

Probing the EOS of Neutron-Rich Matter with Heavy-Ion Reactions

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• **EOS of neutron-rich matter**

• **Terrestrial lab probes**

Two examples: Isospin diffusion as a probe at low densities

Pion production as a probe at high densities

• **Summary**

EOS of Asymmetric Nuclear Mater

At density ρ and neutron-excess $\delta \equiv (\rho_n - \rho_p)/\rho$

$$e(\rho, \delta) \equiv E/A = T_F(\rho, \delta) + V_0(\rho) + \delta^2 V_2(\rho)$$

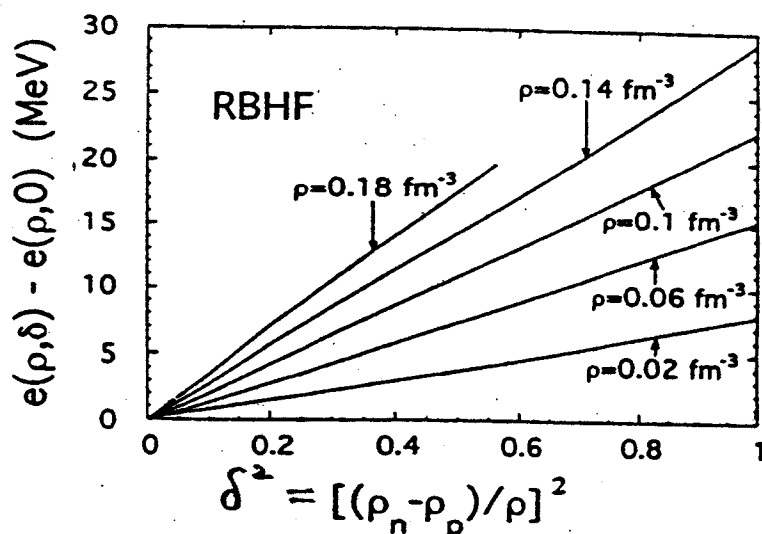
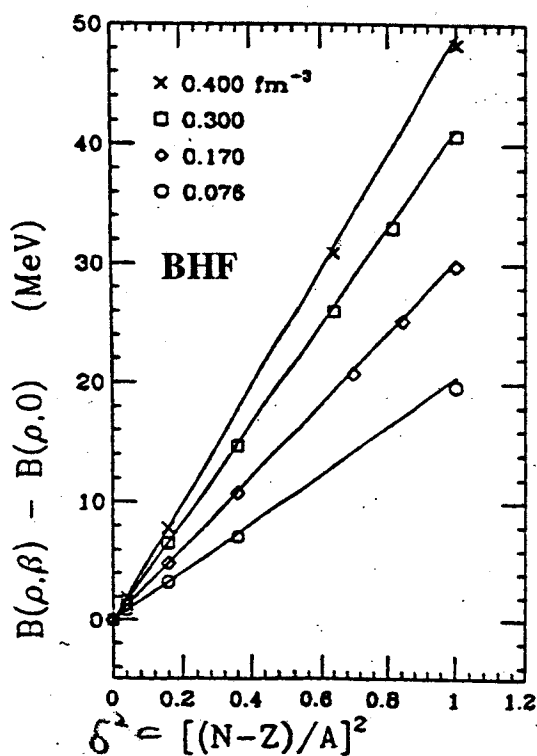
Empirical parabolic law:

$$e(\rho, \delta) = e(\rho, 0) + e_{sym}(\rho)\delta^2 + \mathcal{O}(\delta^4)$$

$$\begin{aligned} \Rightarrow e_{sym}(\rho) &= e(\rho, \text{pure neutron matter}) - e(\rho, \text{symmetric nuclear matter}) \\ &\equiv \frac{1}{2} \left(\frac{\partial^2 e}{\partial \delta^2} \right)_\rho \end{aligned}$$

Theoretical evidence: essentially all many-body calculations

Two examples:



