

Observation of the Giant monopole resonances in the Sn isotopes via (α, α') reactions at 400 MeV at RCNP
- Latest results and implications -

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Contents

- 1) Our aim : Incompressibility and Giant Resonances
(ISGMR, ISGDR)**
- 2) Experiments**
- 3) Analysis with MDA**
- 4) Results: Peak positions of the ISGMR in the Sn
isotopes with $A=112 - 124$.**
- 5) Summary**

Our Aim

In the supernova explosion processes, the iron core absorbs electrons via the electron capture process, and the core is dominated with neutron excess nuclei.

→ This process proceeds at? ρ_0

→ In the region? ρ_0 , the core become hard

→ Rebounding happens

→ Explosion



<http://heasarc.gsfc.nasa.gov/docs/snr.html>

Thus, an important factor for supernova explosion is the hardness of the core with neutron excess nuclei.

Incompressibility

Isotope dependence of Incompressibility



$$K_{sym} \left(\frac{N-Z}{A} \right)^2$$

Two Major Unsolved Issues in Nuclear Incompressibility

1. Different K_A (K_{\dots}) values from ISGMR and ISGDR

$$E_{ISGMR} \doteq \sqrt{\frac{K_A}{m\langle r^2 \rangle}}$$

$$E_{ISGDR} \doteq \sqrt{\frac{7K_A + (27/25)\epsilon_F}{3m\langle r^2 \rangle}}$$

2. From the same GMR data, Non-relativistic and Relativistic calculations gave different K_{\dots} values;

220 MeV non-rel.

270 MeV rel.

The first of these has been resolved. With the background-free spectra, ISGDR strength at higher E_x than before. Now, same calculations give reasonable agreement with E_{GMR} and E_{ISGDR} .

The second issue still remained unsolved.

	$E_{\text{GMR-NR.}}$	$E_{\text{GMR-Rel.}}$
^{112}Sn	17.7 MeV	18.3 MeV
^{124}Sn	16.9 MeV	17.0 MeV
$\Delta[^{124}\text{Sn}-^{112}\text{Sn}]$	0.8 MeV	1.3 MeV

We need precise numbers for E_{GMR} for the whole series of Sn isotopes to fully constrain the values of K_{sym} .

E_{GMR} for several Sn isotopes with uncertainties of 0.1 MeV



High statistics data required.



to be well done.

There was the consensus among the theorists that the Primary difference between the non-relativistic and relativistic calculations comes from the “**symmetry energy**” term.

$$K_A \sim K_{\dots} (1 + cA^{-1/3}) + K_{\text{sym}} ((N-Z)/A)^2 + K_{\text{Coul}} Z^2 A^{-4/3}$$



$K_{\text{sym}} = -400 \sim +466 \text{ MeV}$; not well obtained

B.A. Li, PRL 85, 4221 (2000),

B.A. Li, C.M.Ko, and W. Bauer, Int. J. Mod. Phys. E7, 147 (1998).

B.A.Li, W.Udo, Nova Science Publishers.

R.J. Furnstahl, nucl-th/0112085.

Clearly the (N-Z)/A term is very important

in nuclear structure, heavy ion collision, astronuclearphysics.

The widest range of (N-Z)/A in an isotope series (in medium and heavy mass nuclei) is in Sn:

^{112}Sn **0.107**

^{124}Sn **0.194**

Incompressibility of nuclear matter



Impossible to be determined from the observation.



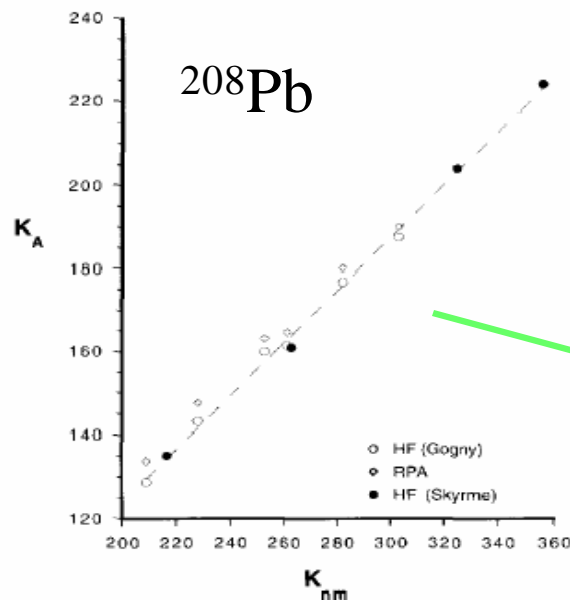
Need help from theory

$$K_{\infty} = \left[9r^2 \frac{d^2(E/A)}{dr^2} \right]_{r=r_0}$$

E/A : Binding energy /A

r : nuclear density

r_0 : nuclear density at equilibrium



- 1 . Determine K_A for finite nuclei.
- 2 . Obtain the relationship between K_A and K_8

K_A is obtained from the information on the excitation energy of ISGMR , ISGDR) .

$$E_{ISGMR} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

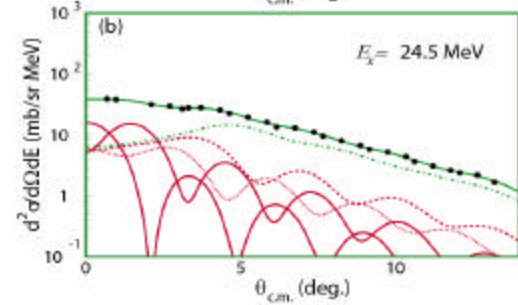
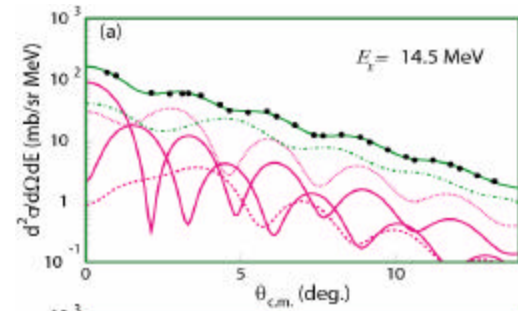
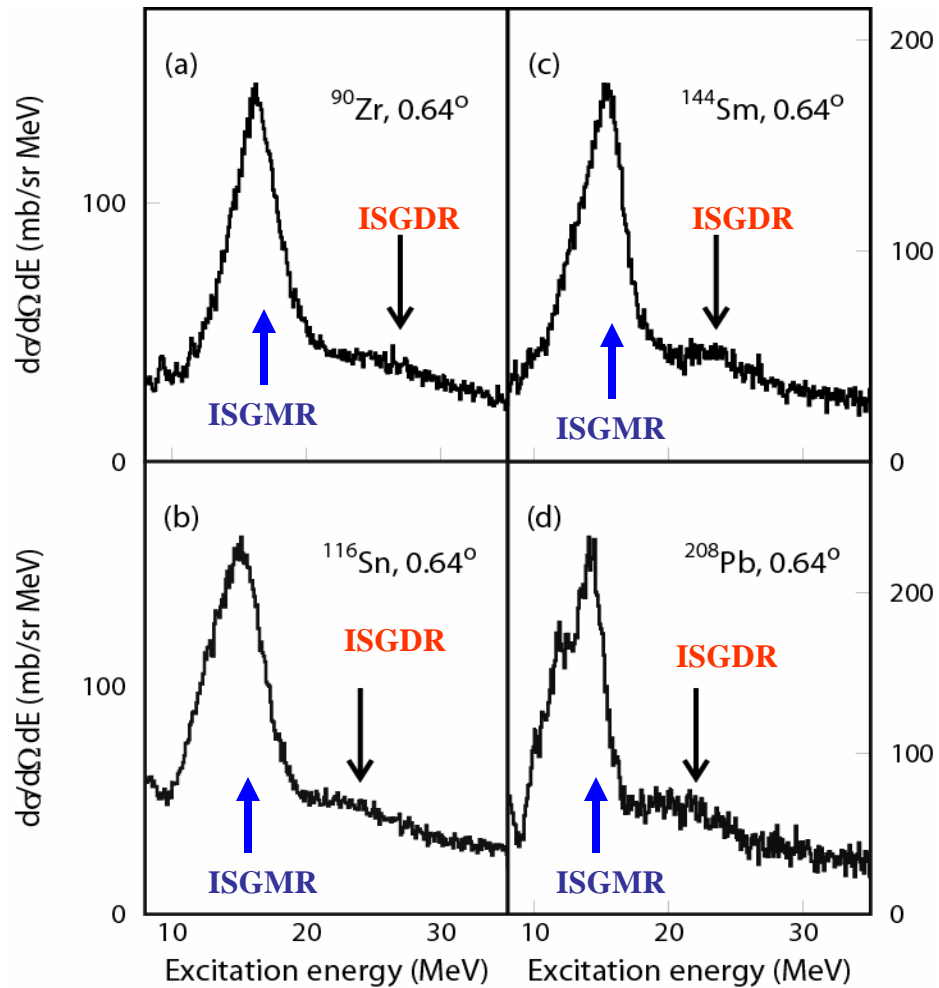
$$E_{ISGDR} = \hbar \sqrt{\frac{3 K_A + (27/25)e_F}{7 m \langle r^2 \rangle}}$$

Relation between K_A and K_8 ,

$$K_A = 0.64K_8 - 3.5$$

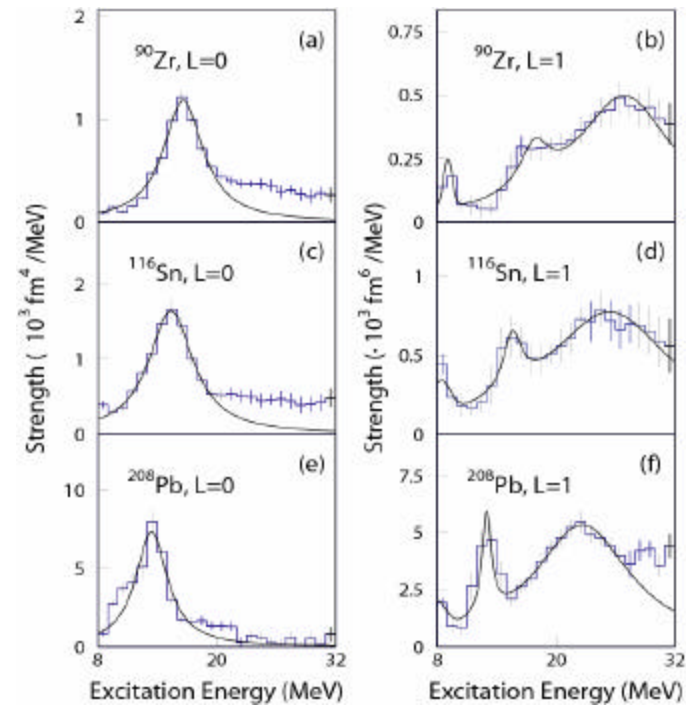
(J.P.Blaizot, NPA591,435,1995)

(α, α') spectra at 386 MeV



^{116}Sn

MDA results for L=0 and L=1



Giant Resonances

• (? T=0)

