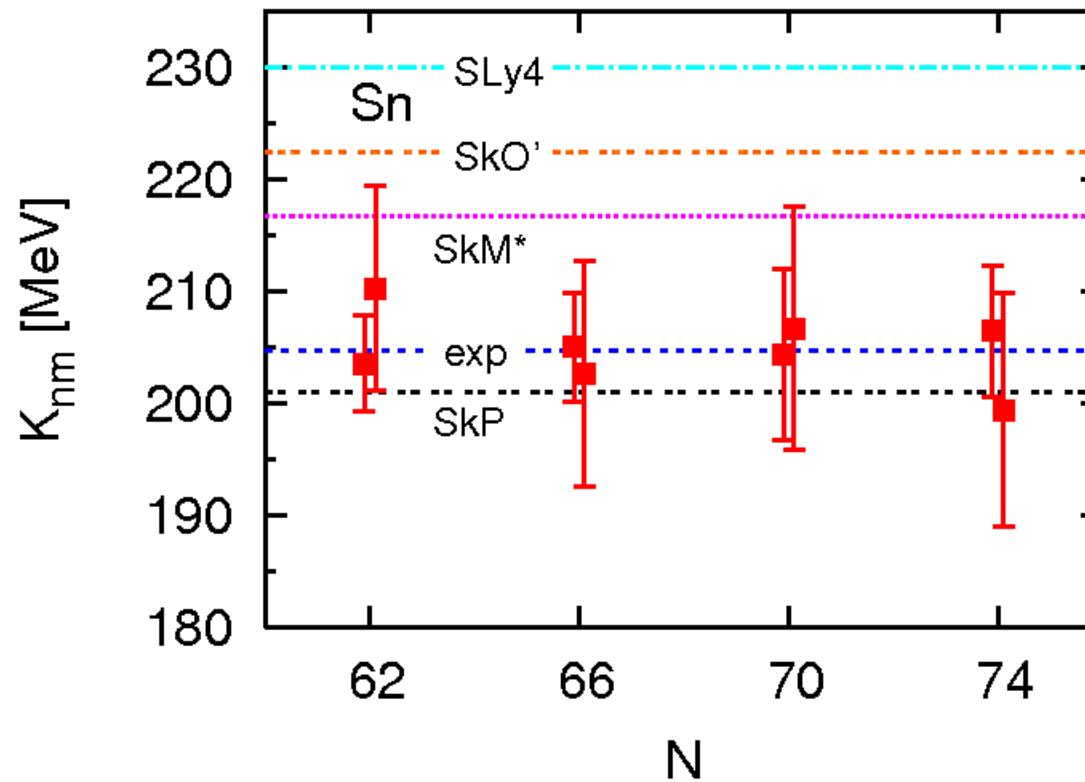
Blaizot et al.,  $\langle K_{nm} \rangle$  for  $^{116}\text{Sn}$