H. Huber, F. Weber and M.K. Weigel, PLB 317 (1993) 485.

I. Bombaci and U. Lombardo, PRC 44 (1991) 1892.

The latest microscopic many-body theory predictions

On the nuclear symmetry energy and the neutron skin in neutron rich nuclei

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arXiv:nucl-th/0312031 v1 10 Dec 2003

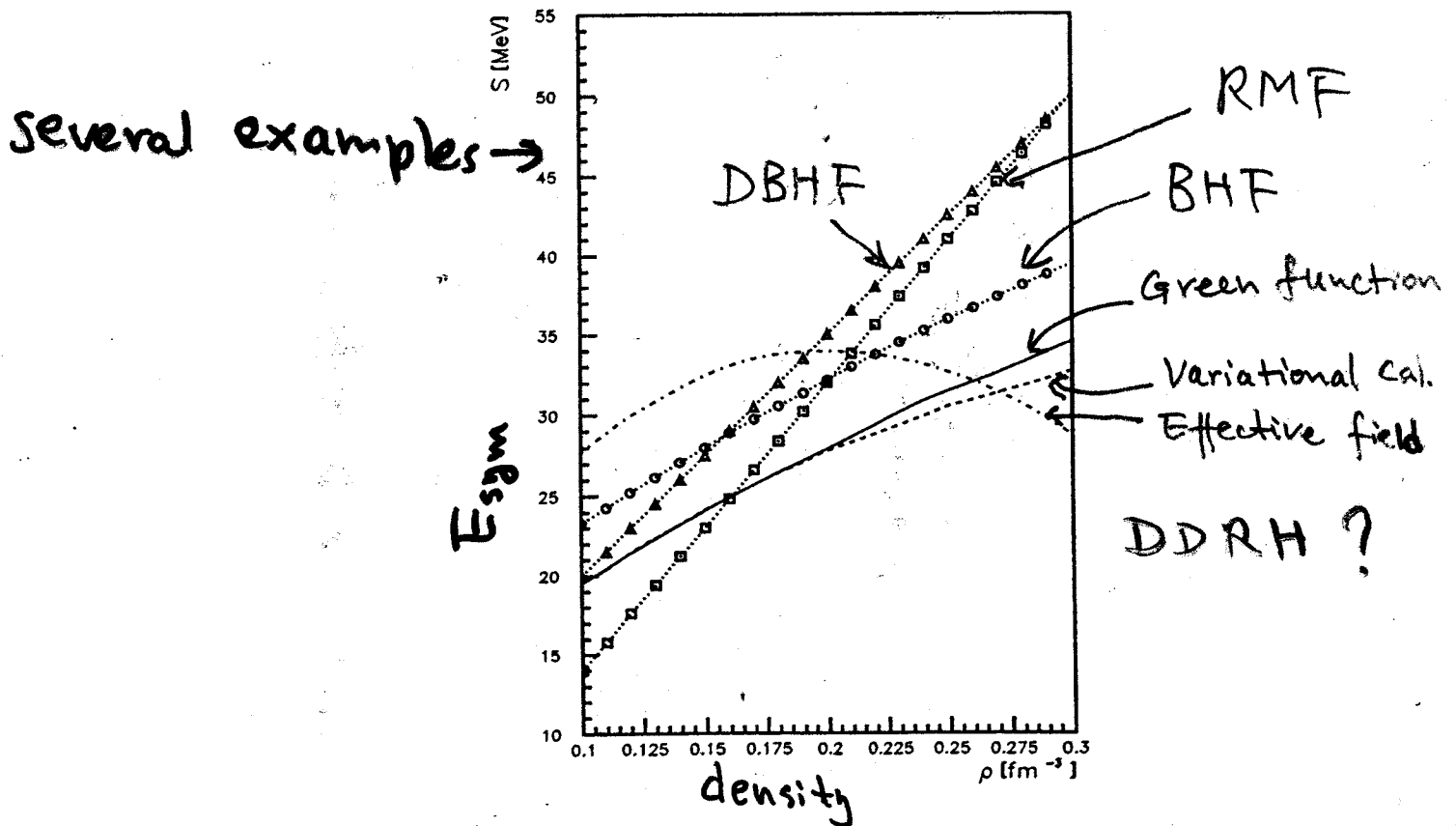
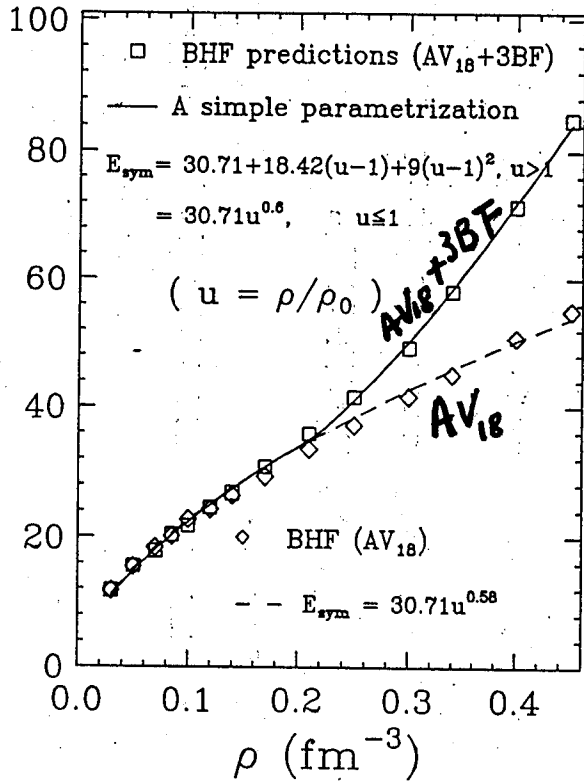
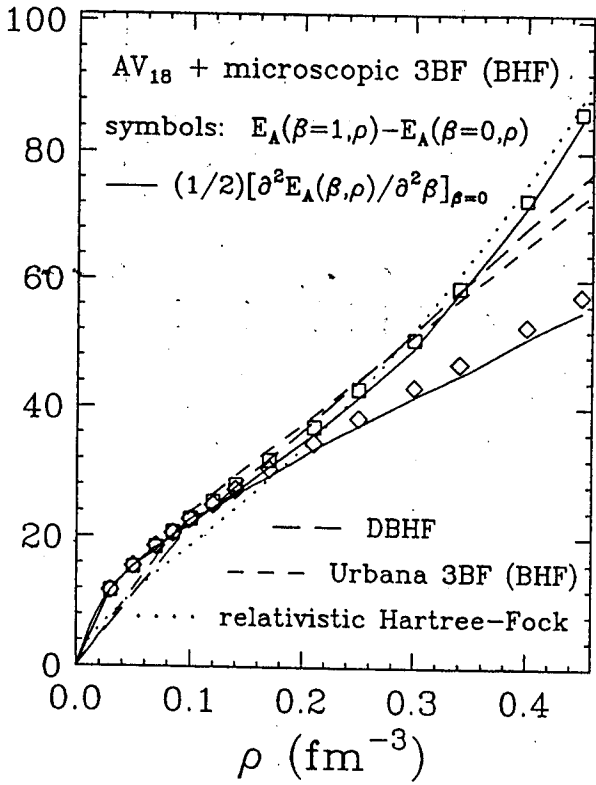


