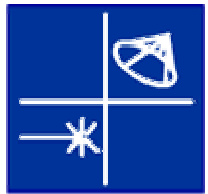


Introduction to Nuclear Science

PIXE-PAN Summer Science Program

University of Notre Dame

2007



J I N A

Tony Hyder, Professor of Physics

The Joint Institute for Nuclear Astrophysics

www.JINAweb.org

Topics we will discuss...

Ground-state properties of the nucleus

size, shape, stability, binding energies, angular momenta

Radioactivity

alpha, beta, and gamma decay

The nuclear force

Nuclear reactions

the compound nucleus, Q values, excited states

Decay modes of an excited nucleus

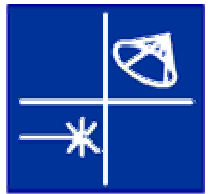
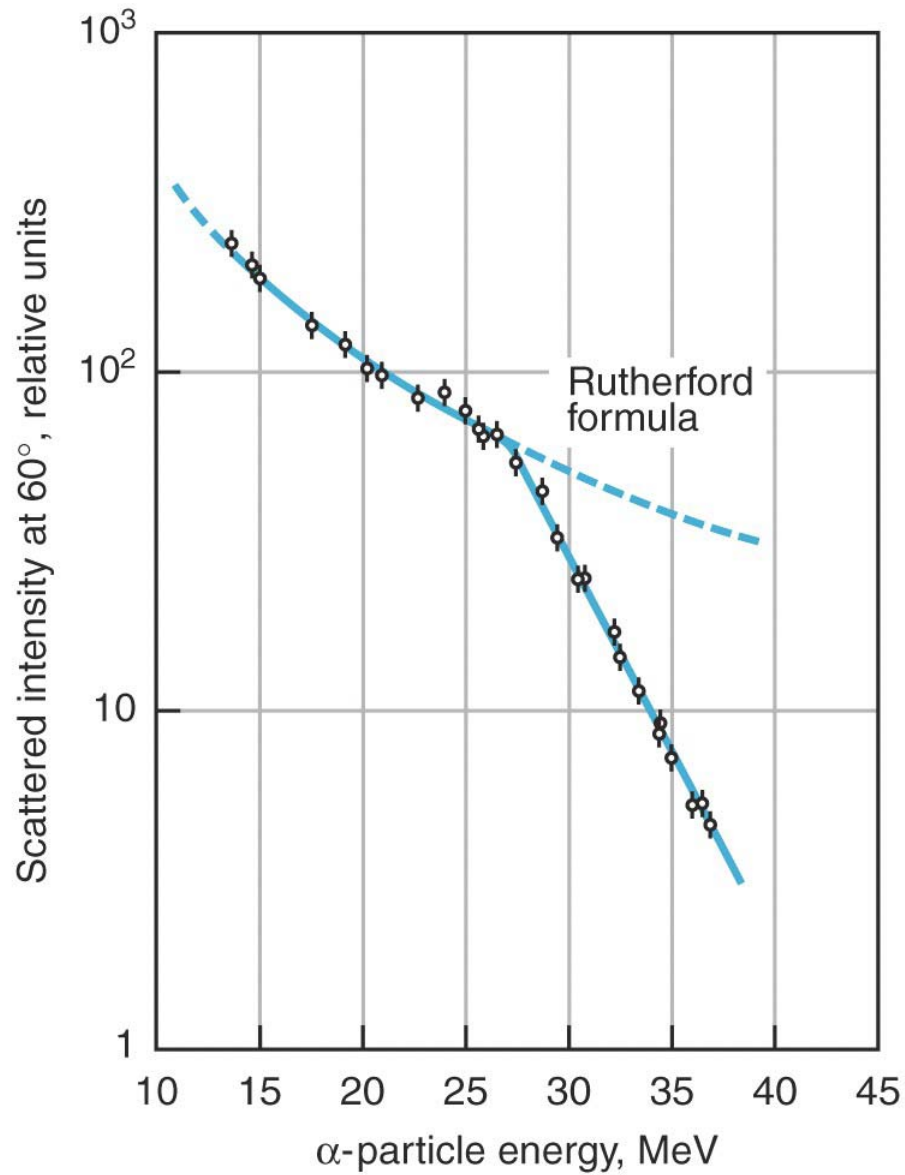
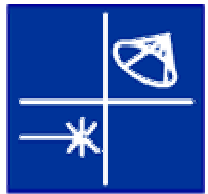


TABLE 11-1 Fundamental properties of atomic constituents

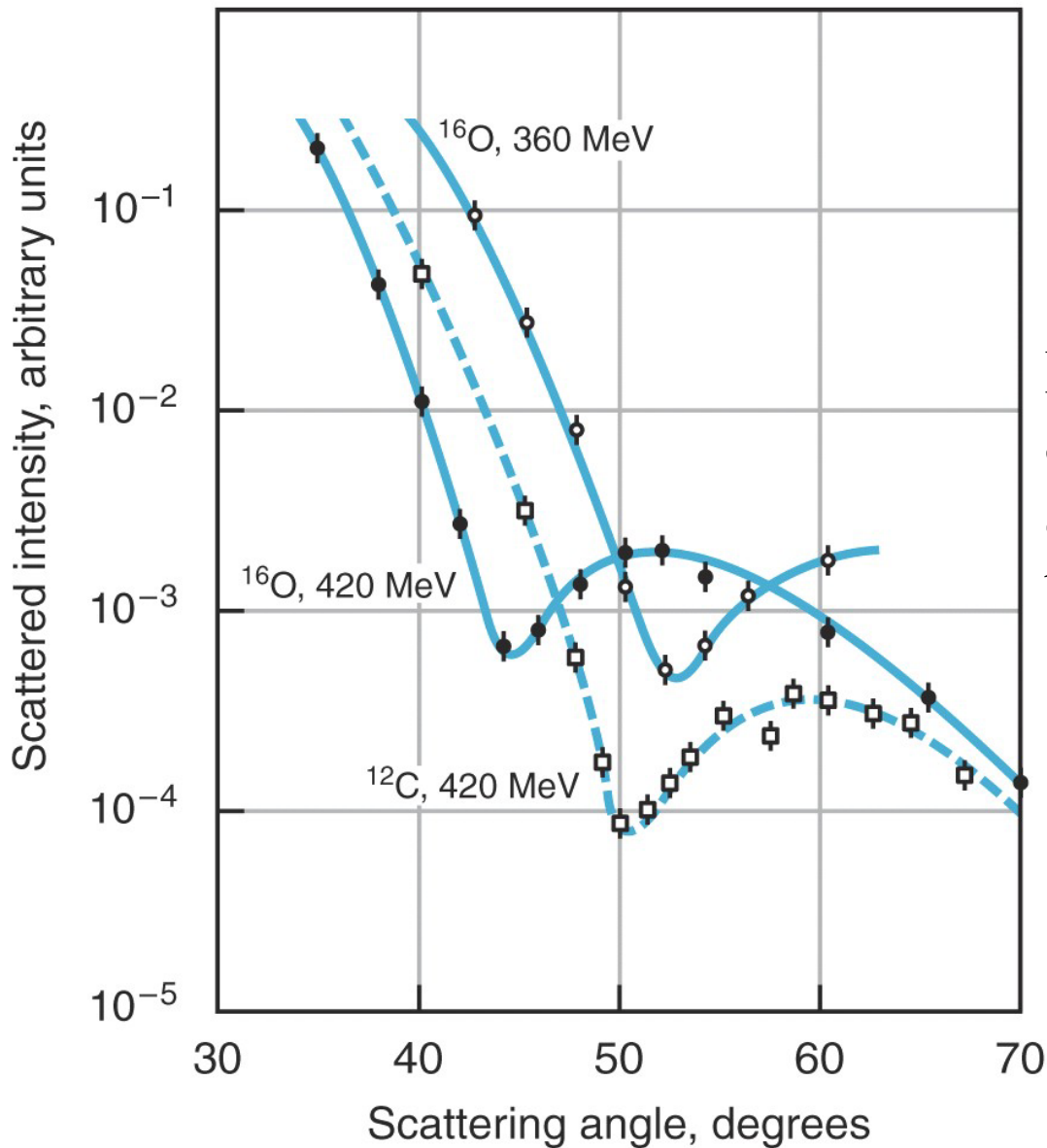
Particle	Charge	Mass (u)	Mass (kg)	Spin	Magnetic moment
Proton	$+e$	1.007276	1.6726×10^{-27}	1/2	$2.79285 \mu_N$
Neutron	0	1.008665	1.6749×10^{-27}	1/2	$-1.91304 \mu_N$
Deuteron	$+e$	2.013553	3.3436×10^{-27}	1	$0.85744 \mu_N$
Electron	$-e$	5.4858×10^{-4}	9.1094×10^{-31}	1/2	$1.00116 \mu_B$



The radius
of the
nucleus...

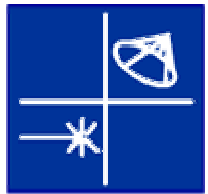


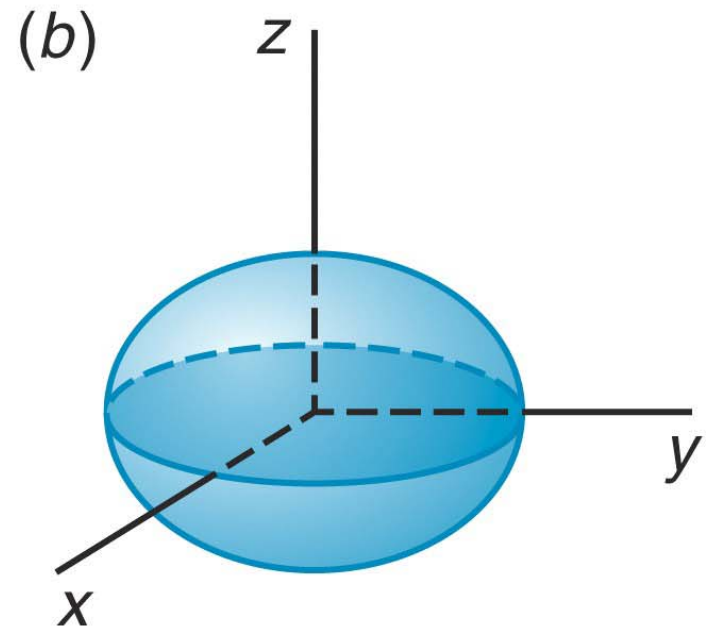
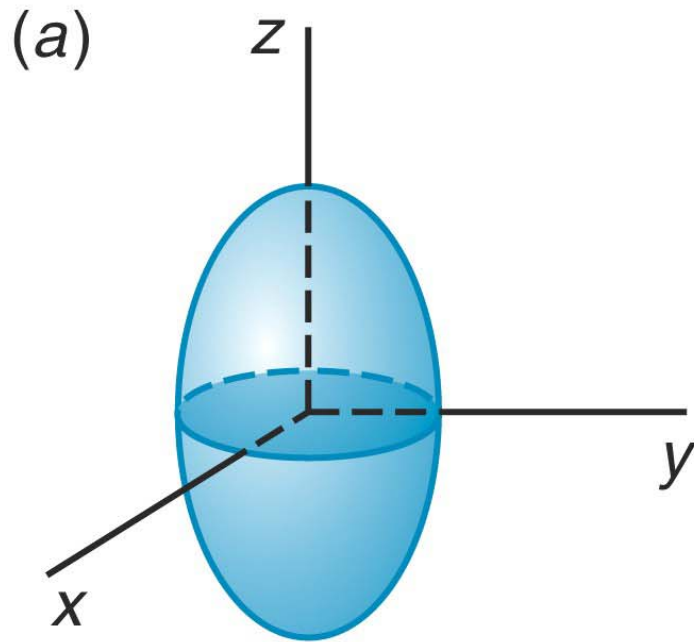
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Diffraction pattern of high-energy electrons scattered by ^{16}O and ^{12}C