Farine et al.,  $K_V = 215 \pm 15 \text{ MeV}$

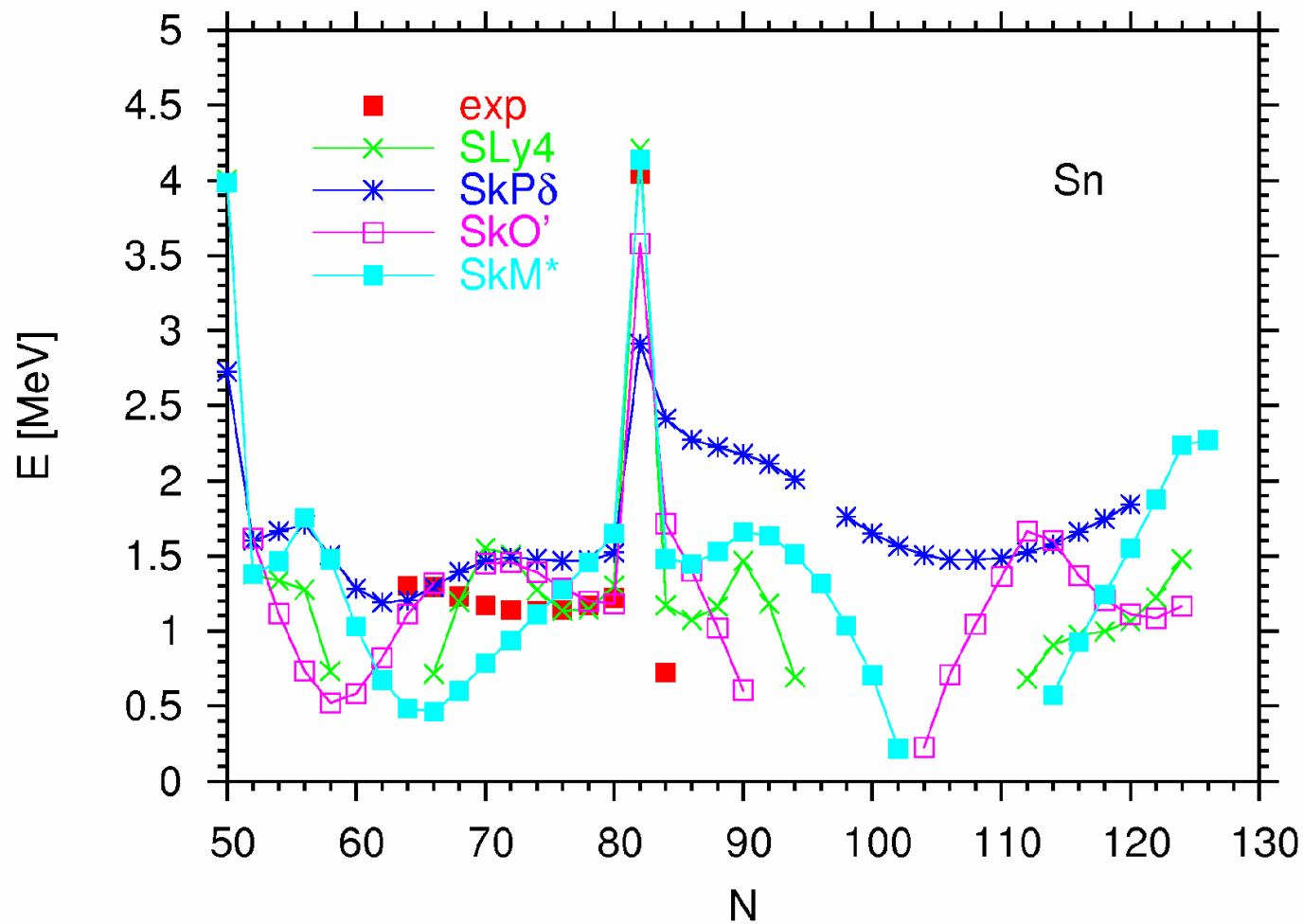
Colo` et al., "Bone fide Skyrme forces can either predict 230-240 MeV ..."