FIG. 3: Overview of several theoretical predictions for the symmetry energy S : Brueckner-Hartree-Fock (continuous choice) with Reid93 potential (circles), self-consistent Green function theory with Reid93 potential (full line), variational calculation from [17] with Argonne Av14 potential (dashed line), Dirac-Brueckner-Hartree-Fock calculation from [25] (triangles), relativistic mean-field model from [28] (squares), effective field theory from [30] (dash-dotted line).

Wei Zuo and Umberto Lombardo
 extended BHF



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arXiv:nucl-th/0309012 v1 5 Sep 2003

phenomenological model predictions

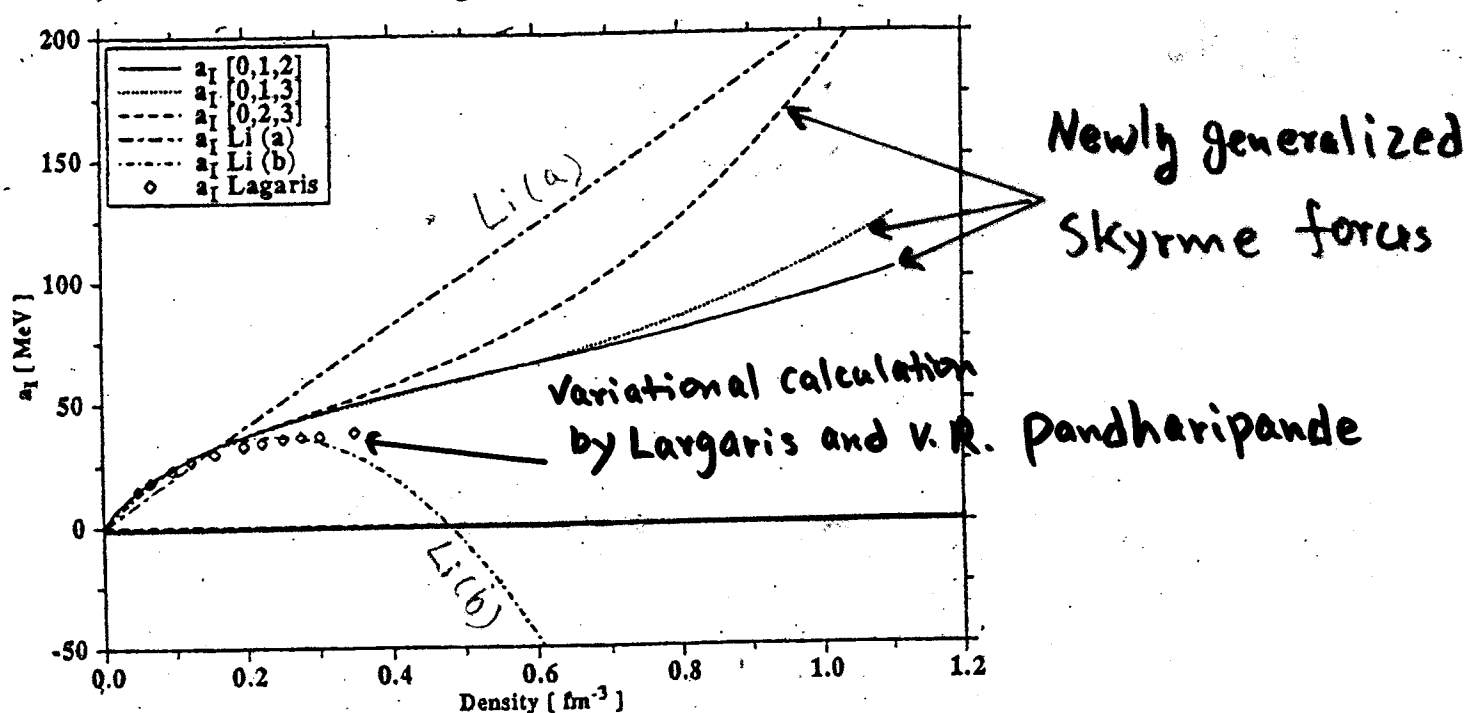


Figure 3. Symmetry energy in infinite matter as a function of the density. Li(a) and (b): phenomenological parametrizations by Bao-An Li [16], see text. Open diamonds: variational calculation of Lagaris et al. [17]. Inset: symmetry energy in the vicinity of equilibrium density of symmetric nuclear matter.

Promising Probes of the $E_{\text{sym}}(\rho)$ in Nuclear Reactions

(an incomplete list !)