$$\sin\theta = 0.61 \lambda/R$$



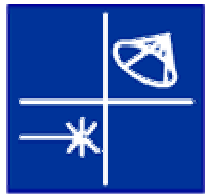


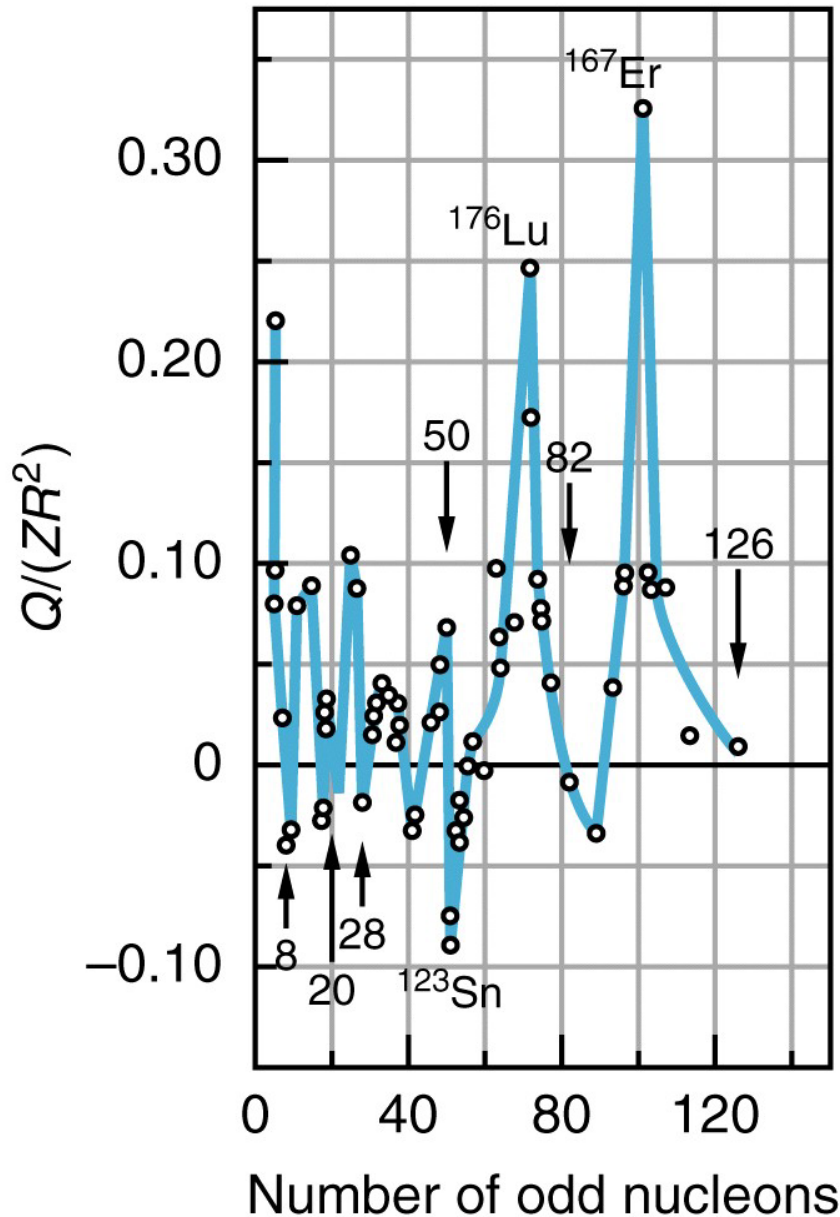
Non-spherical nuclear shapes. The electric quadrupole moment is given by $3z^2 - x^2 - y^2 - z^2$

> 0 for a football on end

$= 0$ for a sphere

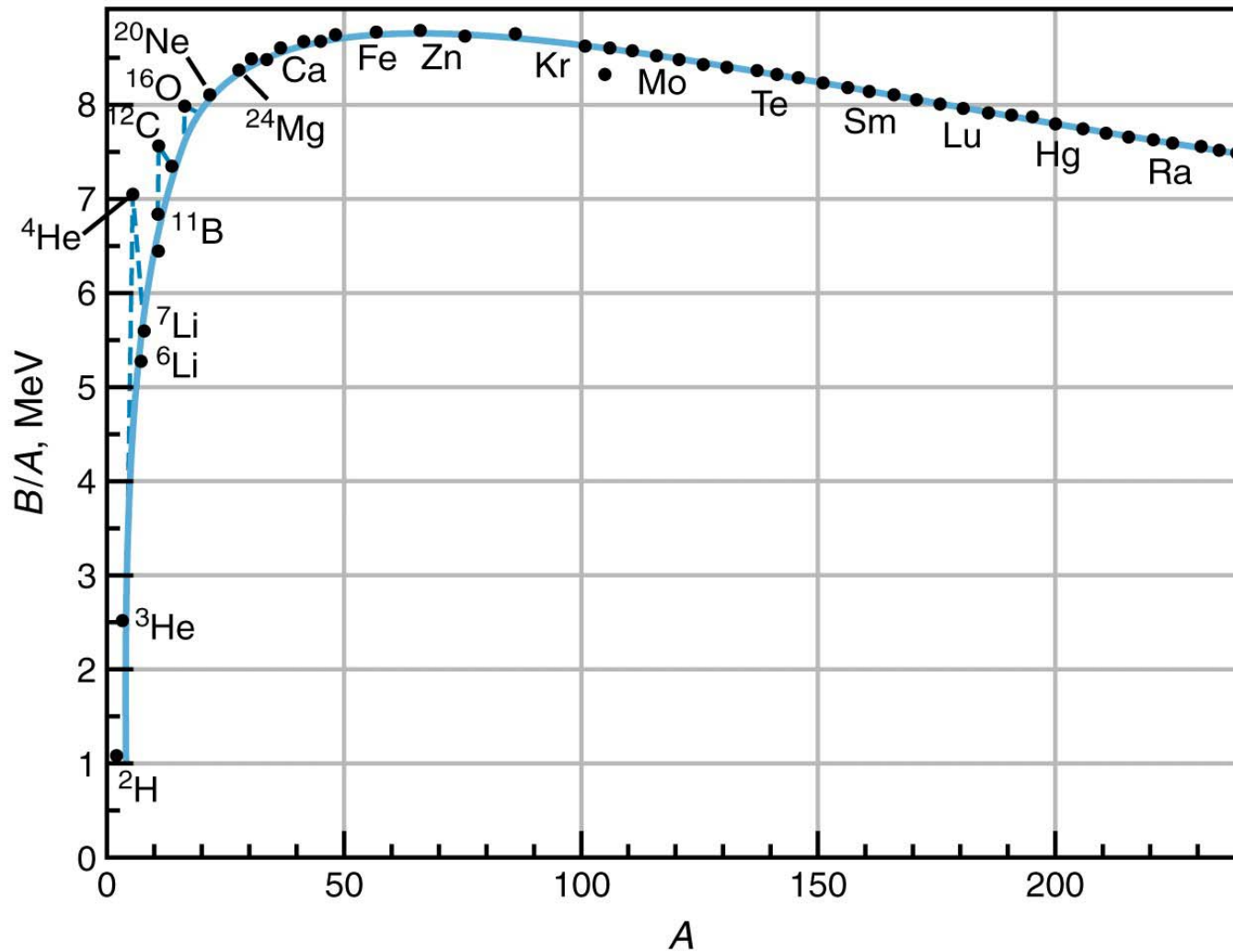
< 0 for an egg on the table



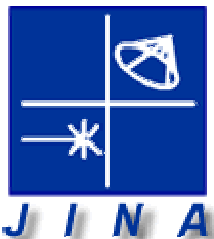


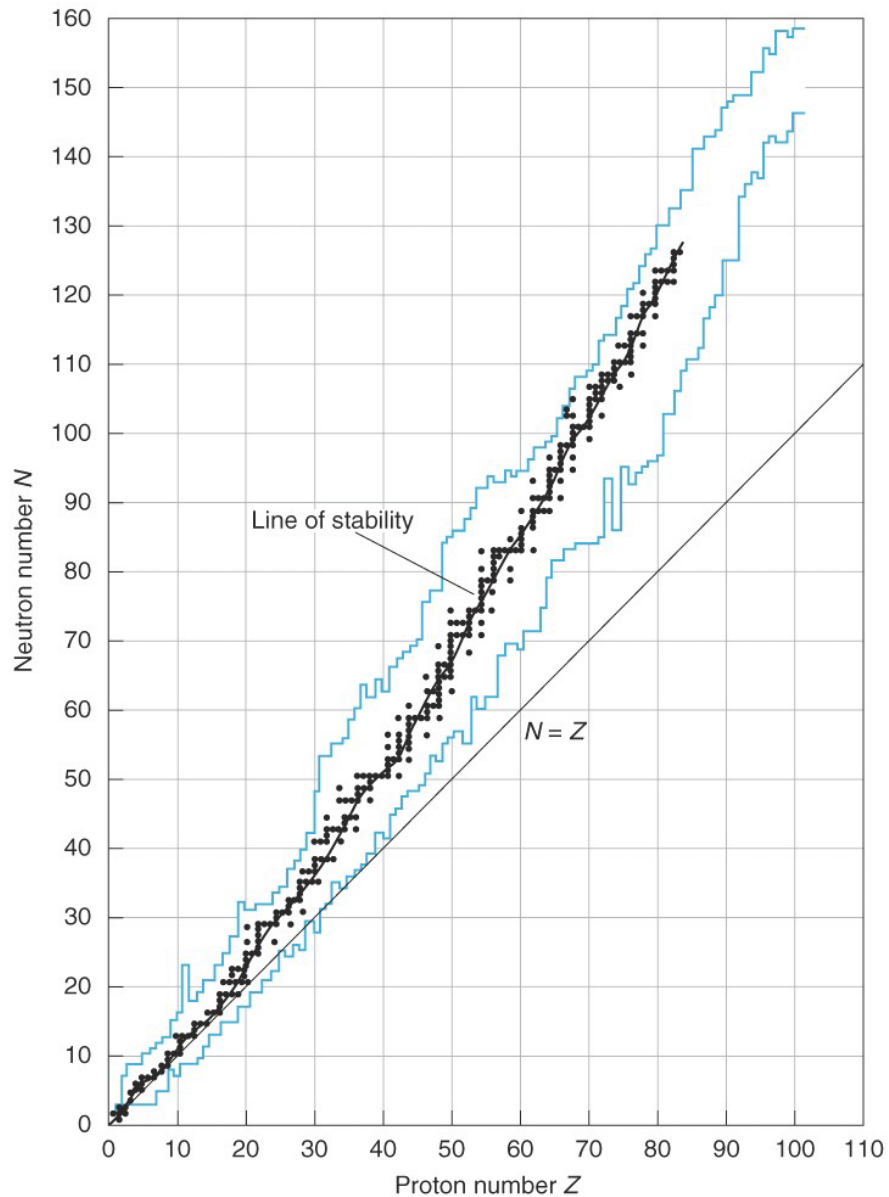
Quadrupole moments of a number of odd-A nuclei. The arrows point to spherical nuclei.

The numbers are the shell-model 'magic numbers' about which we will talk later.