Vretenar et al.,  $K_{nm} \gg 250 - 270 \text{ MeV}$

$K_{nm}$  obtained

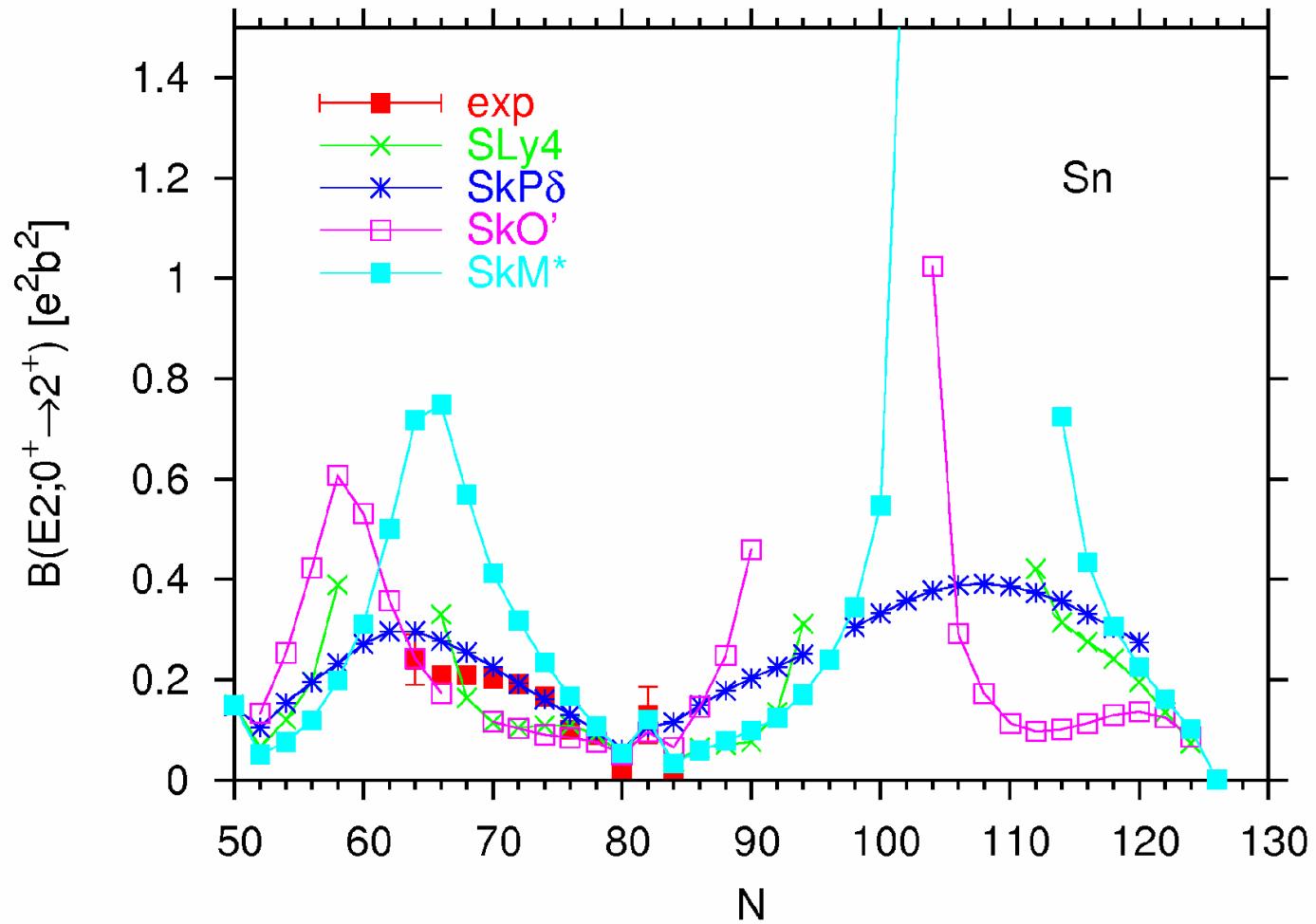


## Energies of the lowest $2^+$ states of Sn

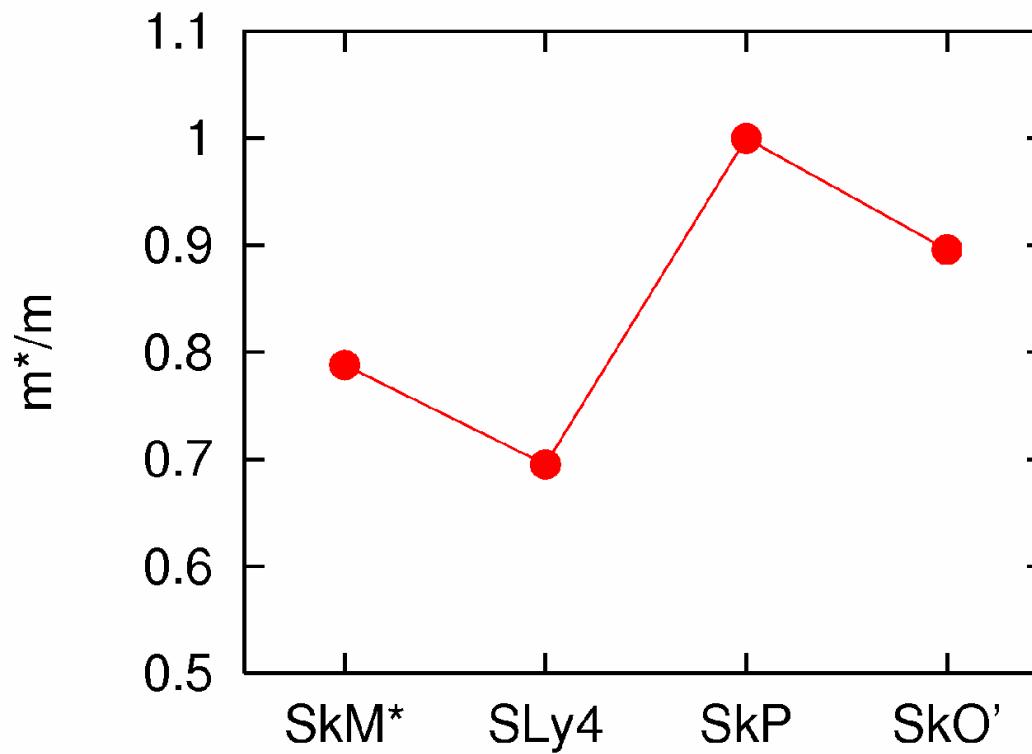


Exp.: S.Raman et al., Atom.Dat.Nucl.Dat.Tab. **78**, 1 (2001),  
D.C.Radford et al., P.R.L. **88**, 222501 (2002); talk in ENAM04

## Transition probabilities of the lowest $2^+$ states of Sn



## Effective mass (nuclear matter) / bare mass



## Summary

- Systematic QRPA calculations have been done for even Sn, Ni, and Ca with  $J^\pi=0^+, 1^-,$  and  $2^+$  from the proton drip line to the neutron drip line with a few parameter sets of the Skyrme interaction.
- Compression modulus was deduced from the experimental data of  $^{112,116,120,124}\text{Sn}$  using the QRPA calculation;

$$K_{nm} = 204.8 \pm 7.9 \text{ MeV.}$$

SkP seems good.

I would like to thank

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