L=0 ISGMR }
L=1 ISGDR }

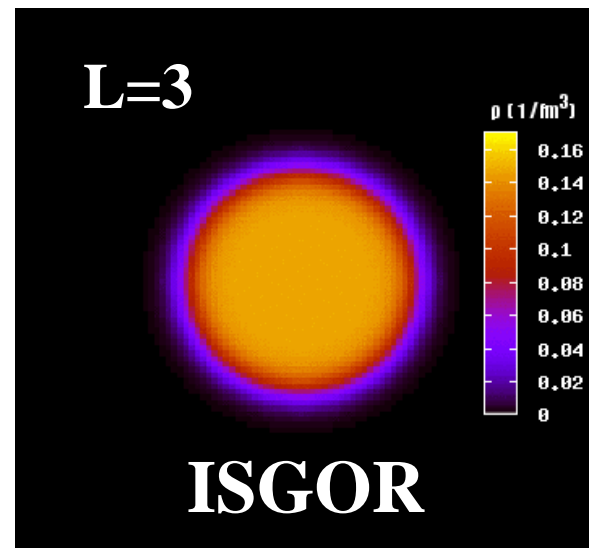
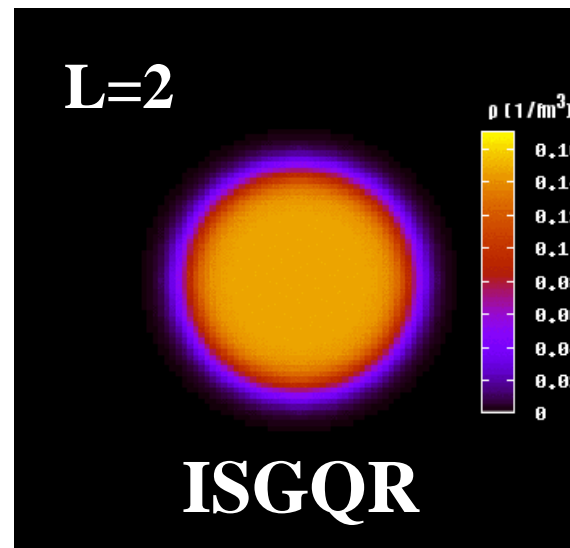
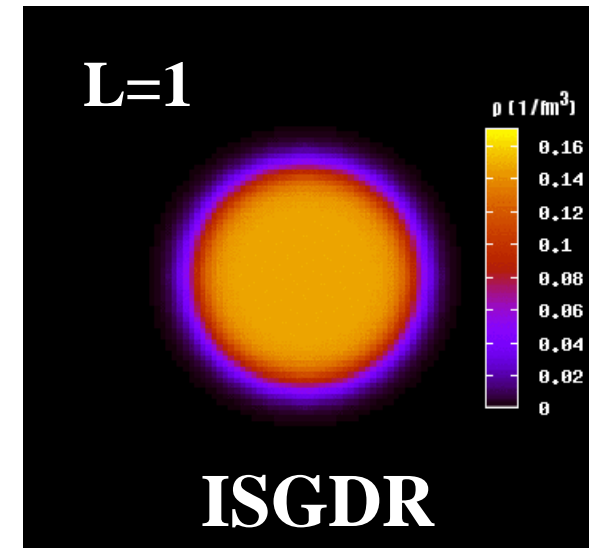
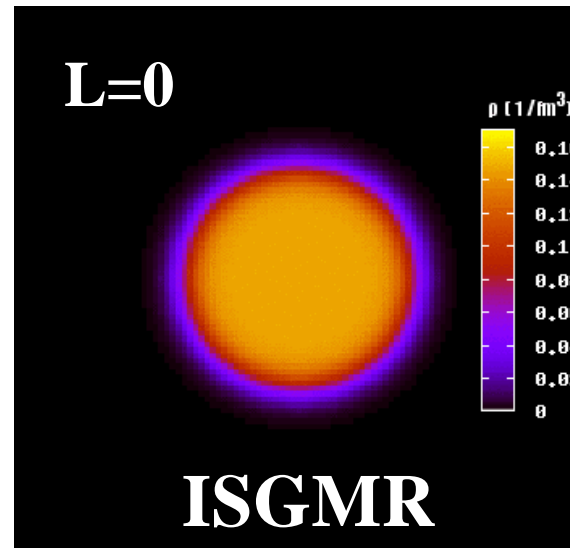
L=2 ISGQR

L=3 ISGOR

⋮ ⋮

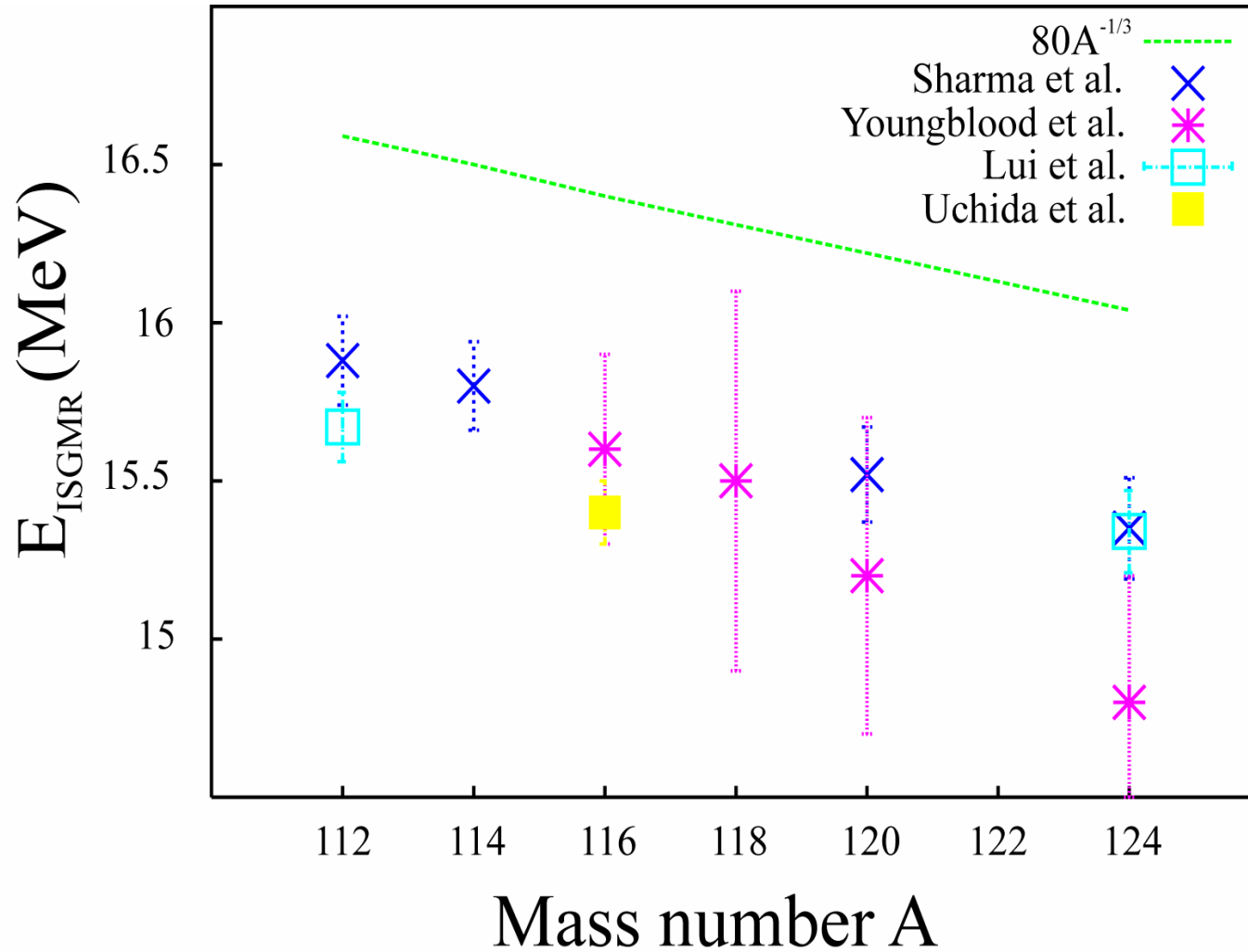
• (? T=1)