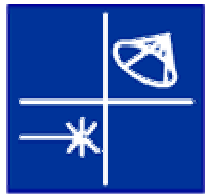


The curve of binding energy: binding energy per nucleon. Note that above mass 40 or so, it is constant

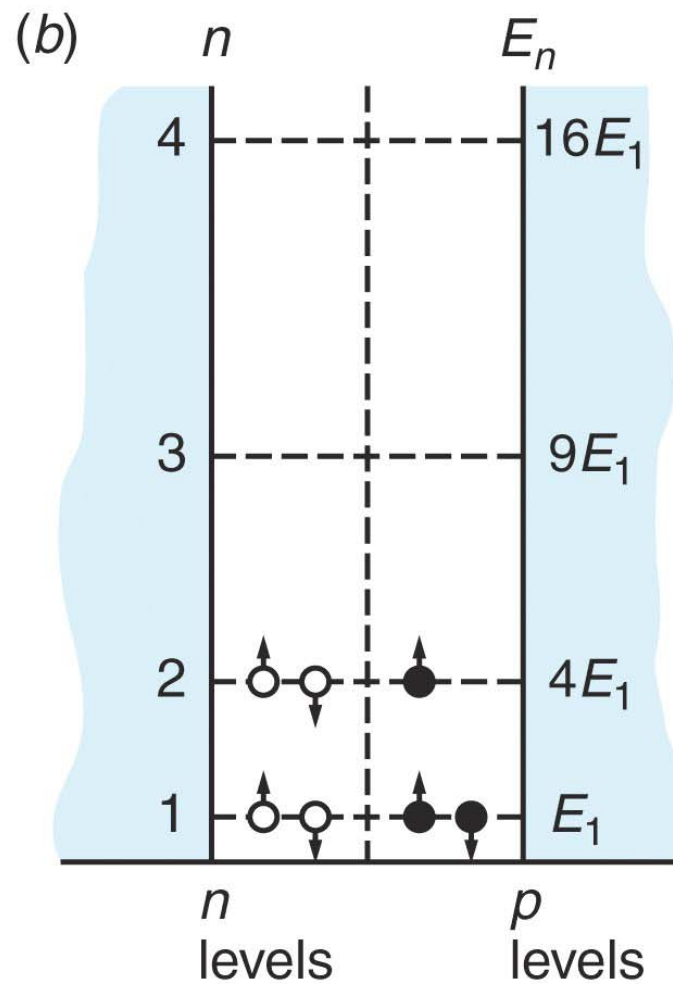
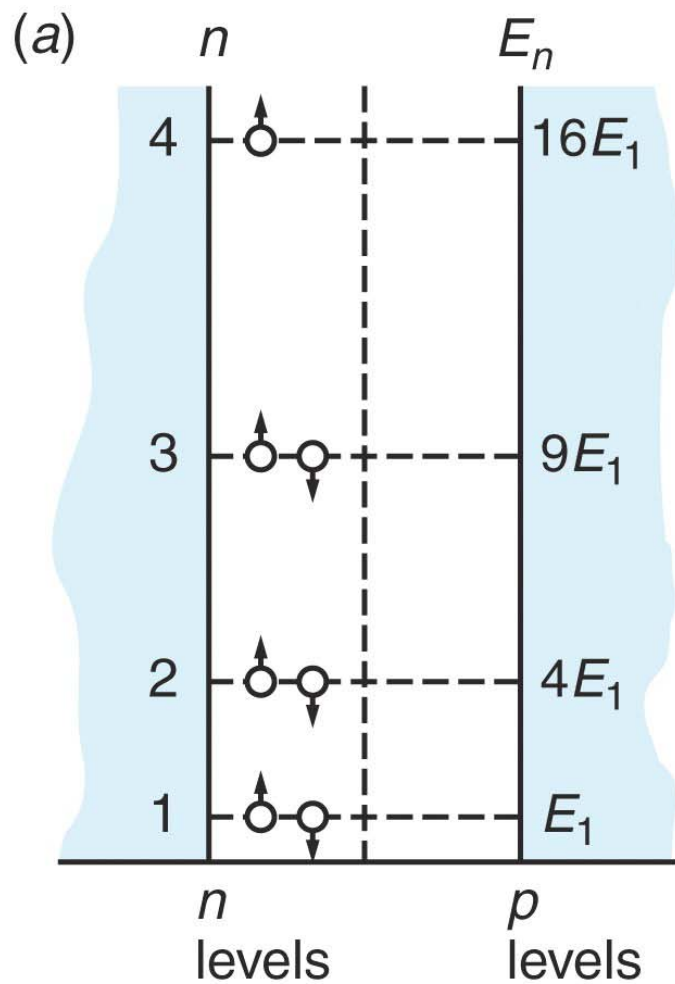




Of the 3000 or so known nuclides, there are only 266 whose ground states are stable. The rest are radioactive



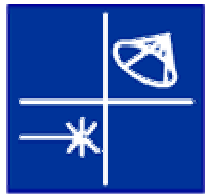
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Any idea what this is about?

TABLE 11-2 N versus Z for stable isotopes

		Z	
N		Even	Odd
Even		159	50
Odd		53	4

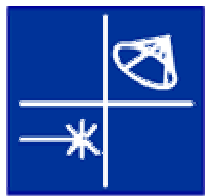


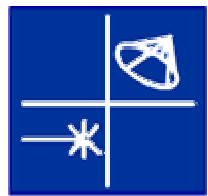
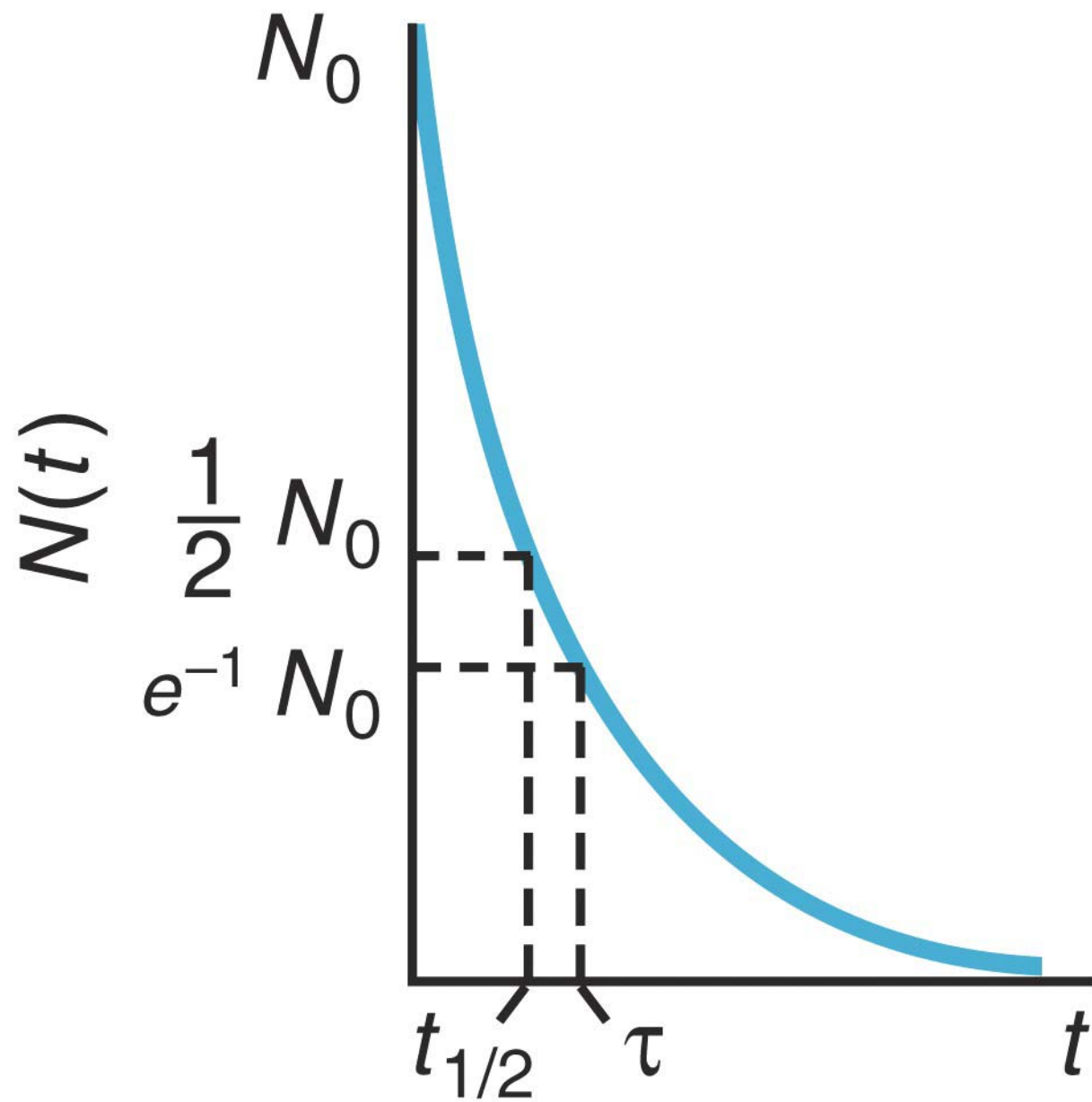
Radioactivity

For a nucleus to be radioactive at all, its mass must be greater than the sum of the masses of the decay products.

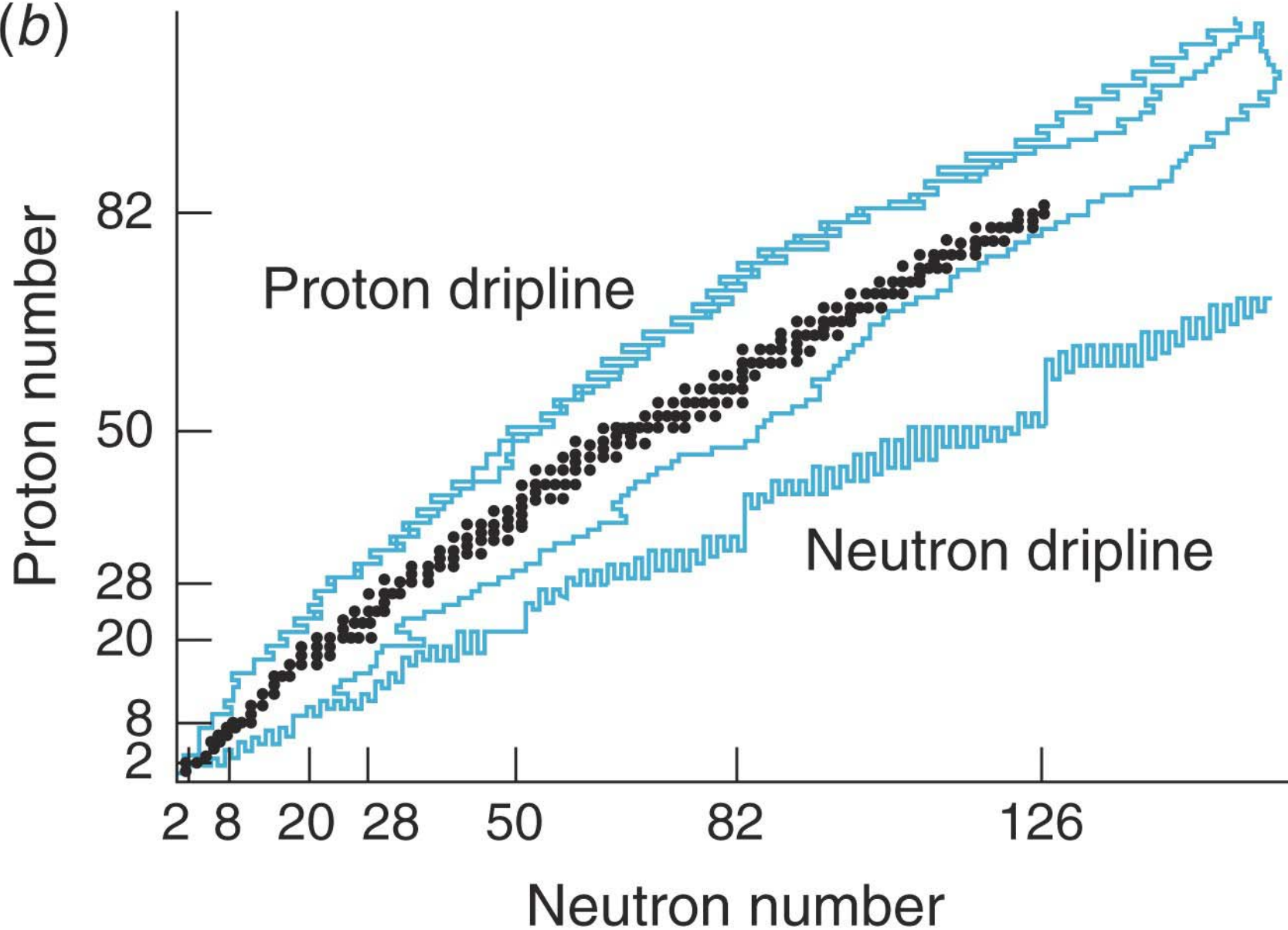
We will look briefly at three types of decay: alpha, beta, and gamma