(a) At low densities

- Sizes of n-skins of unstable nuclei from total reaction xsections
- Parity violating electron scattering studies of the n-skin in ^{208}Pb
- Multiplicity and spectrum of pre-equilibrium neutrons/protons
- Isospin fractionation and isoscaling in multifragmentation
- Isospin diffusion
- Proton differential elliptic flow at high transverse momenta
- Isospin dependence of transverse flow and balance energy
- Neutron-proton correlation functions at low relative momenta
- $t/{}^3\text{He}$ ratio
- Deflection function

(b) Towards high densities

- π^-/π^+ ratio in heavy-ion collisions induced by high-energy radioactive beams
- Neutron-proton differential flow in heavy-ion collisions
- Precursor of *isospin separation instability* in the excitation function of nuclear collective flow
- Asymmetric nuclear matter induced unique ρ^0 - ω mixing and its effects on dilepton and photon production

*** **The most sensitive observables require simultaneous measurement of neutrons and protons, i.e., neutron-proton differential quantities.**

Status on the density-, momentum- and isospin-dependent single nucleon potential $U_{\tau}(k,\rho,\delta)$

(1) Microscopic many-body theories

A lot of work, $U_{\tau}(k,\rho,\delta)$ is normally plotted as a function of density at few ρ and δ in all published papers.

No analytical expression for $U_{\tau}(k,\rho,\delta)$ is given in any of the publications, the predictions are thus hard to be used in heavy-ion studies.

(2) Two phenomenological expressions for $U_{\tau}(k,\rho,\delta)$ exist. Both are derived using the HF with Gogny forces as a guide

Both contains parameters to mimic microscopic many-body predictions on the $E_{\text{sym}}(\rho)$ and $M^*_{\tau}(\rho,\delta)$

(A) Ignazio Bombaci

Two parameters (x_0 and x_3) to mimic the $E_{\text{sym}}(\rho)$

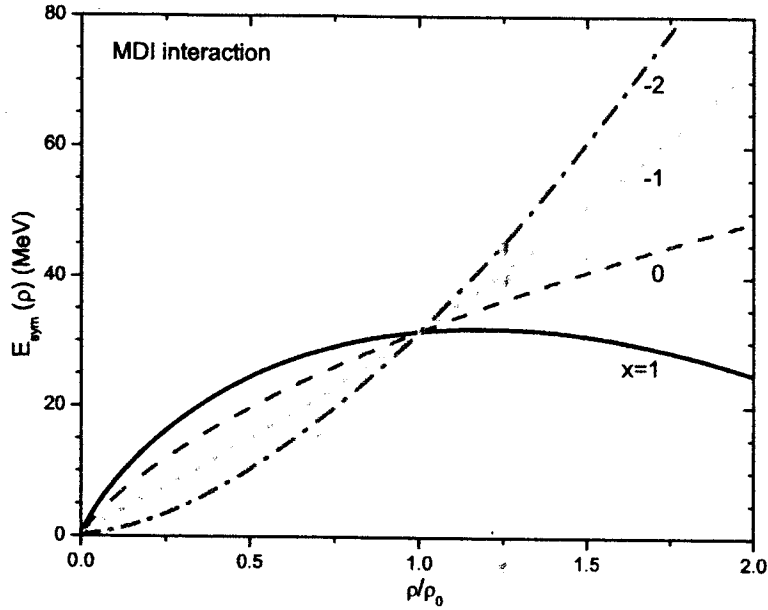
Three parameters (z_1, x_0 and x_3) to mimic the nucleon effective mass $M^*_n(\rho,\delta)$ and $M^*_p(\rho,\delta)$

(B) Subal Das Gupta et al.

Three parameters (x, A_{like} and A_{unlike}) to mimic the $E_{\text{sym}}(\rho)$

$M^*_n(\rho,\delta)$ and $M^*_p(\rho,\delta)$ are fixed, they are independent of the parameters x, A_{like} and A_{unlike} since they are all in the density-dependent parts of the potential.