L=1 IVGDR



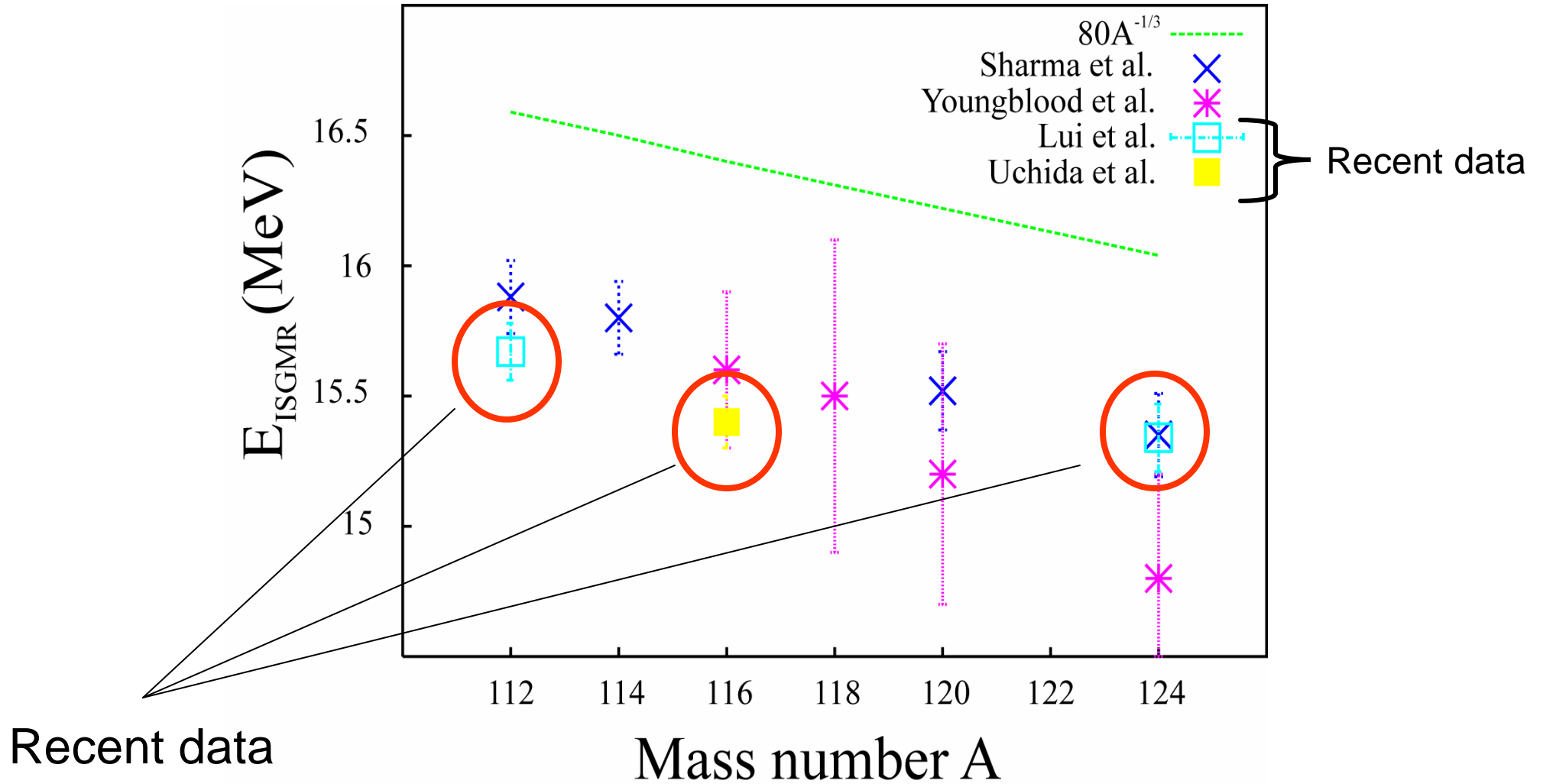
Experimental data on ISGMR

ISGMR energy E_{ISGMR}



Experimental data on ISGMR

ISGMR energy E_{ISGMR}

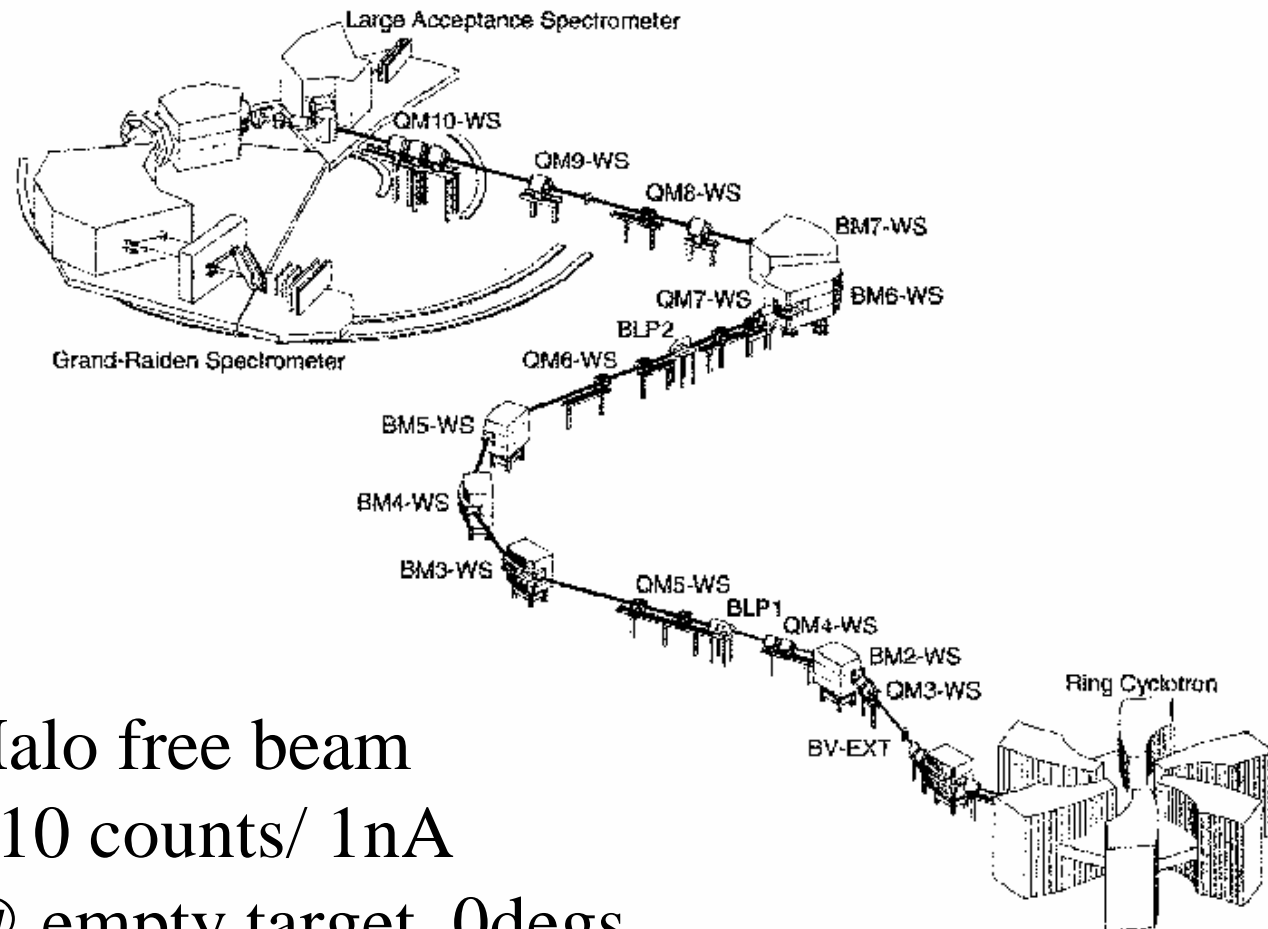


In the present experiment

- Obtain good (α, α') spectra including 0 degrees for ^{112}Sn to ^{124}Sn
- Using MDA analyses, we obtain the $L=0$ cross section distribution for ISGMR, and determine the peak location in excitation energy.
- Obtain K_{sym} .

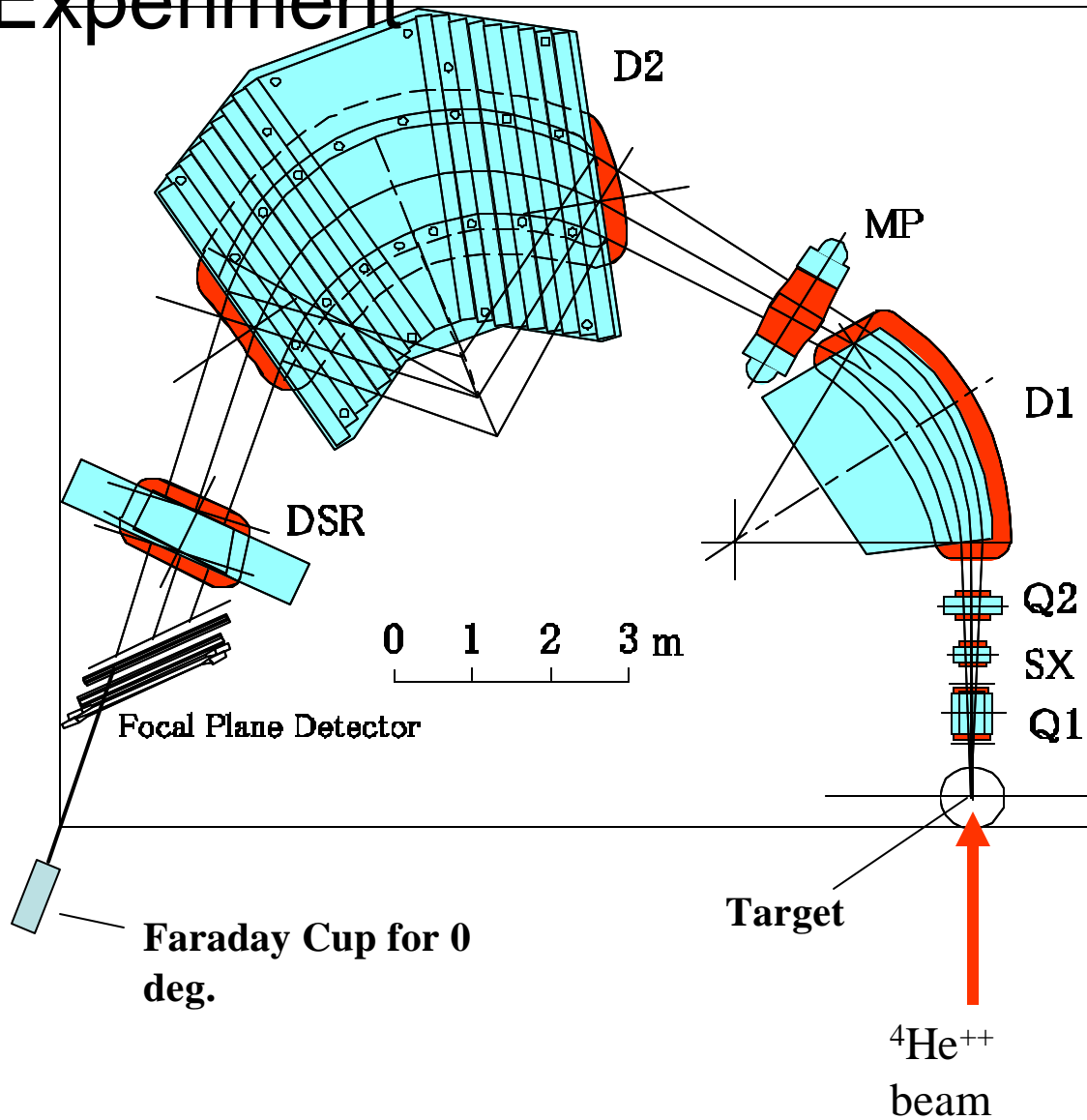
RCNP ring cyclotron facility

$$E_{\alpha}=386 \text{ MeV}$$



Halo free beam
~10 counts/ 1nA
@ empty target, 0degs.

Experiment



Incidence ${}^4\text{He}^{++}$ beam

Energy 386 MeV

Measured angles

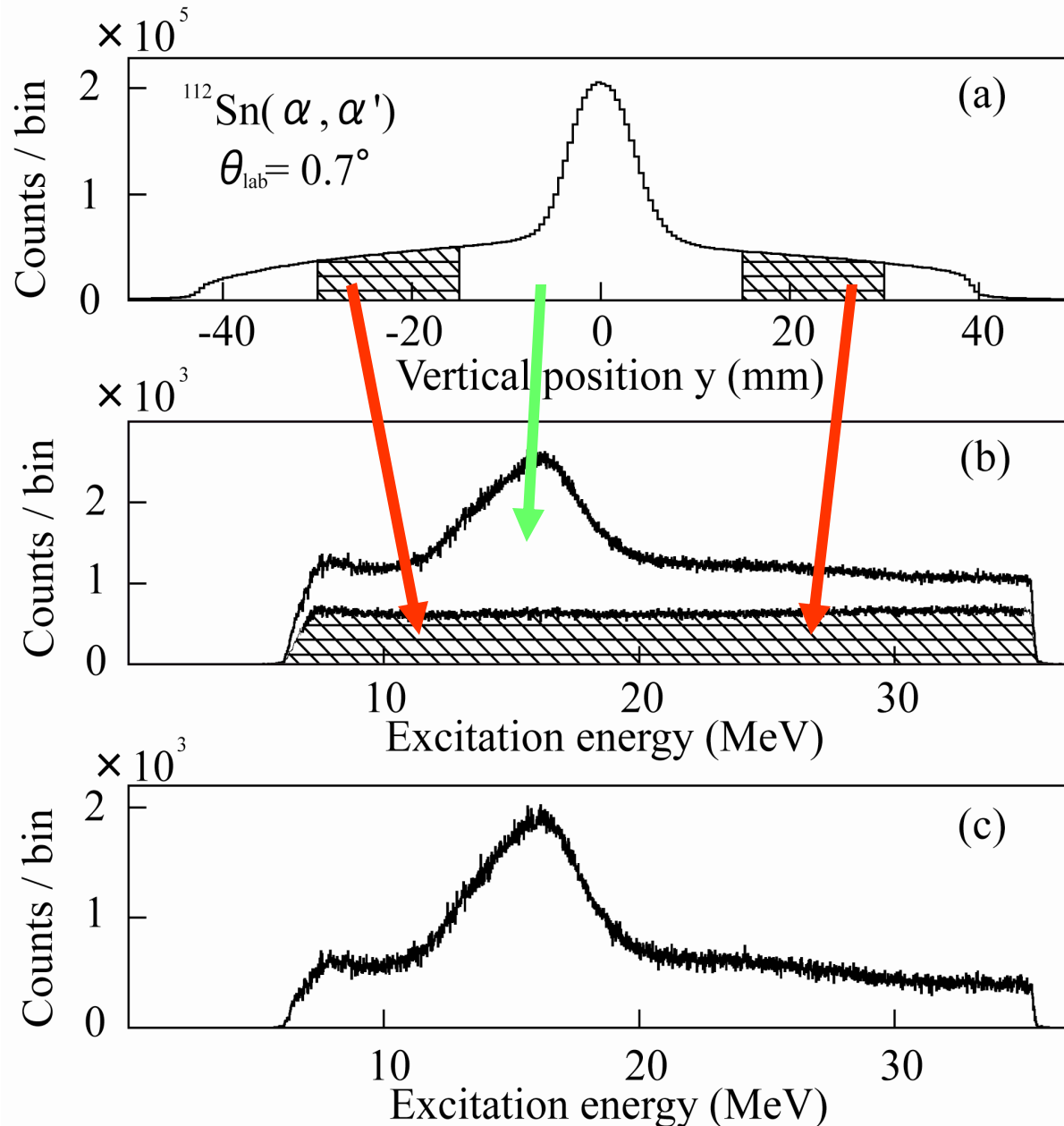
1 1 2 ? 1 1 4 ? 1 1 8 ? 1 2 0 ? 1 2 2 ? 1 2 4 Sn

$0^\circ \sim 8^\circ$

${}^{116}\text{Sn}$ Uchida et al.

Targets	Thickness (mg/cm ²)	Enrichment (%)
${}^{112}\text{Sn}$	5.5	99.5
${}^{114}\text{Sn}$	7.5	87.1
${}^{116}\text{Sn}$	10.0	95.6
${}^{118}\text{Sn}$	5.7	96.5
${}^{120}\text{Sn}$	5.1	98.4
${}^{122}\text{Sn}$	9.3	92.2
${}^{124}\text{Sn}$	5.0	96.7

Data Analysis



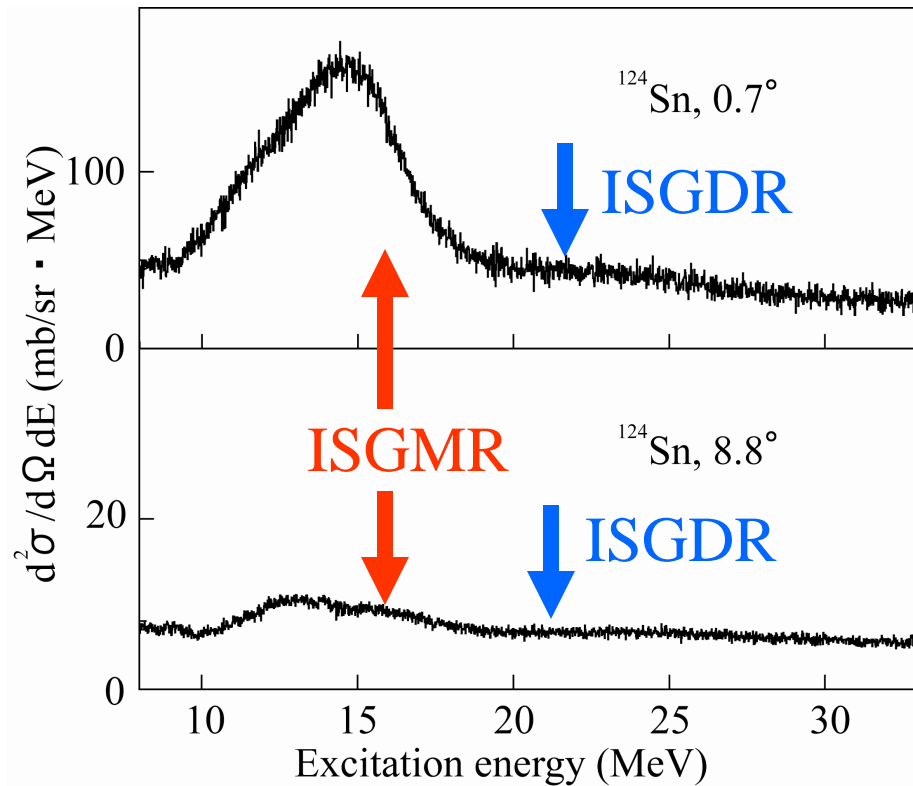
Background rejection with the focal plane detector system of the spectrometer Grand Raiden.