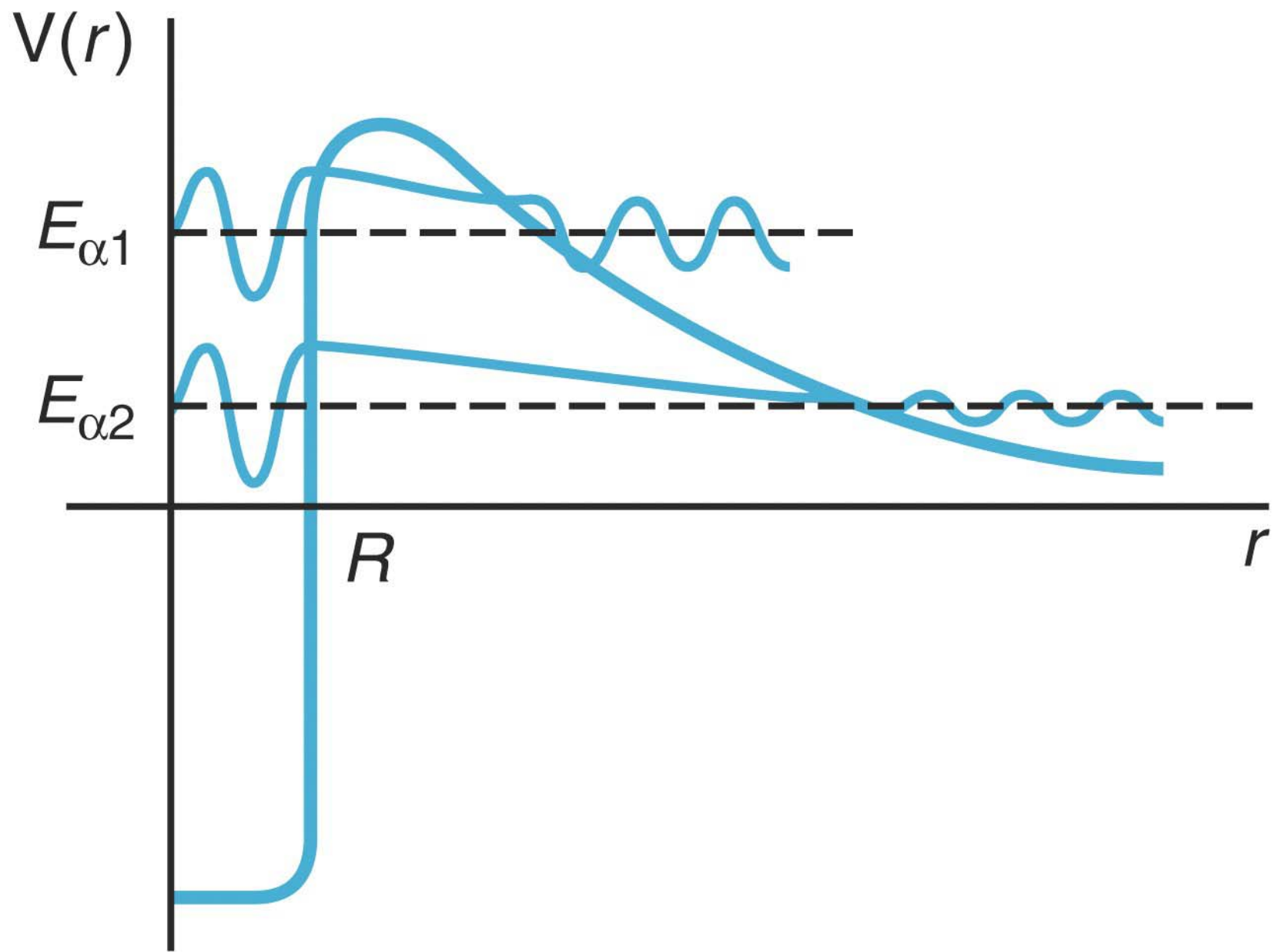
Many of the heavy nuclei are unstable to alpha decay, and because the Coulomb barrier inhibits the decay process, the half life for alpha decay can be very long if the decay energy is small. All very heavy nuclei ($Z > 83$) are theoretically unstable to α decay since the mass of the parent is greater than the sum of the masses of the decay products

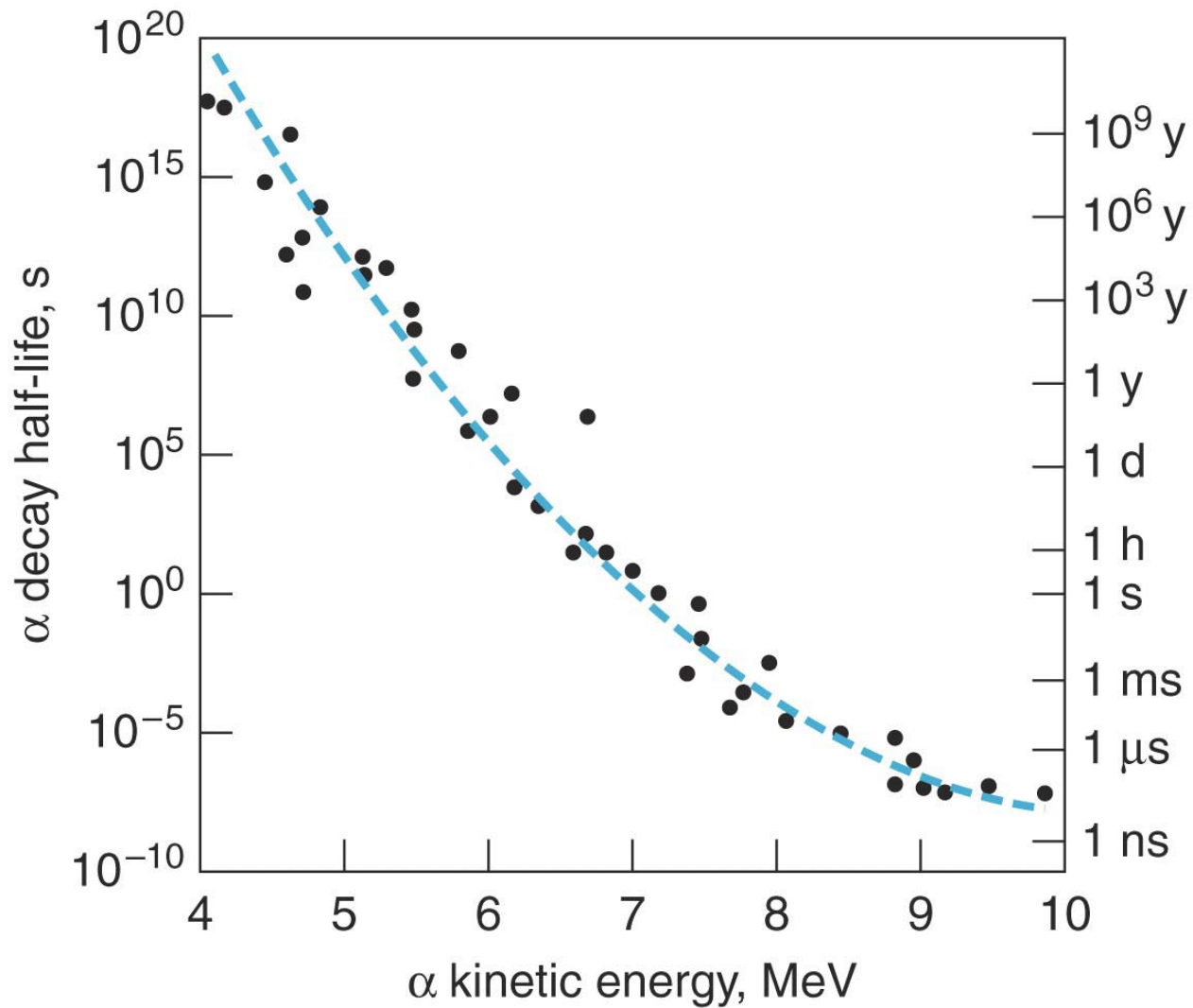




(b)