*** Akira Ono uses one parameter x in the Gogny 2-body potential to vary the $E_{\text{sym}}(\rho)$ in AMD calculations



Momentum and isospin dependent nucleon potential used in the IBUU transport model

$$U(\rho, \delta, \vec{p}, \tau) = A_u \frac{\rho_{\tau'}}{\rho_0} + A_l \frac{\rho_{\tau}}{\rho_0} + B \left(\frac{\rho}{\rho_0} \right)^{\sigma} (1 - x \delta^2) - 8 \tau x \frac{B}{\sigma + 1} \frac{\rho^{\sigma-1}}{\rho_0^{\sigma}} \delta \rho_{\tau'}$$

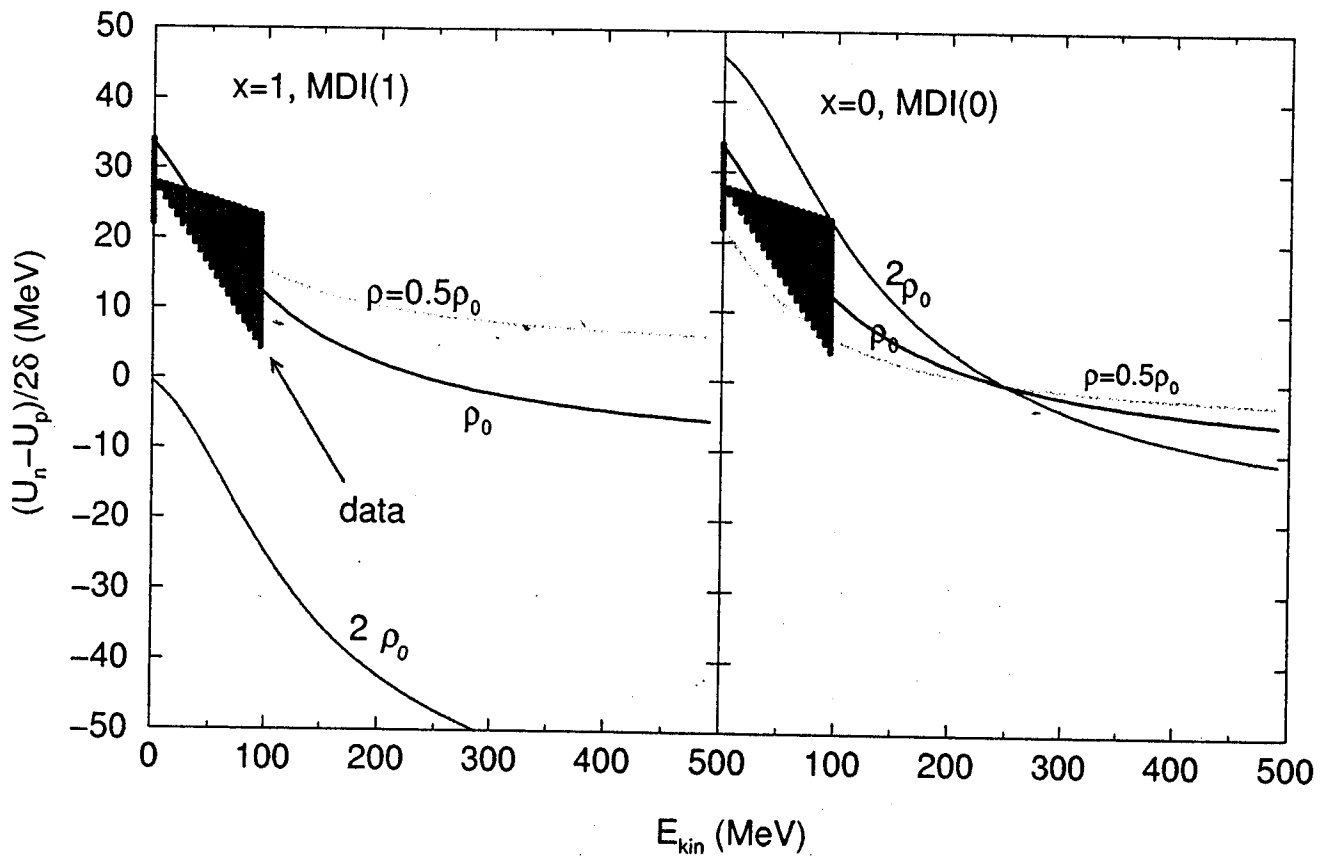
$$+ \frac{2C_{\tau, \tau}}{\rho_0} \int d^3 p' \frac{f_{\tau}(\vec{r}, \vec{p}')}{1 + (\vec{p} - \vec{p}')^2 / \Lambda^2} + \frac{2C_{\tau, \tau'}}{\rho_0} \int d^3 p' \frac{f_{\tau'}(\vec{r}, \vec{p}')}{1 + (\vec{p} - \vec{p}')^2 / \Lambda^2}$$

$$C_{unlike} = -103.4 \text{ MeV}, C_{like} = -11.7 \text{ MeV}, K_0 = 211 \text{ MeV}, \tau = \pm 1/2 \text{ for n/p}$$

x is used to mimic different density-dependent symmetry energy predicted by various many-body theories, or to vary K_{asy} .