(a) one-dimensional spectrum along the vertical direction. Background events correspond to the hatched area. True and background events are in the central region.

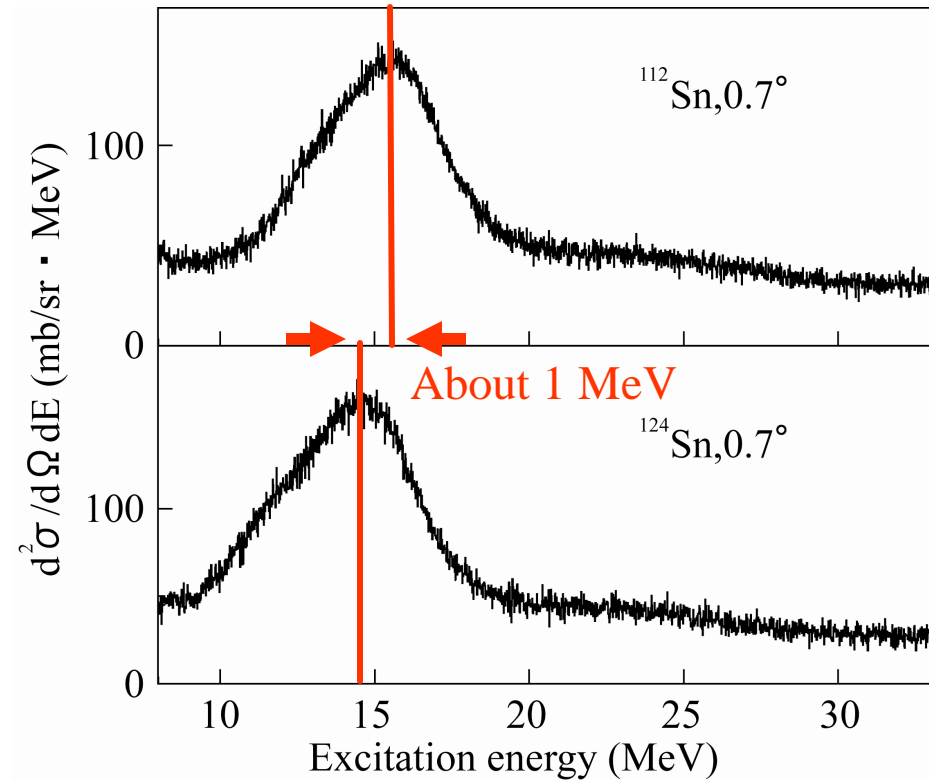
(b) The energy spectra for the true + background events, and for the background Events.

(c) Difference spectrum for true events.

$^{124}\text{Sn}(\alpha, \alpha')$ spectra



$^{112,124}\text{Sn}(\alpha, \alpha')$ spectra

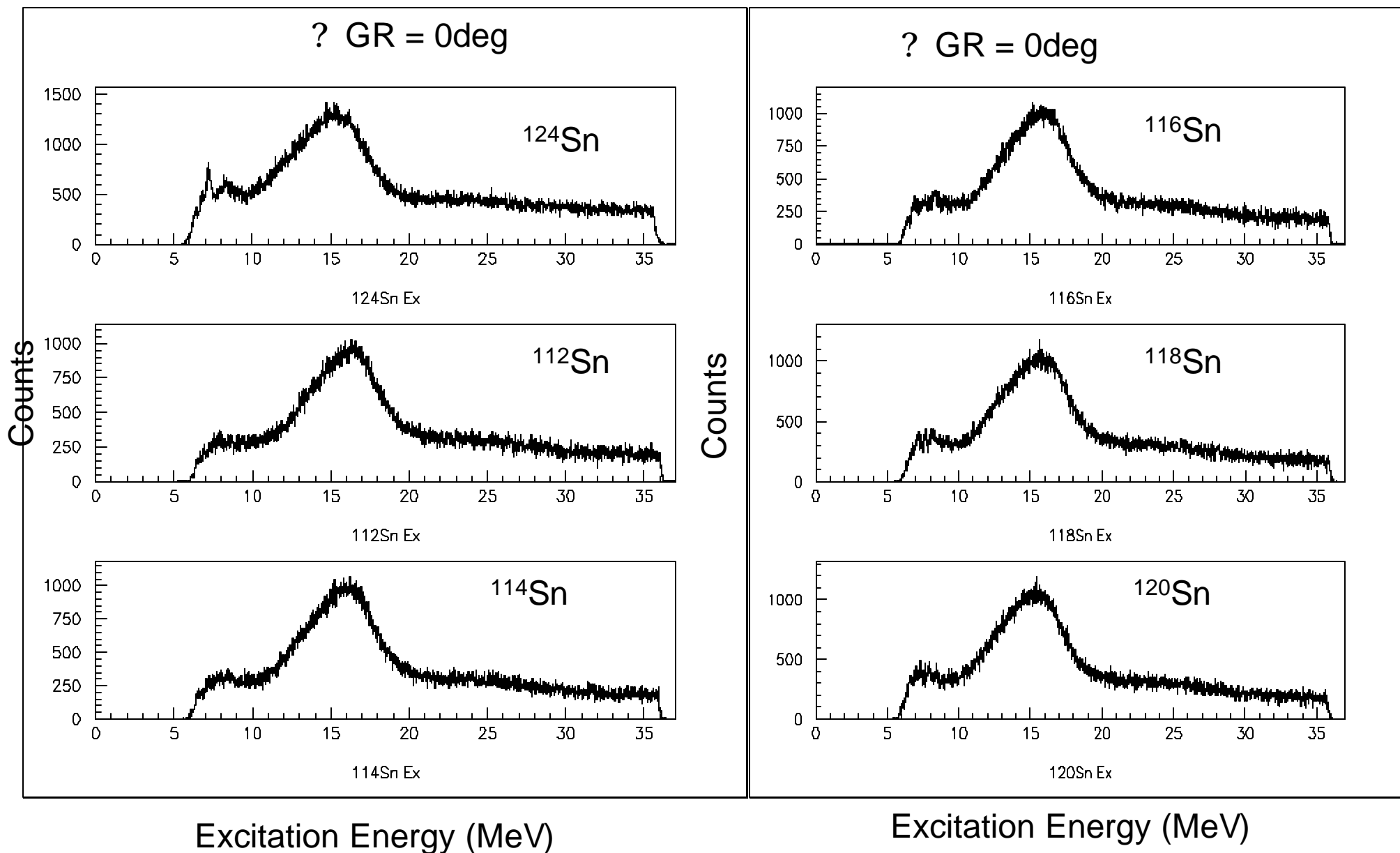


Superposition of components with various L transfer in the $\text{Sn}(\alpha, \alpha')$ spectra.

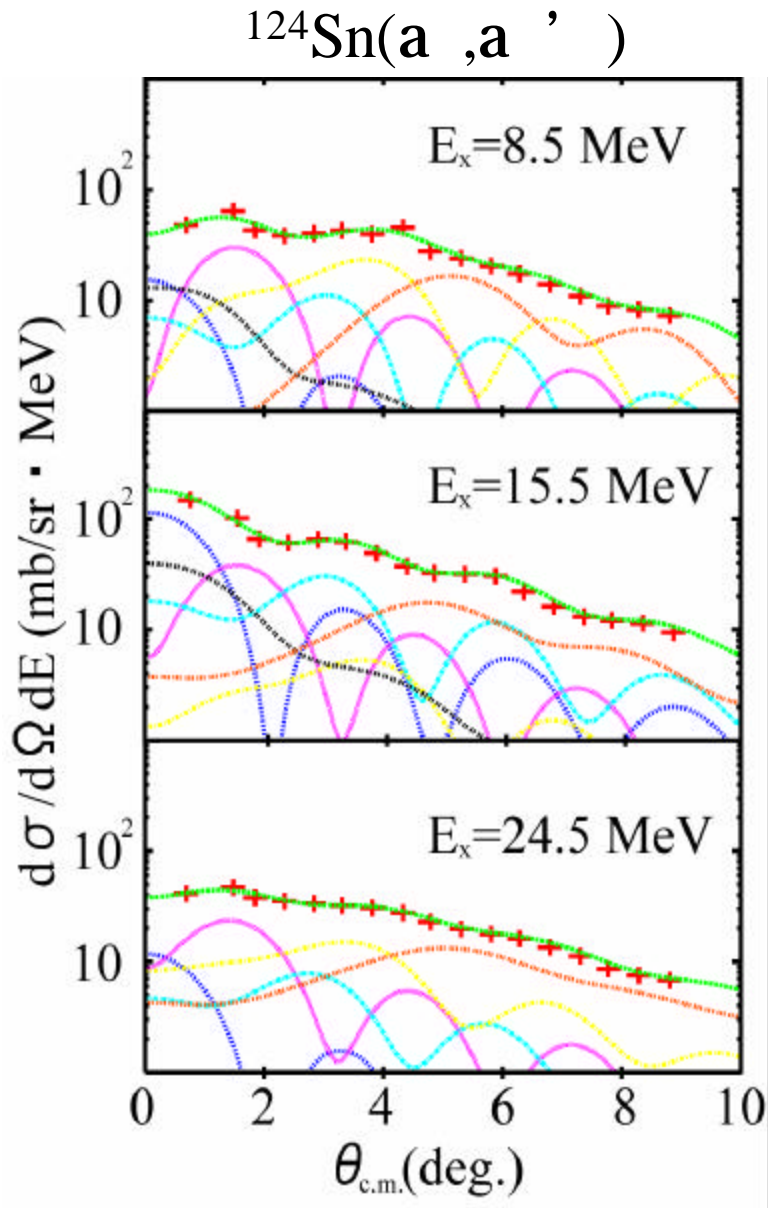


In order to extract the ISGMR peak position, it is necessary for us to extract the $L=0$ component from the excitation energy spectra.