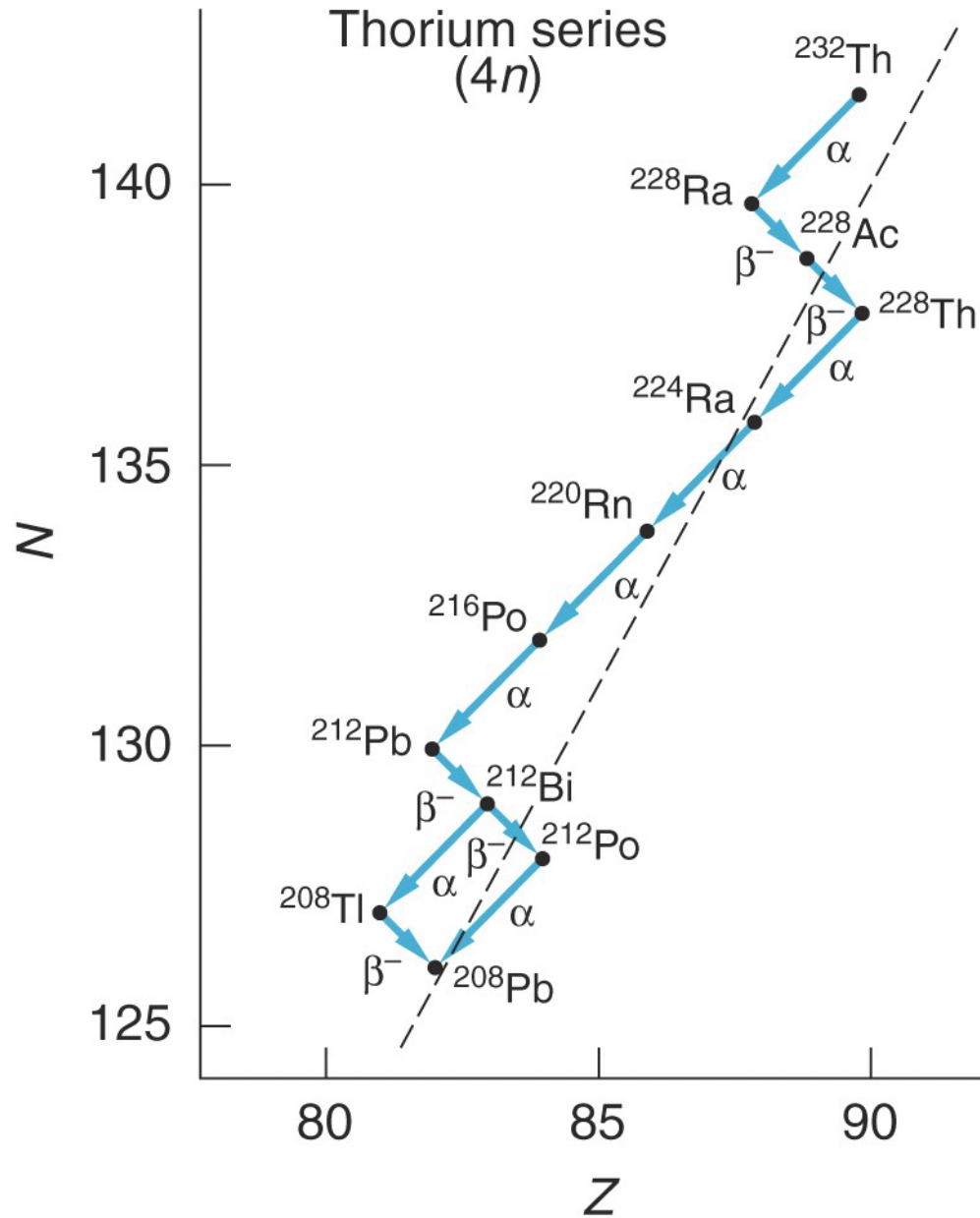




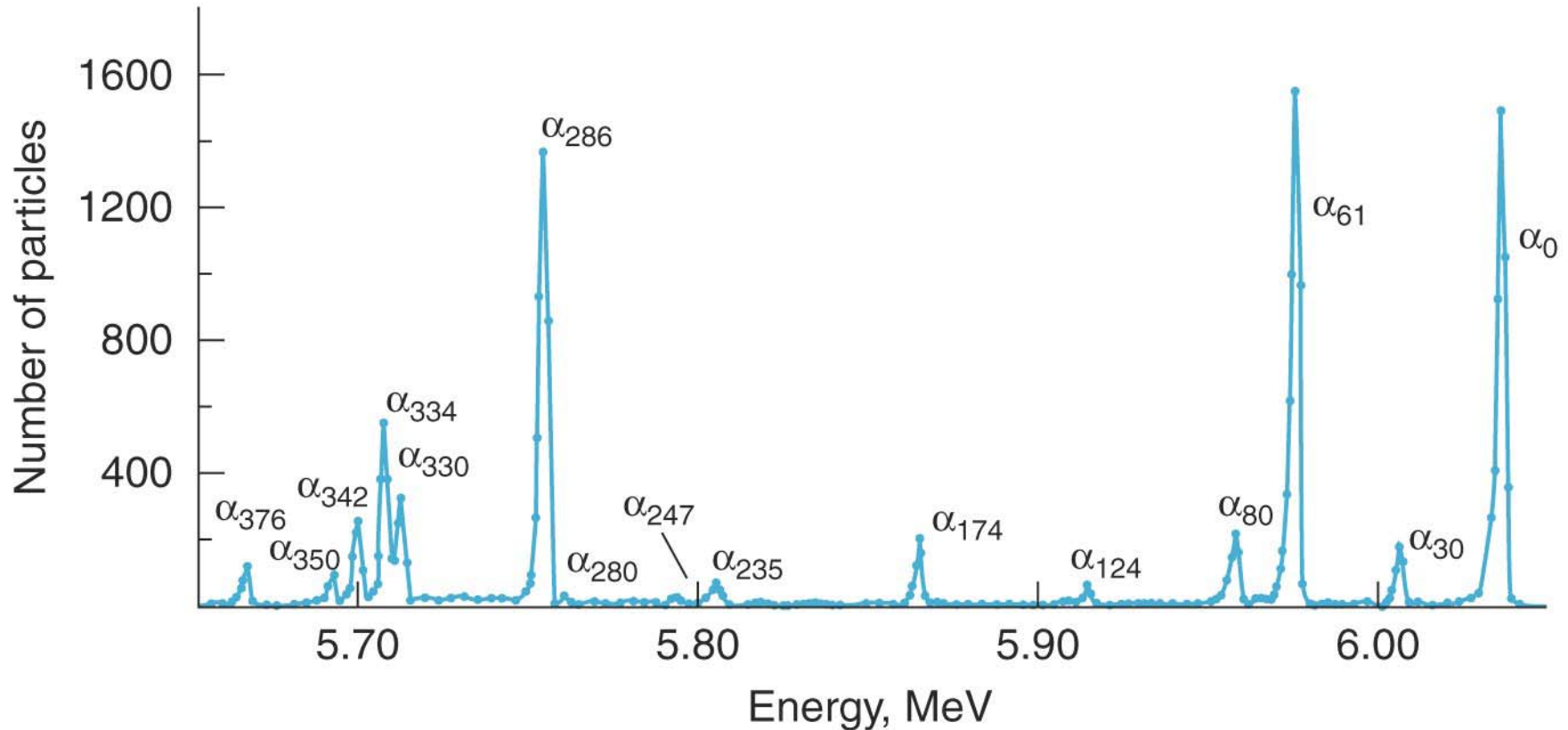
The Geiger-Nuttall relation

Alpha decay

The broken line is the line of stability, i.e., the floor of the energy valley we saw earlier

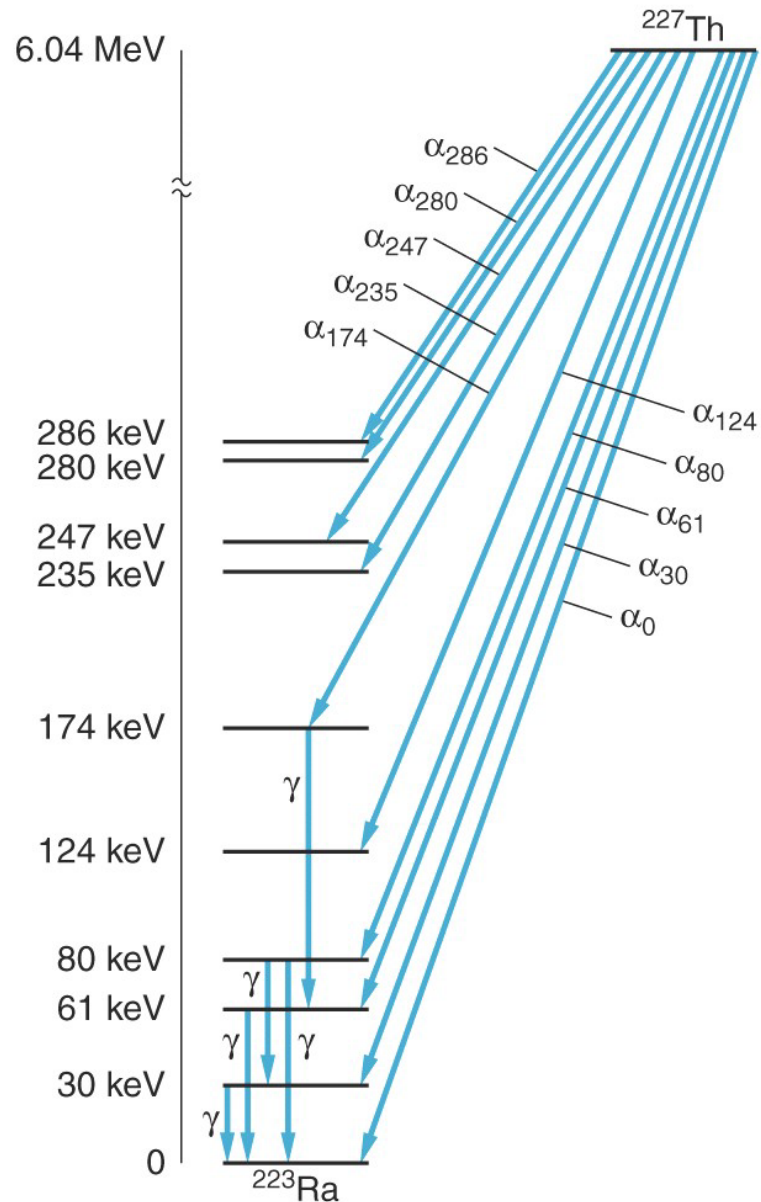


The alpha-particle spectrum from ^{227}Th ... the highest energy alpha particles corresponds to decay to the ground state of ^{223}Ra

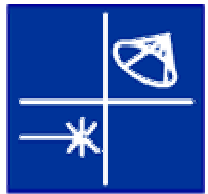
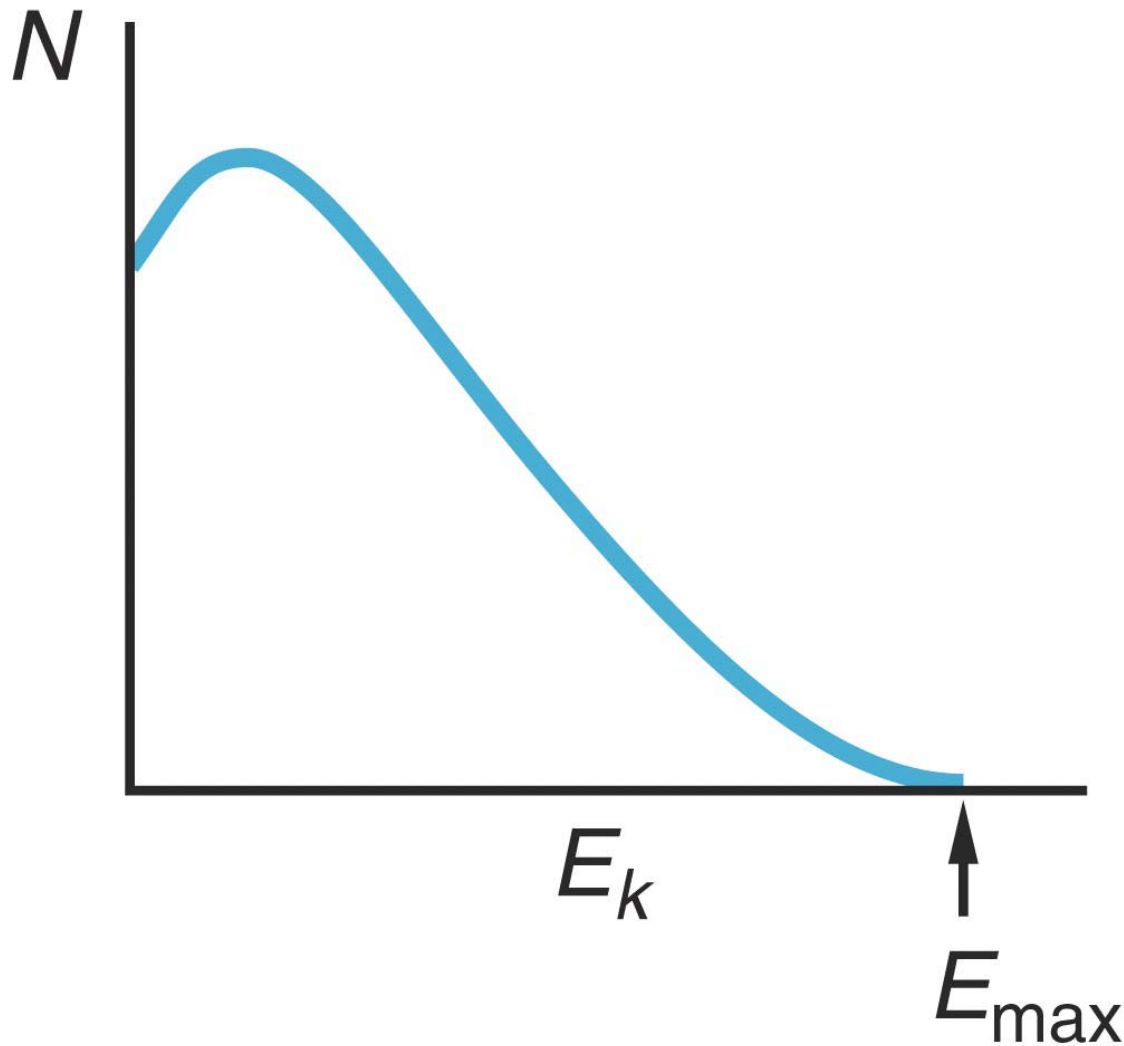


The energy levels of ^{223}Ra can be determined from the measurement of the alpha-particle energies we saw in the pervious slide.