$$A_l(x) = -120.57 + x \frac{2B}{\sigma + 1}, \quad A_u(x) = -95.98 - x \frac{2B}{\sigma + 1}. \quad (1)$$

- The formalism within Hartree-Fock approach using a Gogny force
C.B. Das, S. Das Gupta, C. Gale and B.A. Li, PRC 67, 034611 (2003).
- Applications in heavy-ion collisions induced by neutron-rich nuclei:
 - (1) B. A. Li, C.B. Das, Subal Das Gupta and C. Gale, PRC 69, 034614 (2004).
 - (2) B. A. Li, C.B. Das, Subal Das Gupta and C. Gale, NPA 735, 563 (2004).



Evidence of the momentum-dependence of symmetry potential from the difference of optical potentials for neutron and proton scatterings

(1) Neutron and proton scatterings on the same nucleus at the same energy

$$U_n - U_p = 2\delta U_{\text{sym}} + U_{\text{Coulomb correction}}$$

(2) Neutron or proton scatterings on a sequence of isotopes

$$U_p(\text{isotope 1}) - U_p(\text{isotope 2}) = 2(\delta_1 - \delta_2) U_{\text{sym}}$$

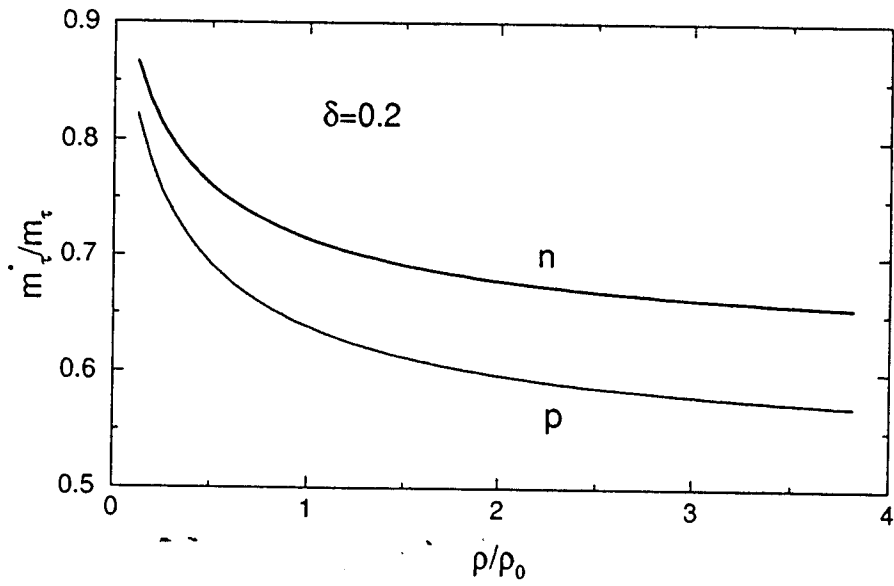
The finding: $U_{\text{sym}} = V_1 - \epsilon_R E$, $V_1 \approx 28 \pm 6 \text{ MeV}$, $\epsilon_R \approx 0.1 - 0.2$