Energy Spectra



Multipole-decomposition analysis



$$\mathbf{s}^{\text{exp}}(\mathbf{q}, E_x) = \sum_L a_L(E_x) \mathbf{s}_L^{\text{calc}}(\mathbf{q}, E_x)$$

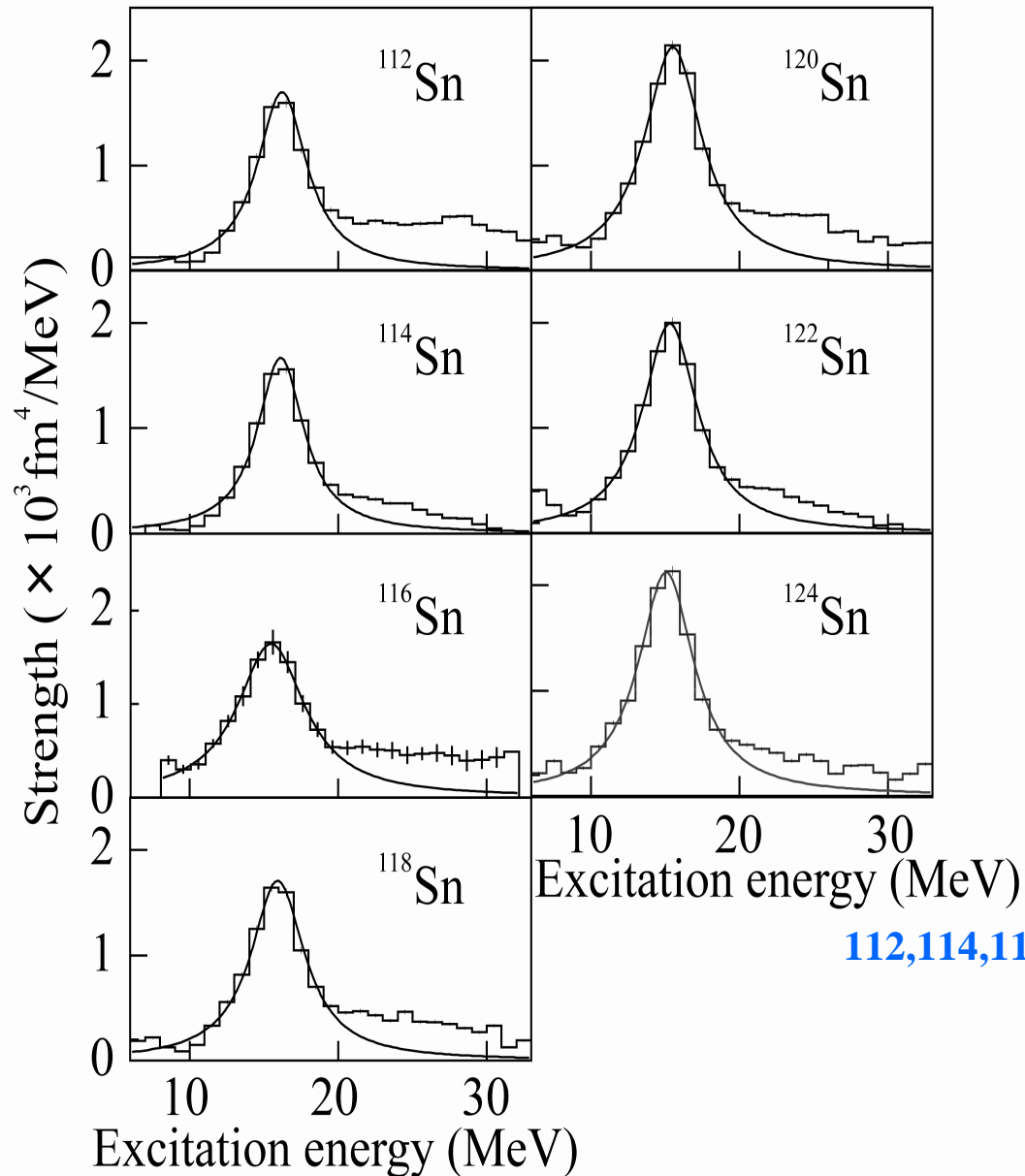
↑
Cross sections

↑
DWBA calculations

(L=0~ 15) and IVGDR cross section

- ? L = 0
- ? L = 3
- ? L = 1
- ? L > 3
- ? L = 2
- ? L = 1, ? T = 1

L=0



Breit-Wigner function

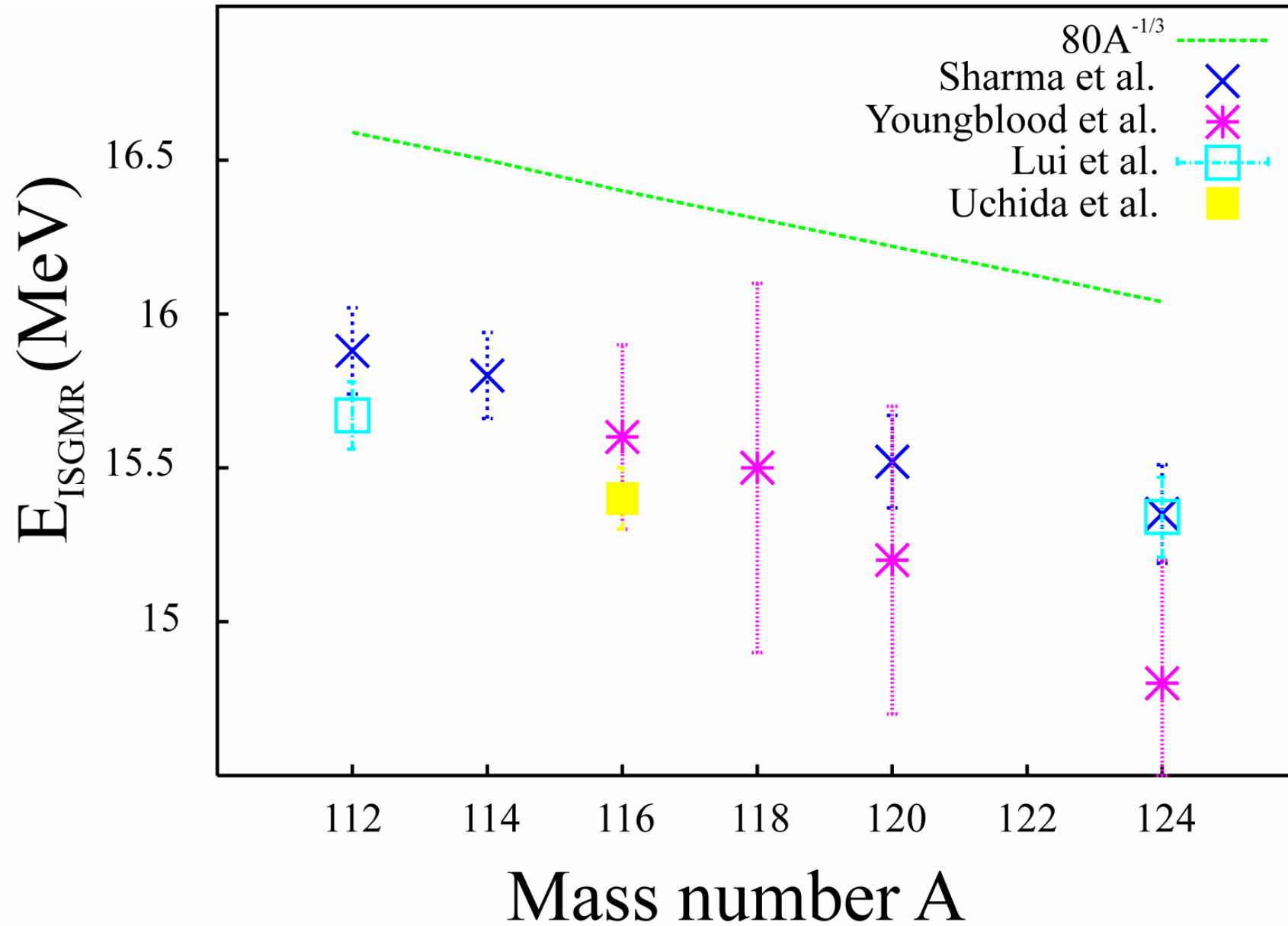
$$S(E) = \frac{S_m}{(E - E_m)^2 + \Gamma_m}$$

	E_m (MeV)	G_m (MeV)
^{112}Sn	16.1 ± 0.1	4.4 ± 0.2
^{114}Sn	16.0 ± 0.1	4.0 ± 0.1
^{116}Sn	15.4 ± 0.1	5.7 ± 0.3
^{118}Sn	15.7 ± 0.1	4.6 ± 0.2
^{120}Sn	15.3 ± 0.1	4.8 ± 0.2
^{122}Sn	15.1 ± 0.1	4.5 ± 0.2
^{124}Sn	14.9 ± 0.1	4.5 ± 0.3

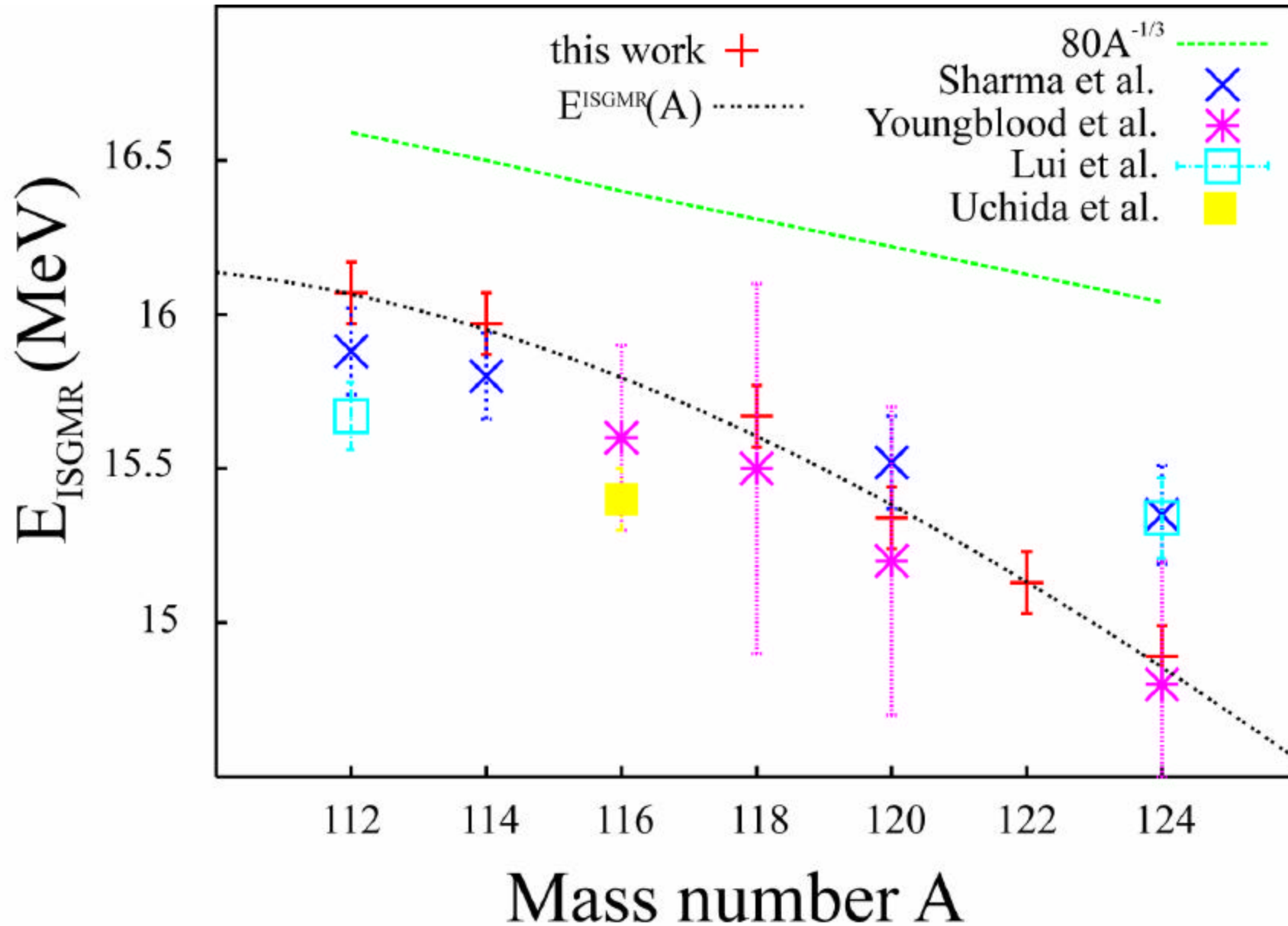
$^{112,114,118,120,122,124}\text{Sn}$: this work

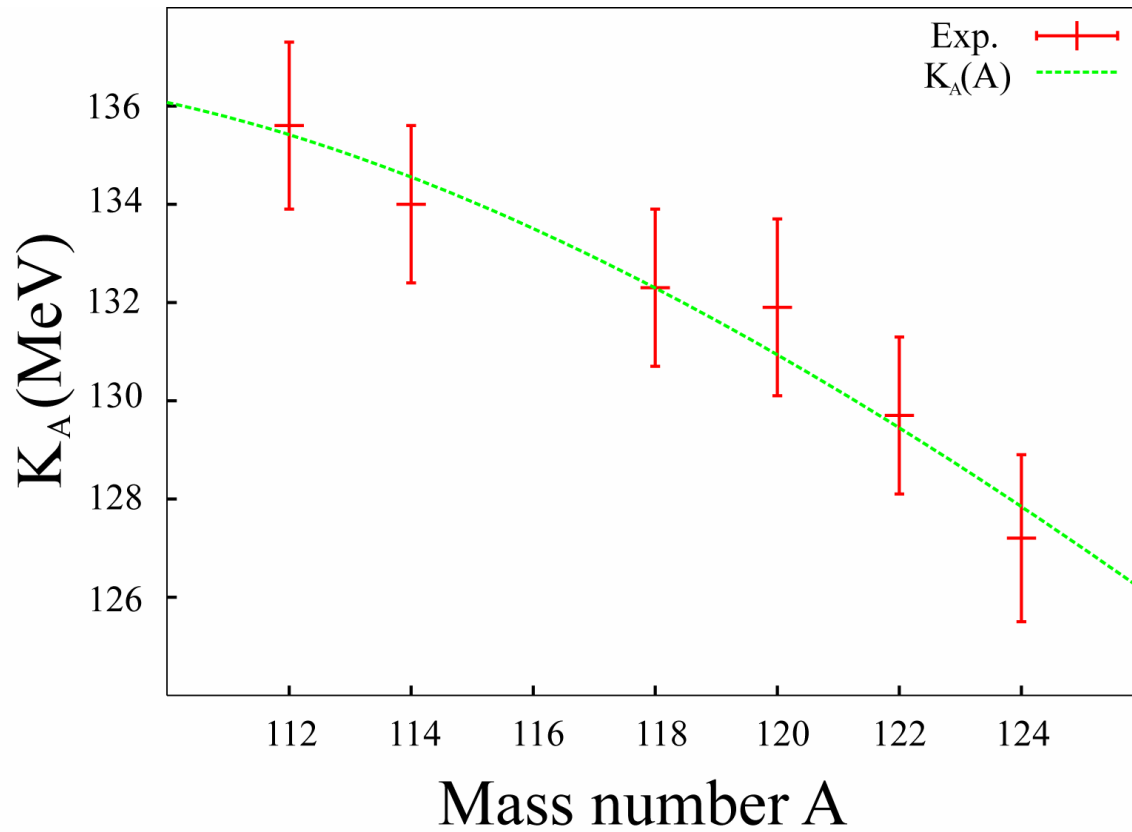
^{116}Sn : Uchida et al.