Not all of the gamma-ray transitions are shown

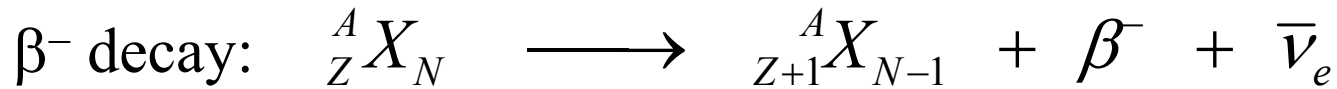


The energy spectrum of electrons emitted in beta decay



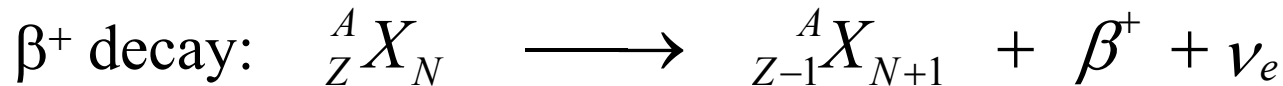
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Beta decay



A neutron changes into a proton and emits an electron

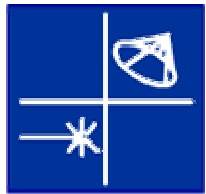
$$Q = (M_P - M_D)c^2$$



A proton changes into a neutron and emits a positron

$$Q = (M_P - M_D + 2m_e)c^2$$

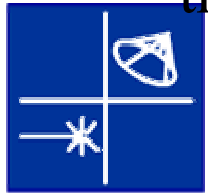
Electron capture: a process that competes with β^+ decay in which a proton in the nucleus captures an atomic electron and changes into a neutron with the emission of a neutrino



Gamma decay

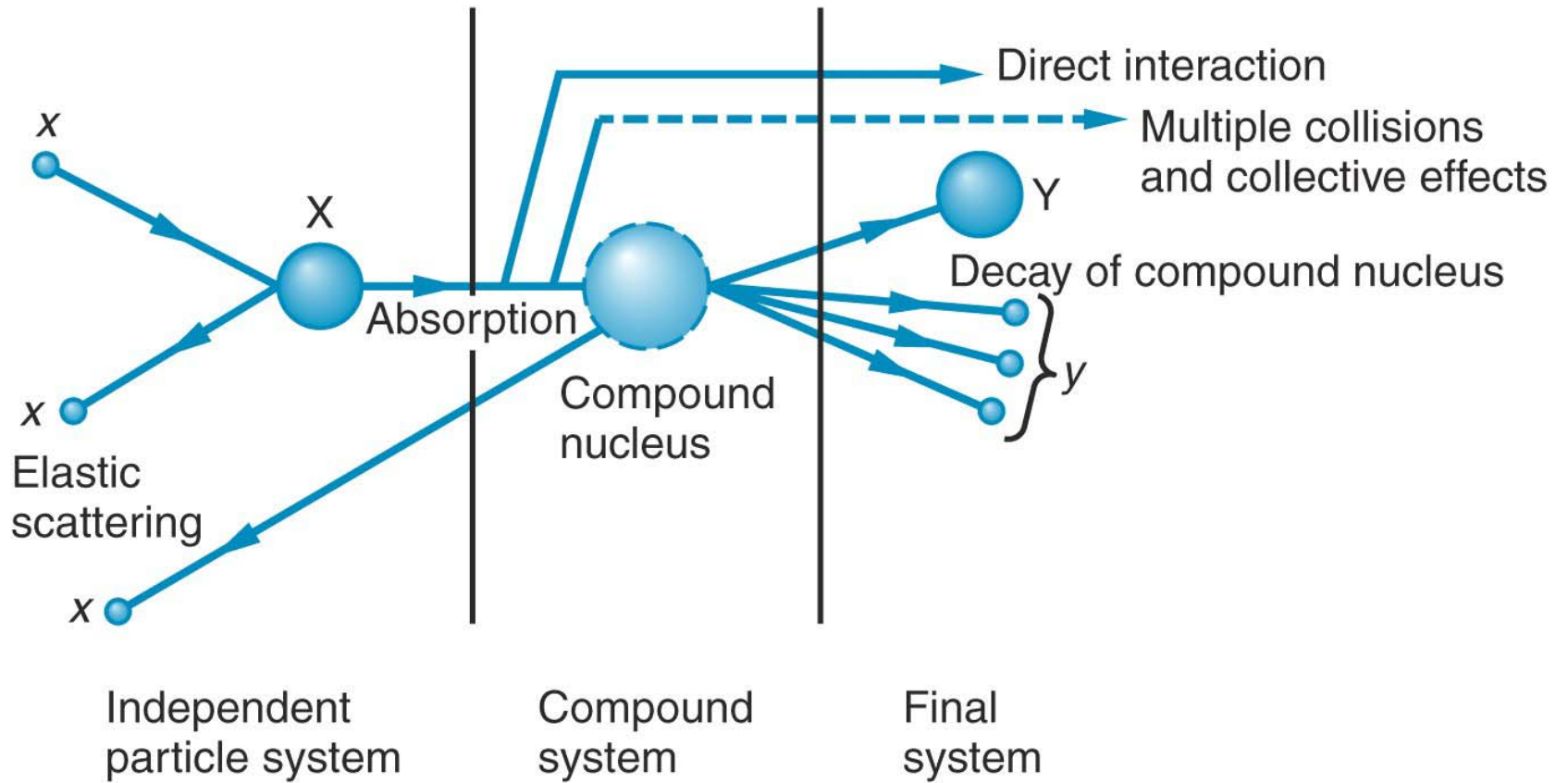
A process in which a nucleus in an excited state decays to a lower energy state of the same isotope by the emission of a photon. We saw an example of this earlier in the decay of ^{223}Ra

Internal conversion is a competing process especially for lower-lying energy states, in which the excitation energy of the state is transferred to an orbital electron which is ejected from the atom. The ejected electron is observed to have a kinetic energy equal to the nuclear transition energy minus the electron's atomic binding energy

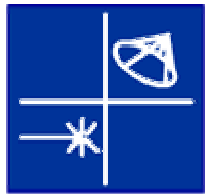


The compound nucleus.

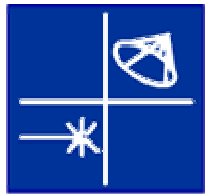
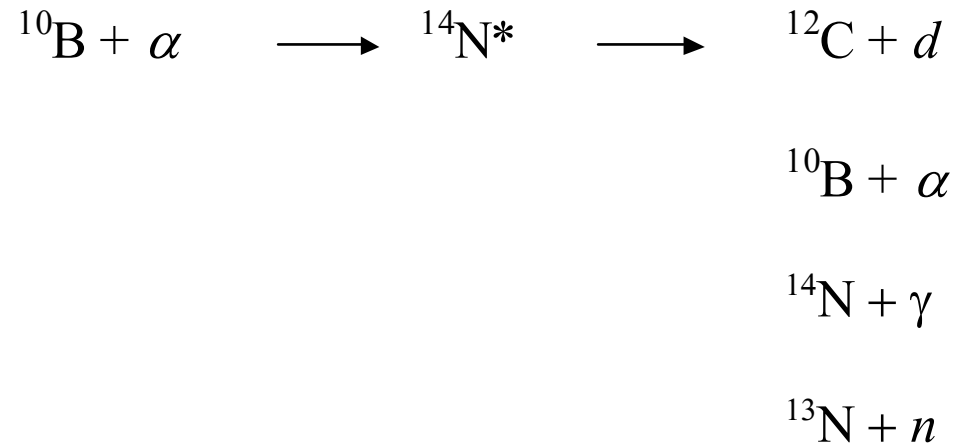
$X(x,y)Y$



$$Q = (m_x + m_X - m_y - m_Y) c^2$$

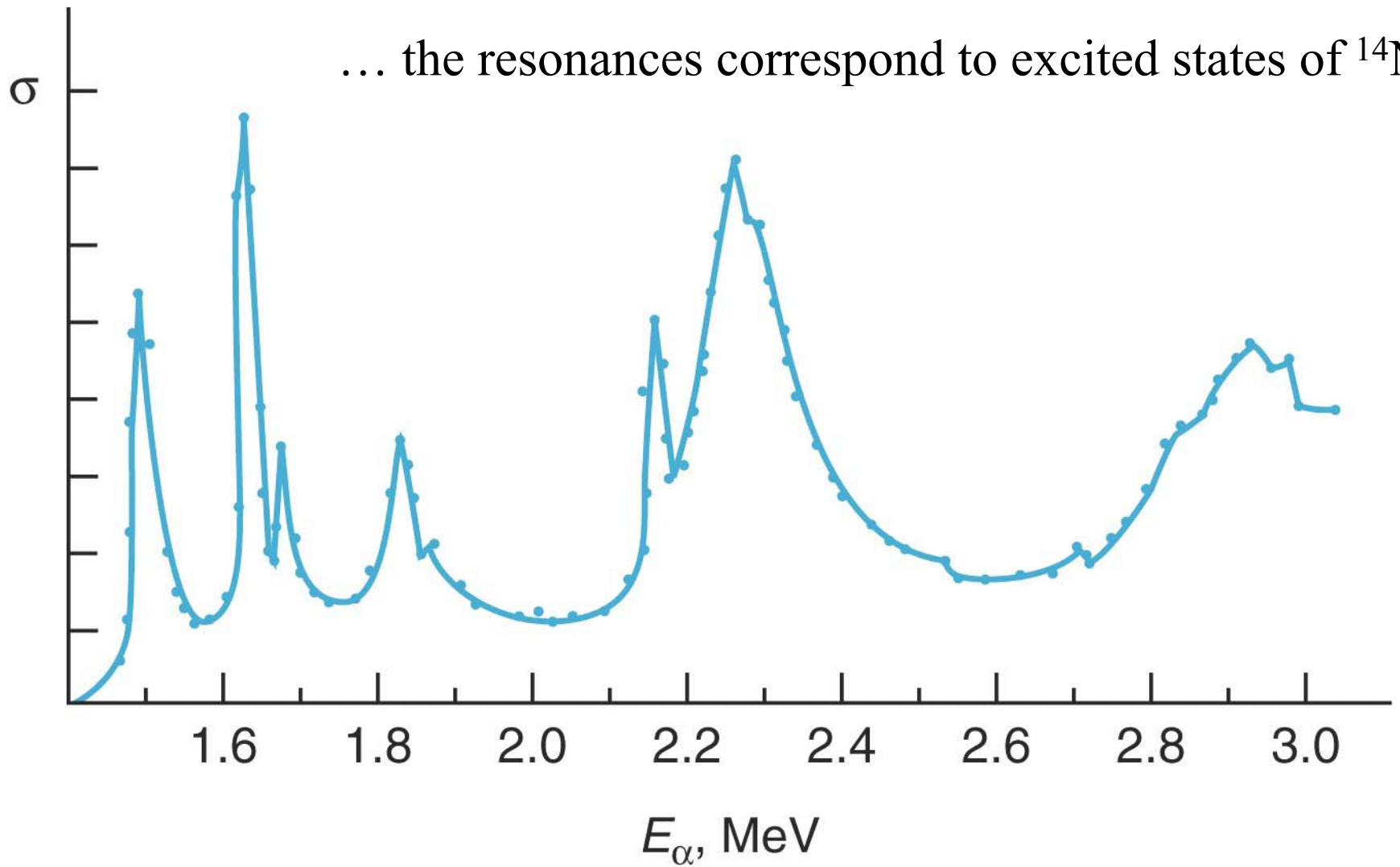


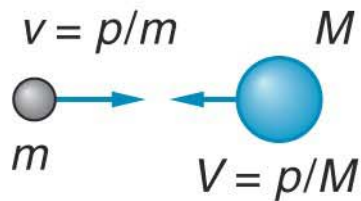
Some examples



Cross section for the $^{10}\text{Be} + \alpha$ reaction

... the resonances correspond to excited states of ^{14}N

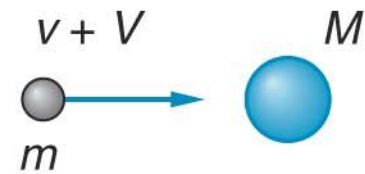




CM system

$$p = mv = MV$$

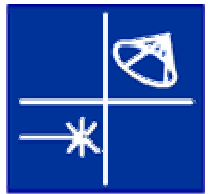
$$E_{\text{CM}} = p^2/2m + p^2/2M = (m + M)p^2/2mM$$

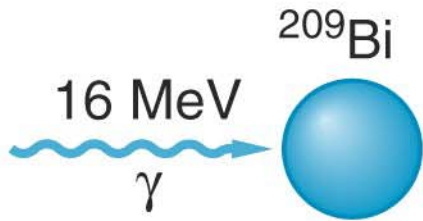


Lab system

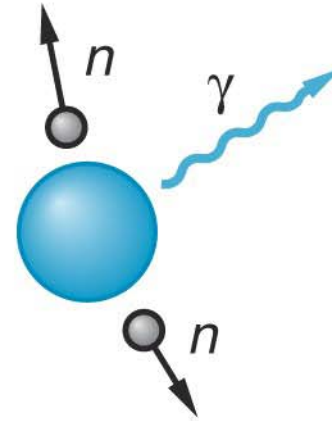
$$p_{\text{lab}} = m(v + V) = mv(1 + m/M) = \frac{M + m}{M} p$$

$$E_{\text{lab}} = \frac{p_L^2}{2m} = \left(\frac{p^2}{2m} \right) \left(\frac{M + m}{M} \right)^2 = \frac{M + m}{M} E_{\text{CM}}$$

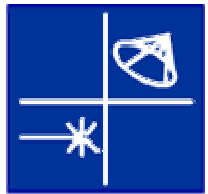




“Hot” or excited
compound
nucleus



Neutrons and photons with
energies of a few MeV boil off

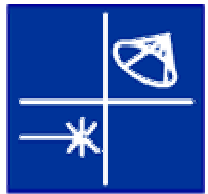
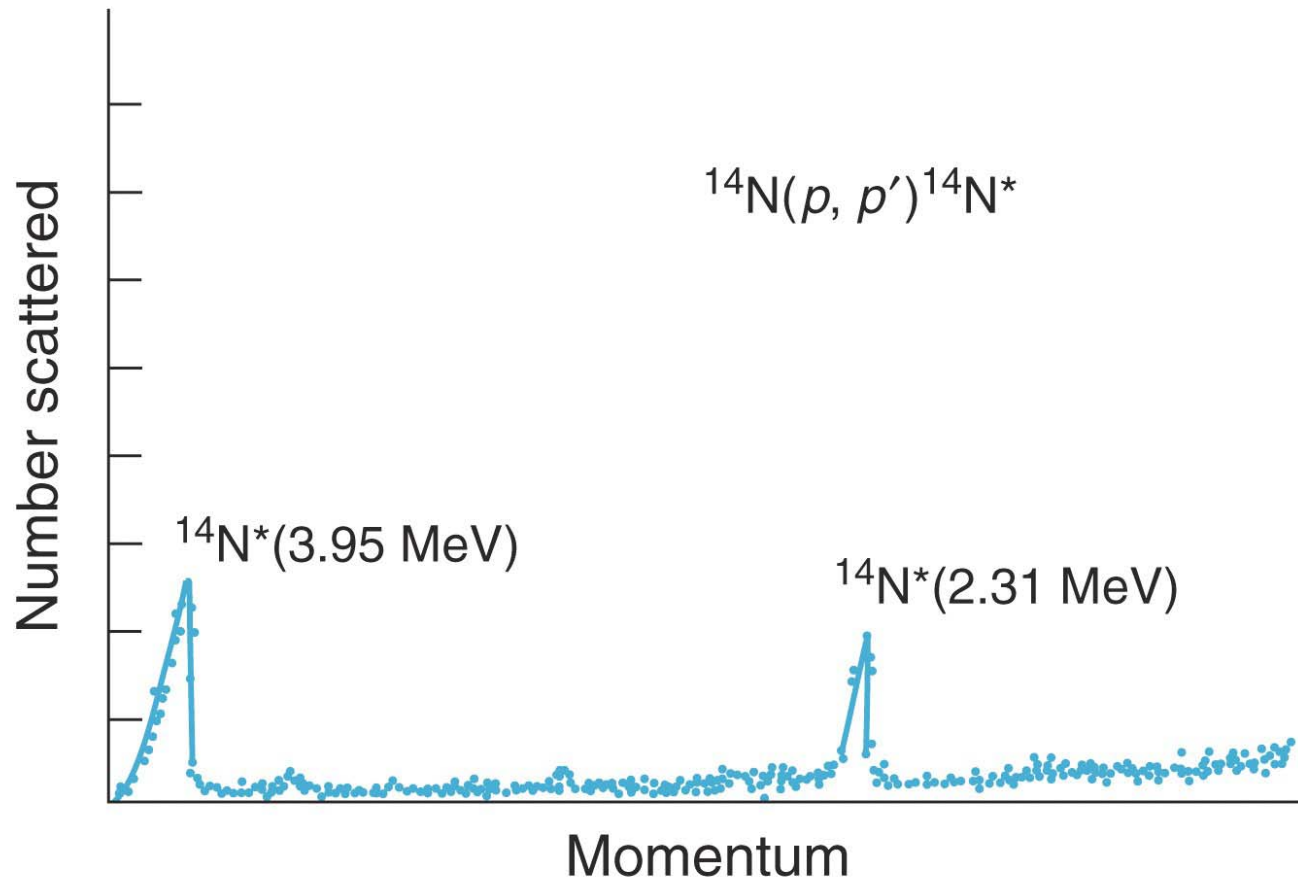


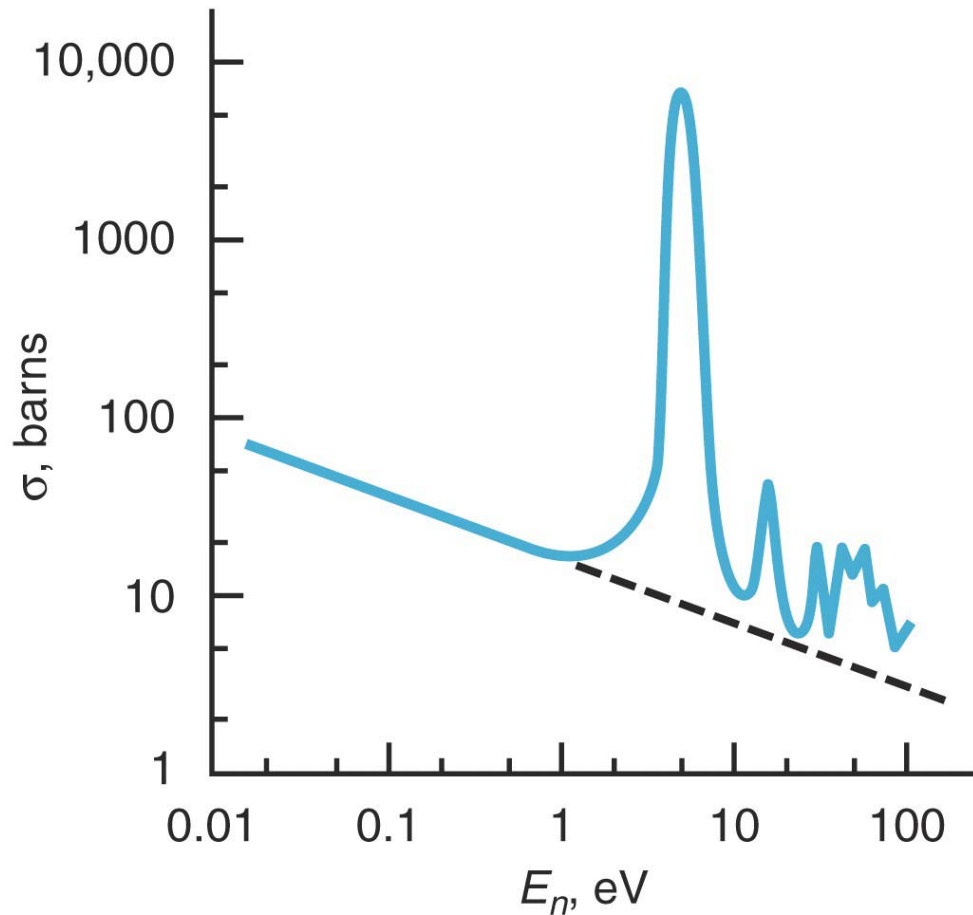
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The Joint Institute for Nuclear Astrophysics