ISGMR energy E_{ISGMR}



ISGMR energy E_{ISGMR}





$$E_{ISGMR} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

$$K_A \sim K_{\dots} (1 + cA^{-1/3}) + K_{\text{sym}} ((N-Z)/A)^2 + K_{\text{Coul}} Z^2 A^{-4/3}$$

$$c = -1$$

$$K_{\text{Coul}} = -5$$

G. Colo et al. PRC **70** 024307
(2004)

$$-580 < K_{\text{sym}} < -380$$

Summary

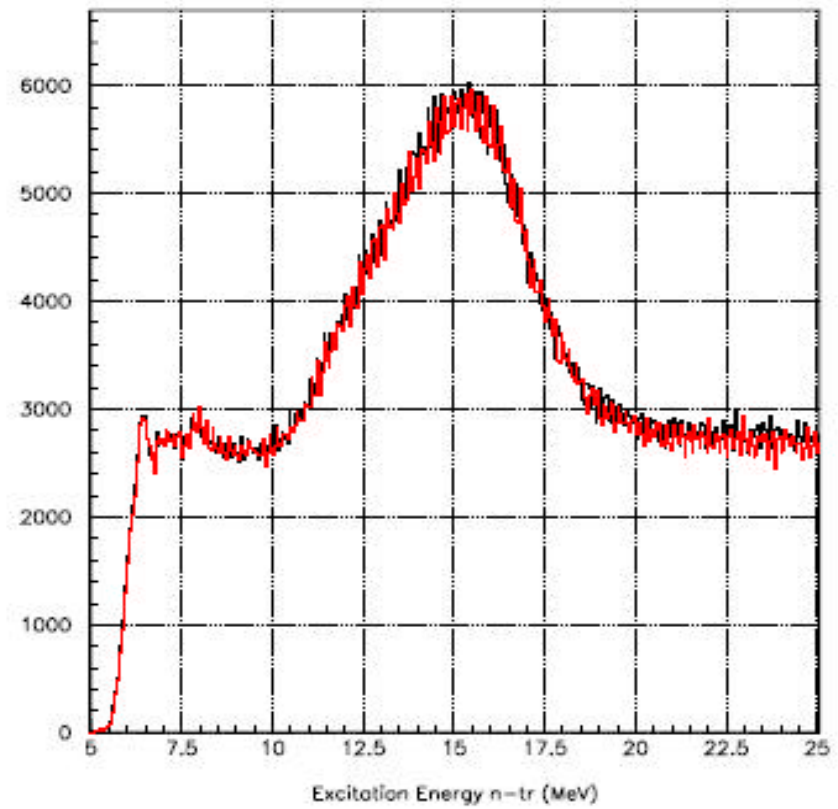
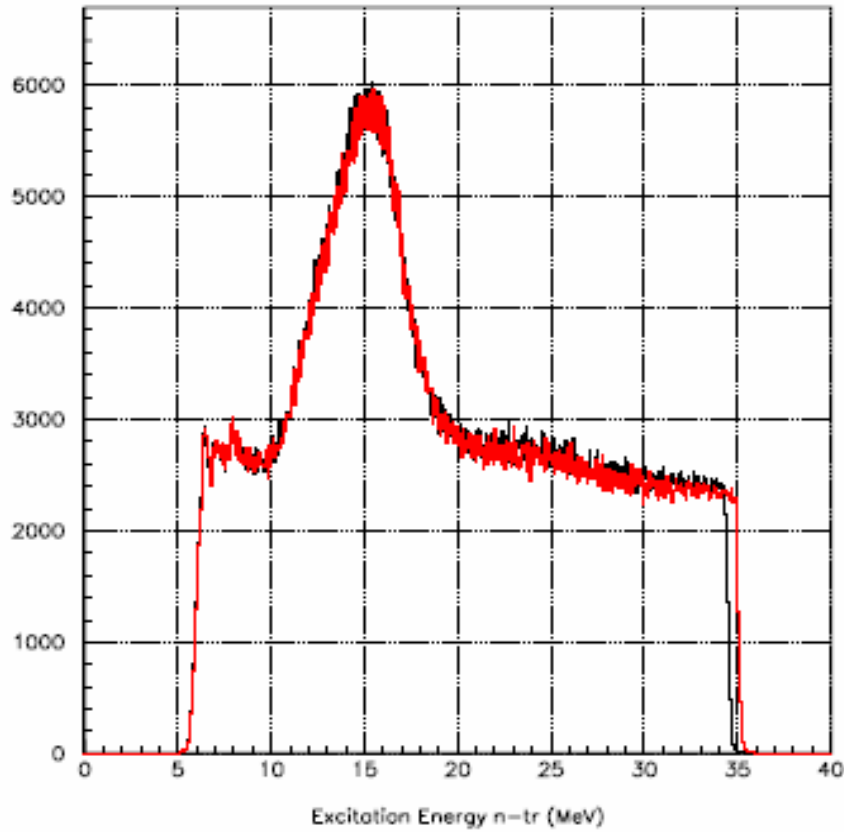
- 1) ISGMR in $^{112,114,118,120,122,124}\text{Sn}$ via (α, α')
- 2) We obtained the ISGMR cross section distribution and peak positions.
- 3) K_{sym} is most likely $-580 < K_{\text{sym}} < -380$ MeV.

	$E_{\text{GMR-NR.}}$	$E_{\text{GMR-Rel.}}$	Exp.
^{112}Sn	17.7 MeV	18.3 MeV	16.1 MeV
^{124}Sn	16.9 MeV	17.0 MeV	14.9 MeV
$\Delta[^{124}\text{Sn}-^{112}\text{Sn}]$	0.8 MeV	1.3 MeV	1.2 MeV

Comparison between two spectra calibrated by Ex and p.

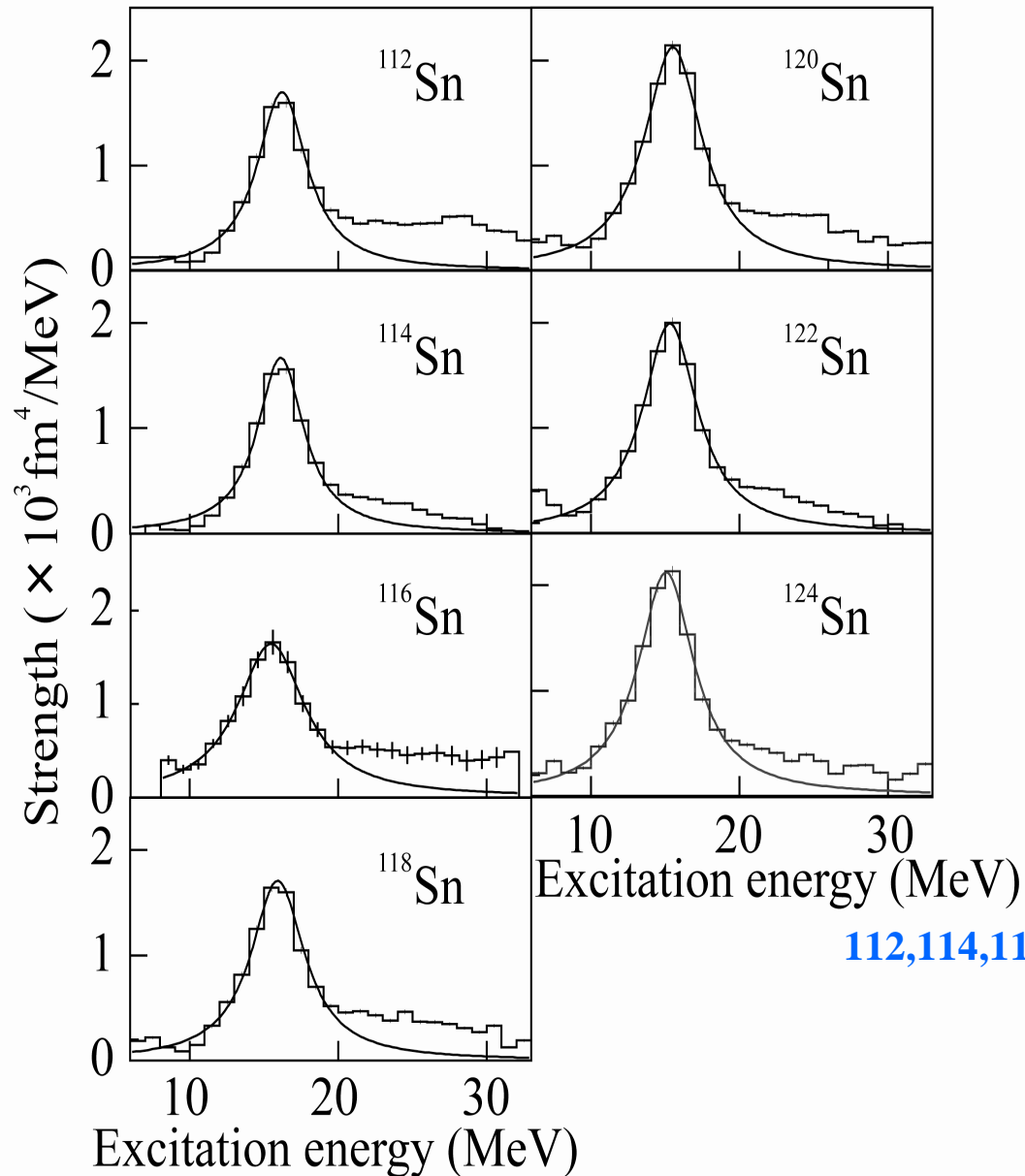
2005/05/06 13.35

2005/05/06 13.35



Red: Ex
Black: momentum (correct method)

L=0



Breit-Wigner function

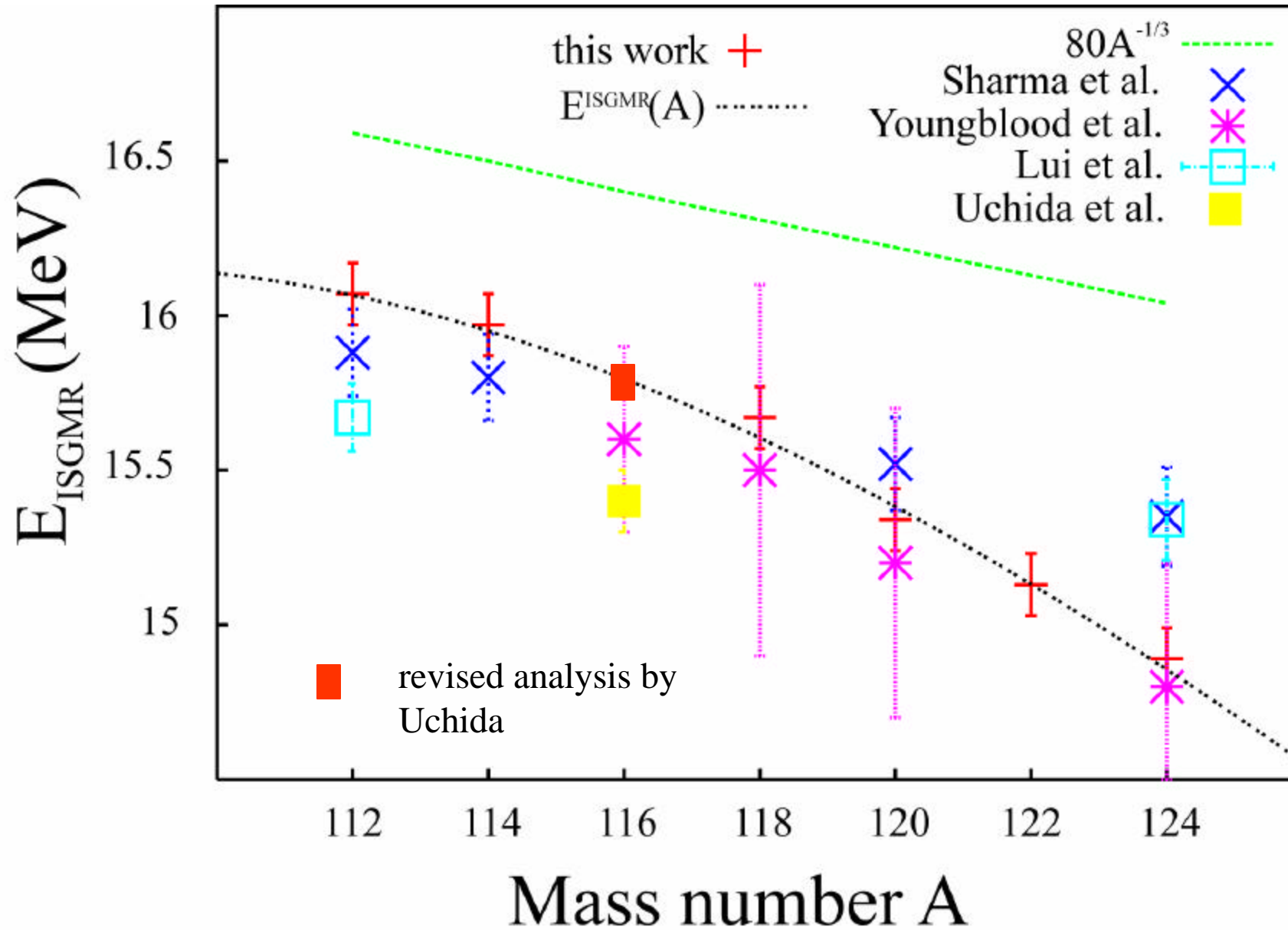
$$S(E) = \frac{S_m}{(E - E_m)^2 + \Gamma_m}$$

	E_m (MeV)	G_m (MeV)
^{112}Sn	16.1 ± 0.1	4.4 ± 0.2
^{114}Sn	16.0 ± 0.1	4.0 ± 0.1
^{116}Sn	15.8 ± 0.1	4.9 ± 0.3
^{118}Sn	15.7 ± 0.1	4.6 ± 0.2
^{120}Sn	15.3 ± 0.1	4.8 ± 0.2
^{122}Sn	15.1 ± 0.1	4.5 ± 0.2
^{124}Sn	14.9 ± 0.1	4.5 ± 0.3

$^{112,114,118,120,122,124}\text{Sn}$: this work

^{116}Sn : Uchida et al.

ISGMR energy E_{ISGMR}

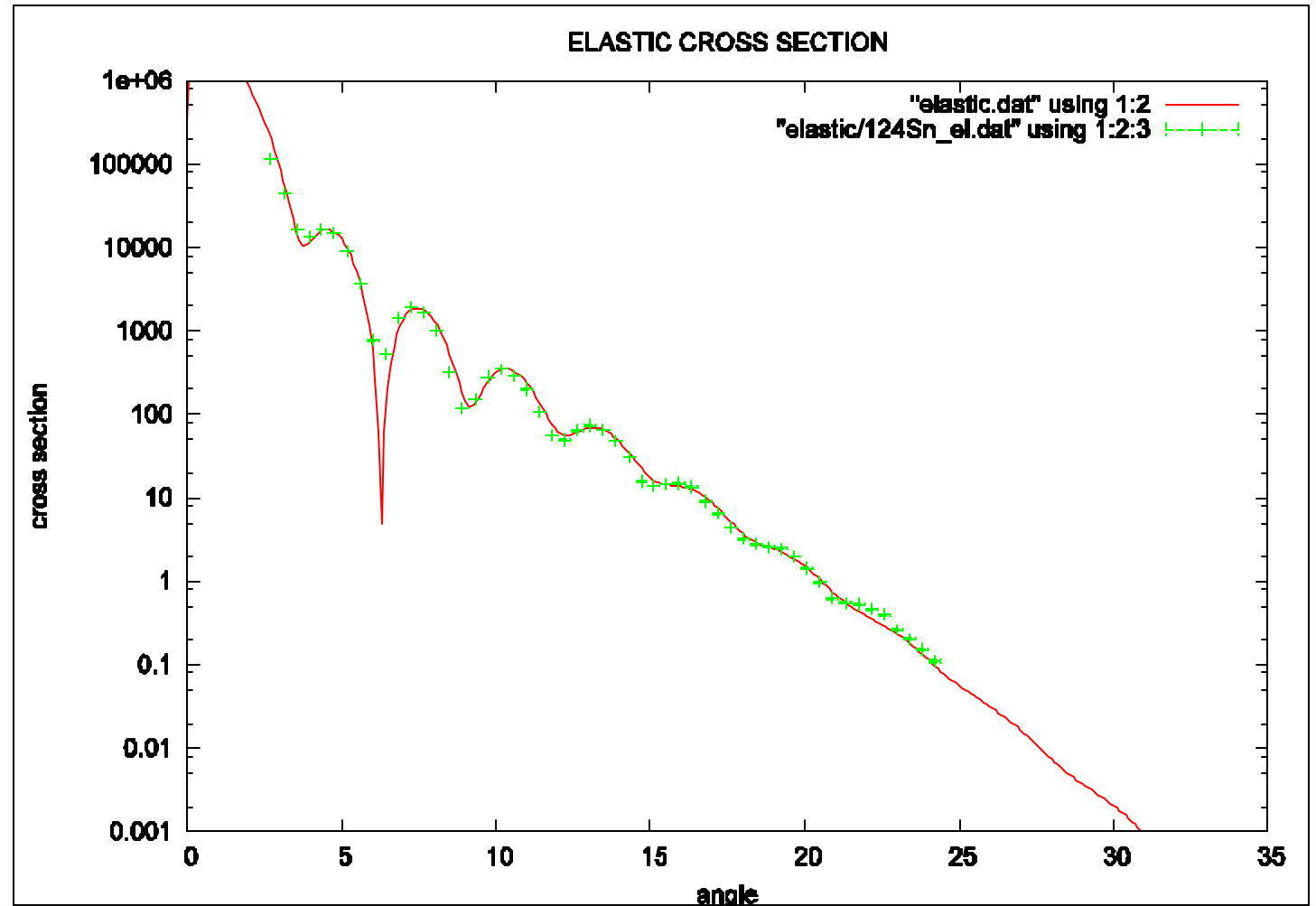


Density-dependent N-a interaction

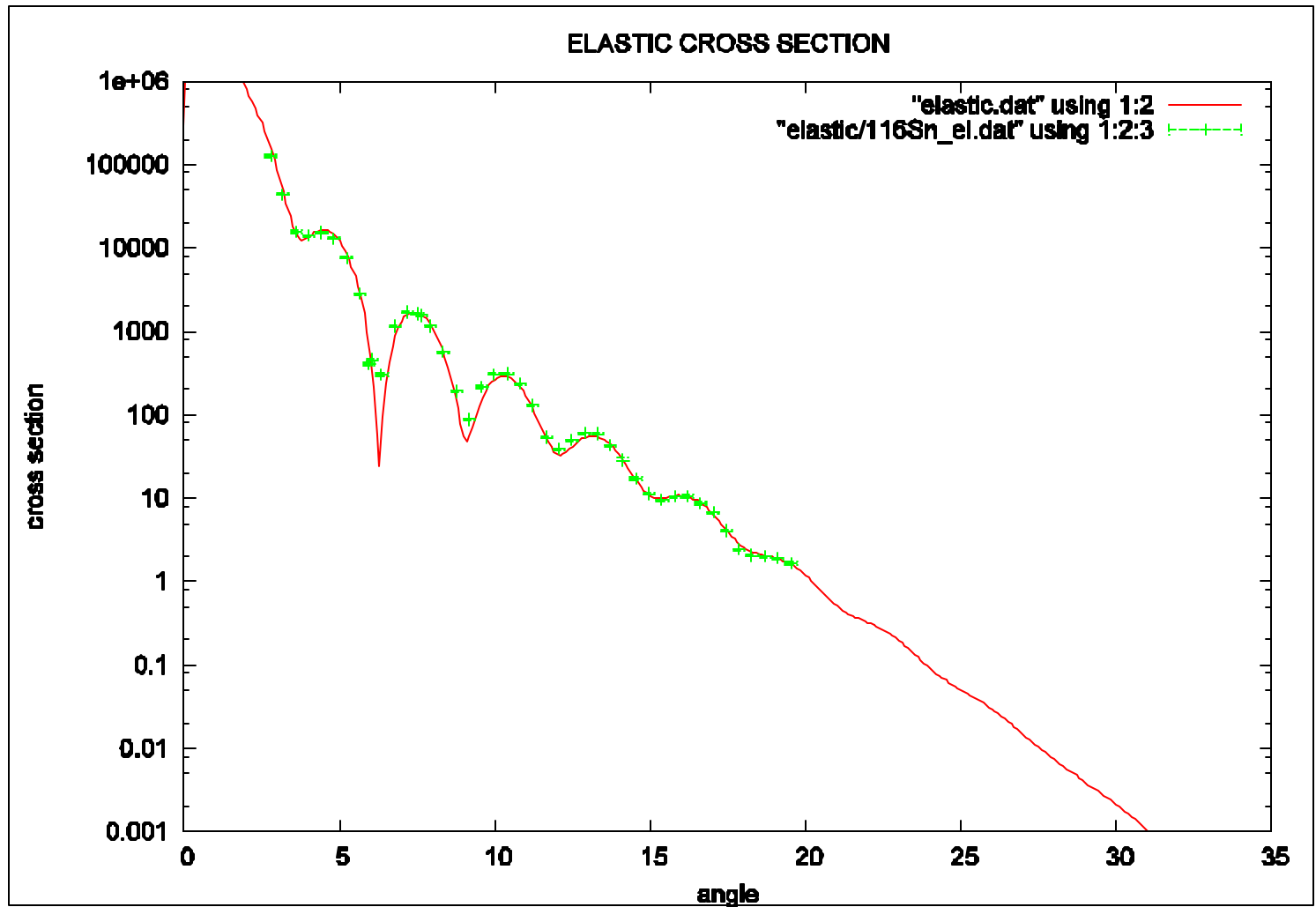
$$V(|\vec{r} - \vec{r}'|, \mathbf{r}_0(r')) = -V(1 + \mathbf{b}_V \mathbf{r}_0(r')^{2/3}) \exp(-|r - r'|^2 / a_V) - iW(1 + \mathbf{b}_W \mathbf{r}_0(r')^{2/3}) \exp(-|r - r'|^2 / a_W)$$

Interaction parameters were obtained for ^{124}Sn by fitting the elastic scattering.