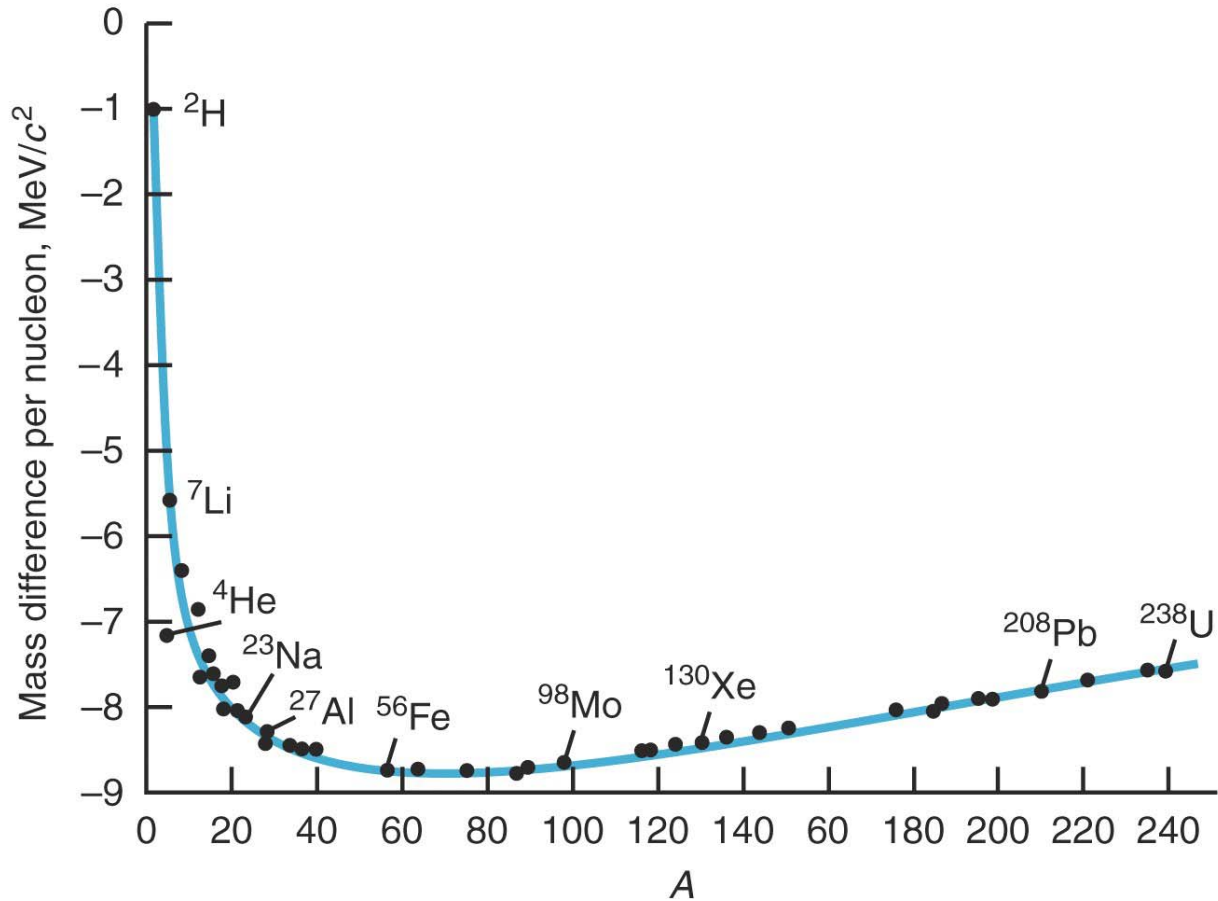
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The same kind of information about excited states of ^{14}N can be obtained by inelastic proton scattering

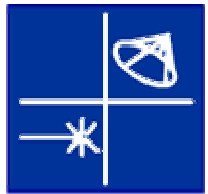




The effect of resonances on the cross section (here the neutron-capture on silver, can be quite dramatic. The dashed line is an extension of the $1/v$ behavior expected in the absence of resonances



Before we leave this, look at the flipped version of the plot of B/A vs A that we saw earlier. Note that the rest energy per nucleon is less for intermediate mass nuclei than for very heavy or light ones....the key to fission



The Nuclear Force

About a hundred times stronger than the Coulomb force

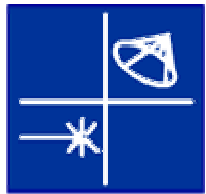
Very short range—goes to zero beyond about 3 fm

Charge independent—does not matter if the particles are protons or neutrons

Saturated—is constant at about 8 MeV/nucleon above $A=20$ or so

Depends on the spin orientation of the nucleons

Suspected to be an exchange force in which the attraction is due to an exchange of pions



The nuclear shell model

It is an independent-particle model, similar to that used for assigning energy states to atomic electrons, but one that makes use of a strong spin-orbit coupling for each nucleon. It accounts for the shell-like structure of protons and neutrons and explains the ‘magic numbers’

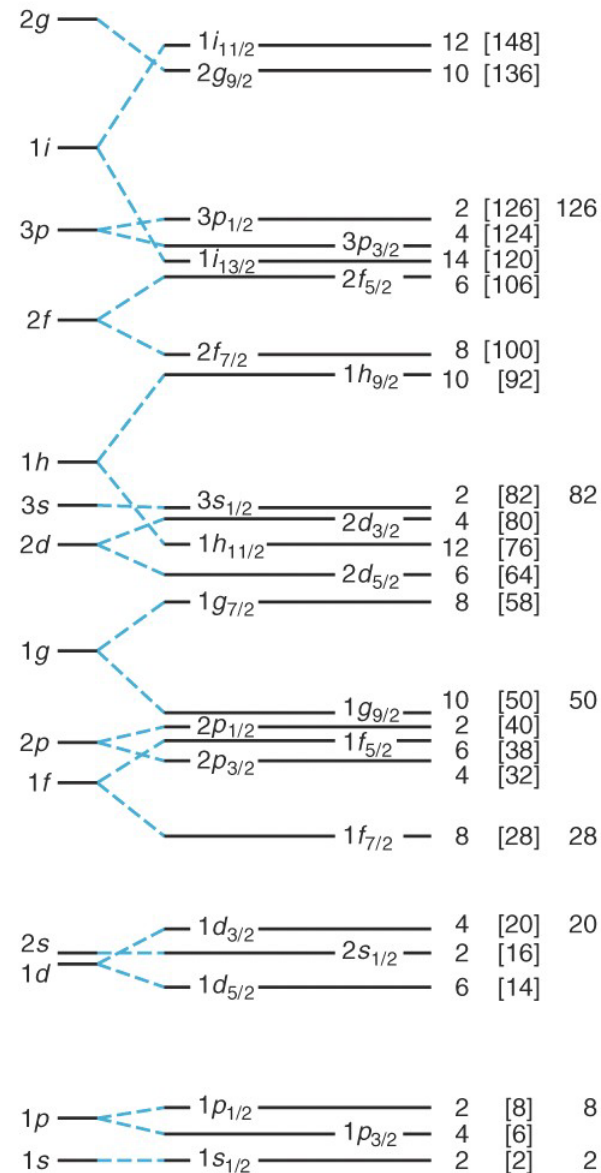
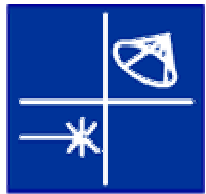


TABLE 11-5 Angular momenta and magnetic moments of selected odd-*A* nuclei

Isotope	Number of odd particles	<i>Z</i> or <i>N</i> , a magic number	Predicted level	Measured spin	Measured magnetic moment (μ_N)
${}^{11}_3\text{B}_6$	5	—	$p_{3/2}$	3/2	+2.689
${}^{13}_6\text{C}_7$	7	—	$p_{1/2}$	1/2	+0.702
${}^{15}_7\text{N}_8$	7	<i>N</i>	$p_{1/2}$	1/2	−0.283
${}^{17}_8\text{O}_9$	9	<i>Z</i>	$d_{5/2}$	5/2	−1.894
${}^{17}_9\text{F}_8$	9	<i>N</i>	$d_{5/2}$	5/2	+4.722
${}^{27}_{13}\text{Al}_{14}$	13	—	$d_{5/2}$	5/2	+3.641
${}^{39}_{19}\text{K}_{20}$	19	<i>N</i>	$d_{3/2}$	3/2	+0.09
${}^{41}_{20}\text{Ca}_{21}$	21	<i>Z</i>	$f_{7/2}$	7/2	−1.595
${}^{41}_{21}\text{Sc}_{20}$	21	<i>N</i>	$f_{7/2}$	7/2	—
${}^{57}_{28}\text{Ni}_{29}$	29	<i>Z</i>	$p_{3/2}$	3/2	—
${}^{91}_{40}\text{Zr}_{51}$	51	—	$g_{7/2}$	7/2	−1.303
${}^{115}_{49}\text{In}_{66}$	49	—	$g_{9/2}$	9/2	—
${}^{205}_{81}\text{Tl}_{124}$	81	—	$s_{1/2}$	1/2	+1.628
${}^{209}_{83}\text{Bi}_{126}$	83	<i>N</i>	$h_{9/2}$	9/2	+4.080



Summary

It's been a quick trip, and I hope a not-too-boring one. I tried to touch on a number of topics:

Ground-state properties of the nucleus

Radioactivity

Nuclear reactions

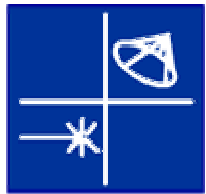
Decay modes of an excited nucleus, and

The nuclear force

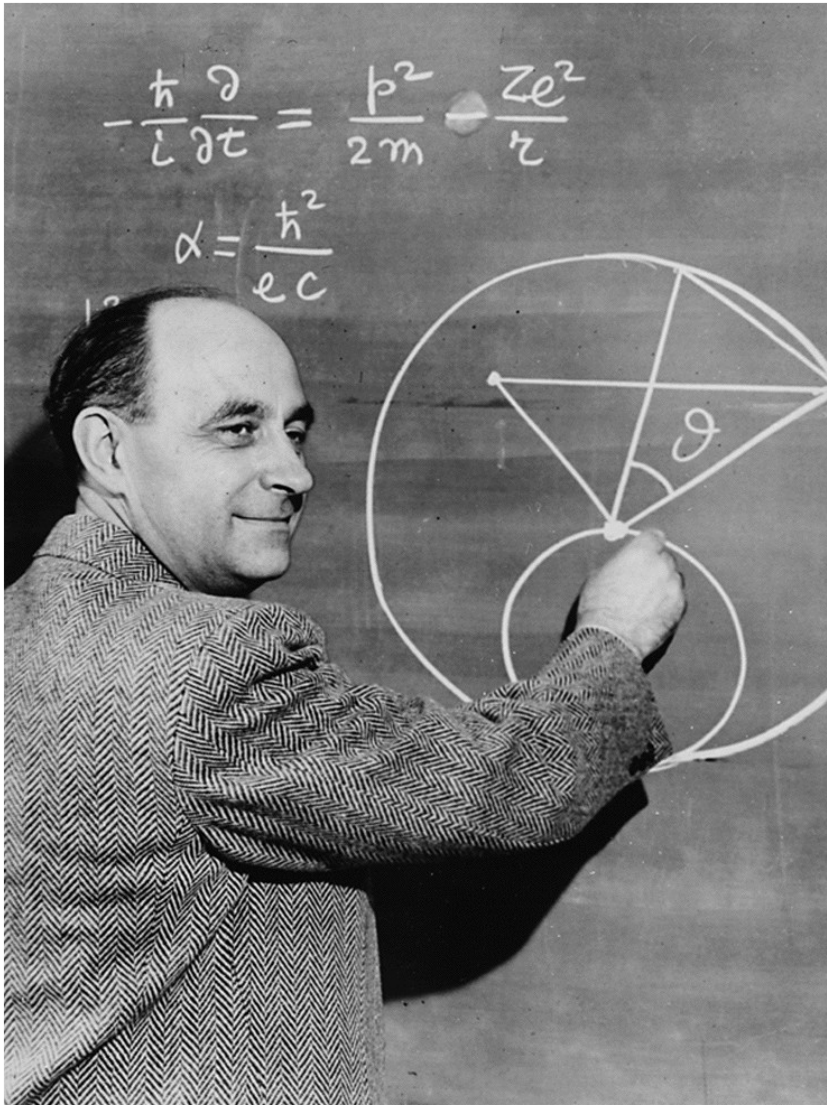
I hope that there was something in there that you found interesting.

If questions arise during the school year, please call (574.631.8591) or drop me a note at ahyder@nd.edu

Tony Hyder

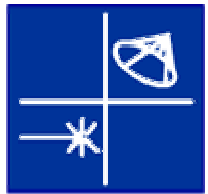


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Bonus:

- Who is this?
- What's wrong with this picture?



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