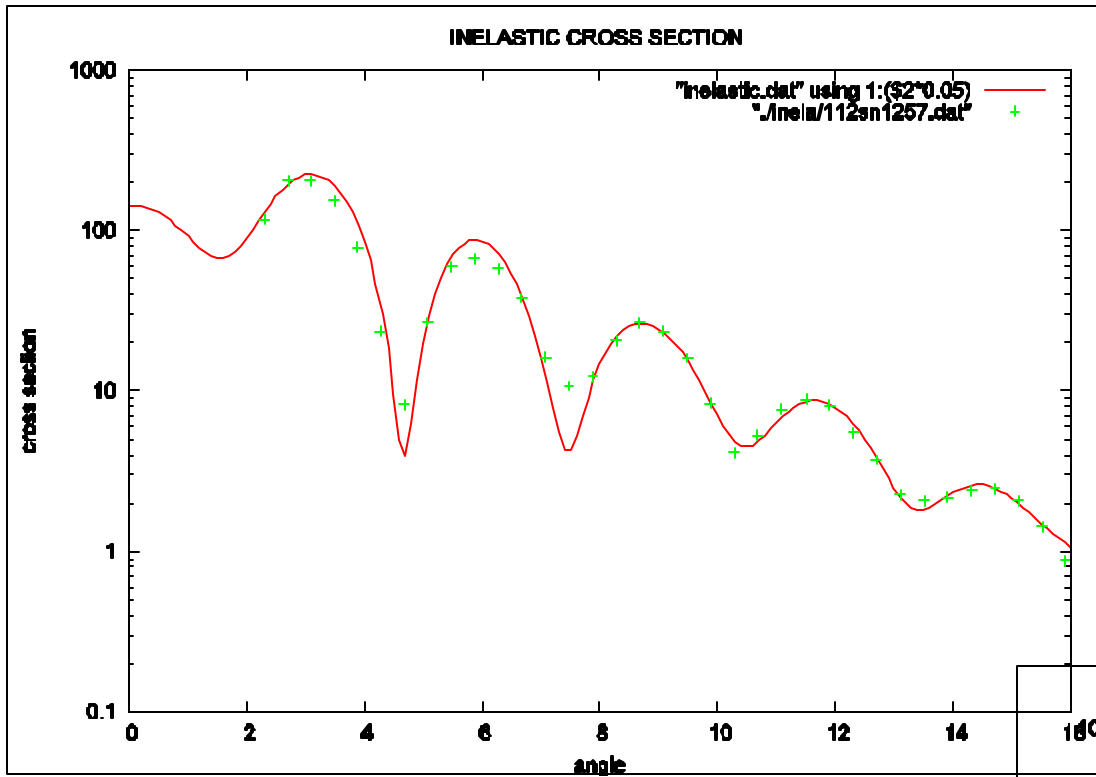
The angular distributions were calculated with the DWBA code "ECIS95"



	V (MeV)	a _v (fm ²)	β _v (fm ²)	W (MeV)	a _w (fm ²)	β _w (fm ²)
^{124}Sn	29.39	4.09	1.9	14.99	4.31	1.9

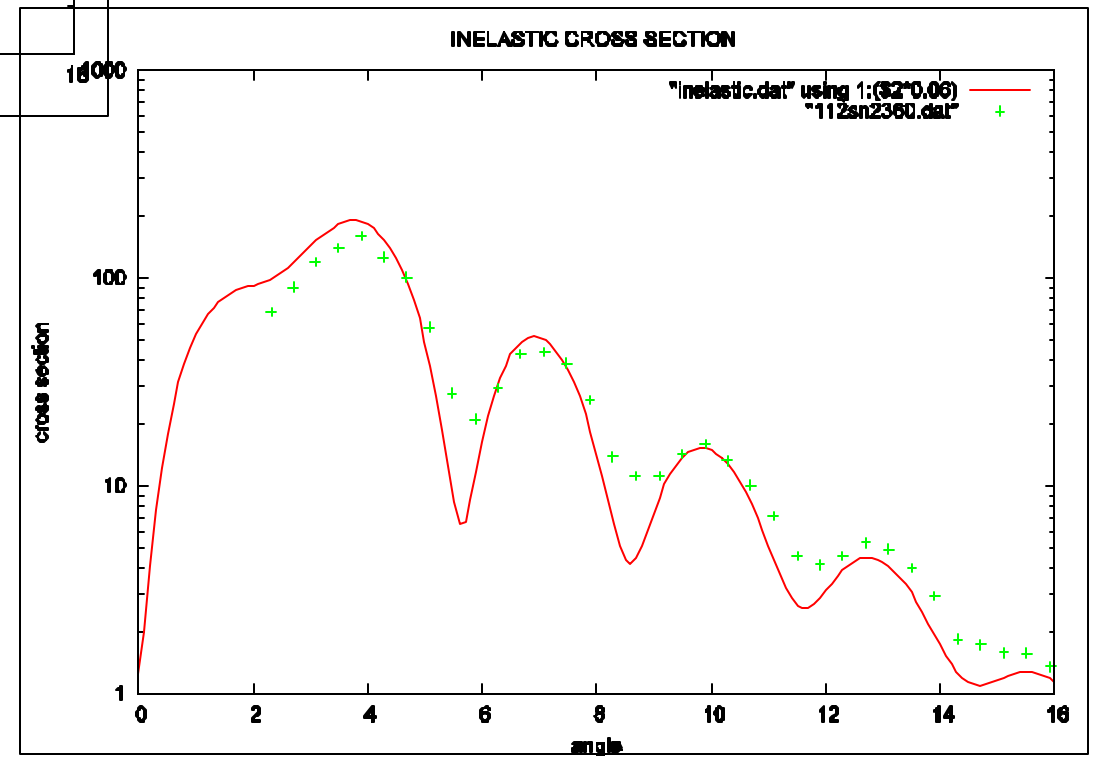


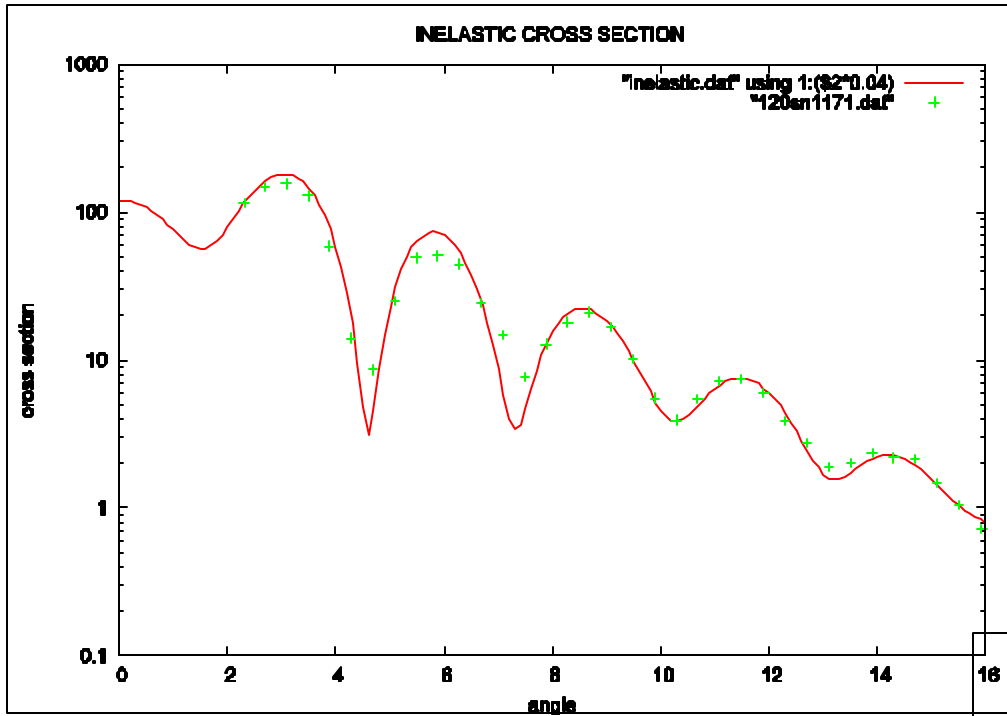
	V (Me V)	a _v (fm ²)	β _v (fm ²)	W (MeV)	a _w (fm ²)	β _w (fm ²)
¹ 16S n	31.19	3.82	1.9	16.47	4.15	1.9
	29.70	3.82	1.9	14.82	4.15	1.9



$^{112}\text{Sn } 2^+ (1.257\text{MeV})$

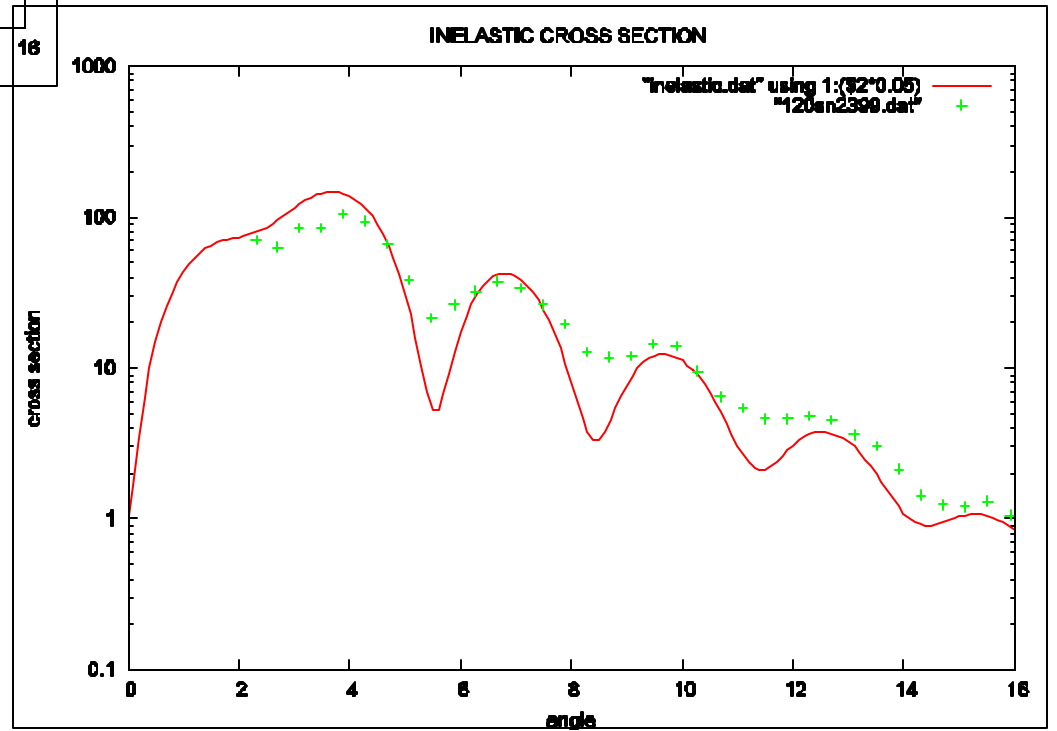
$^{112}\text{Sn } 3^- (2.360\text{MeV})$

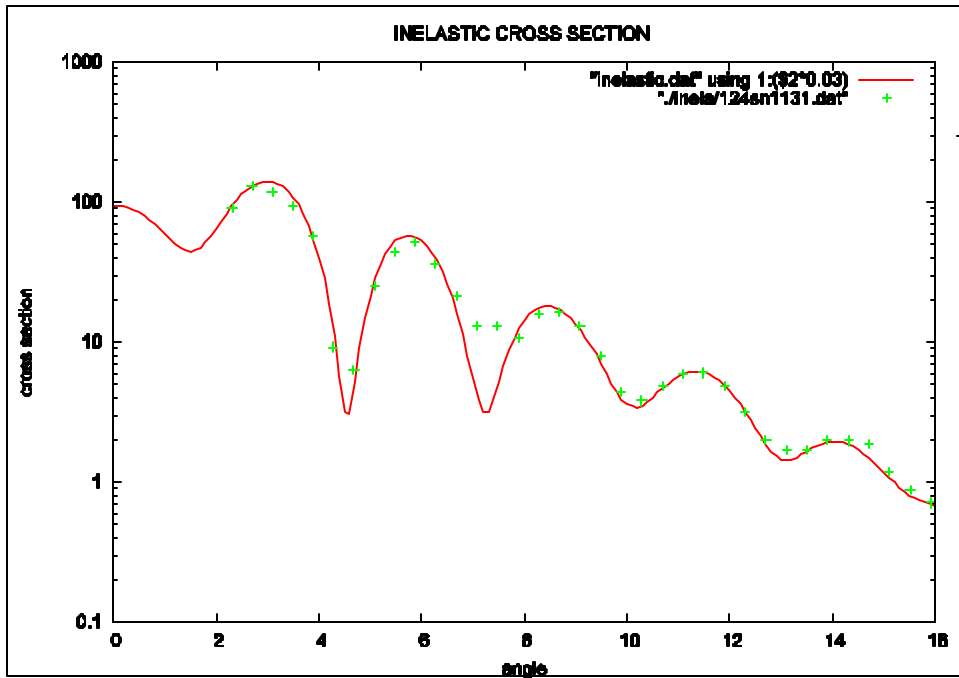




$^{120}\text{Sn } 2^+ (1.171\text{MeV})$

$^{120}\text{Sn } 3^- (2.399\text{MeV})$





^{124}Sn 2^+ (1.131 MeV)

^{124}Sn 3^- (2.614 MeV)

^{124}Sn 4^+ (3.158 MeV)