P.E. Hodgson, the Nucleon Optical Model, World Scientific, 1994

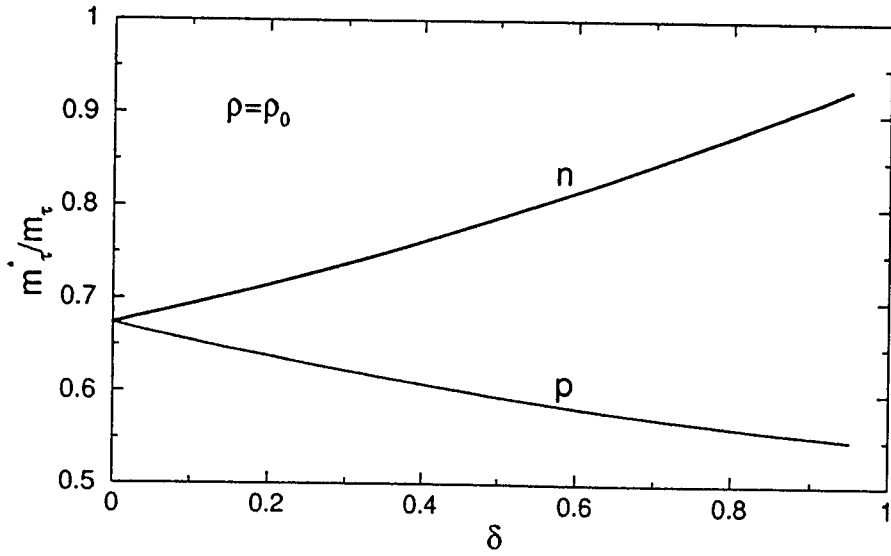
2.4.4. Energy dependence of isospin potential. Some global analyses of neutron elastic scattering have been made that included a linear energy dependence of the asymmetry term. Values of the coefficients ϵ_R and ϵ_I in the expressions $(V_1 - \epsilon_R E)$ and $(W_1 - \epsilon_I E)$ are given in table 12.

Table 12. Values of coefficient of linear energy dependence of isospin potential.

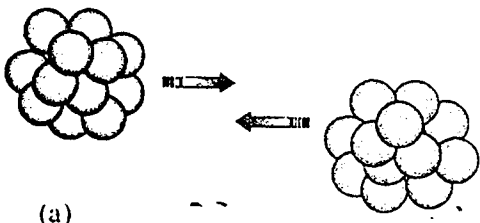
Reaction	E	Nucleus	ϵ_R	ϵ_I	Reference
		Theoretical analyses	0.1	—	Rook (1973)
			0.17	—	Dabrowski and Haensel (1974)
			0.1	—	Jeukenne <i>et al</i> (1977b)
(n, n)	7-26	$^{40}\text{Ca}, ^{208}\text{Pb}$	0.19	—	Rapaport <i>et al</i> (1979a, b)
(p, n)	25-45	$^{48}\text{Ca}-^{208}\text{Pb}$	0.18	—	Patterson <i>et al</i> (1976)
(n, n), (p, p)	10-50	^{208}Pb	0.183 ± 0.008	—	De Vito <i>et al</i> (1981)
(n, n), (p, p)	10-40	^{208}Pb	—	0.178 ± 0.052	De Leo and Micheletti (1981)
(p, p)	100	$^{58}\text{Ni}-^{208}\text{Pb}$	0.12 ± 0.06	—	Kwiatkowski and Wall (1978)



$$\frac{m_1^*}{m_\tau} = \left[1 + \frac{m_\tau}{p} \frac{dU}{dp} \right]^{-1} \frac{p}{p_\tau}$$

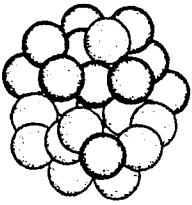


Isospin diffusion



Particle Flux:

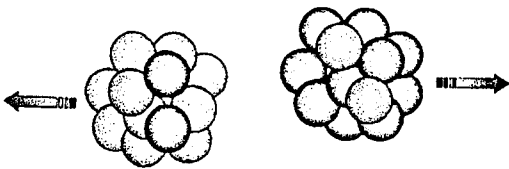
$$\Gamma_i = n_i (\underline{v}_i - \underline{v}),$$



Isospin Flow:

$$\Gamma_I = \Gamma_n - \Gamma_p = -n \underbrace{D_I}_{\text{Isospin diffusion coef } D_I} \frac{\partial \delta}{\partial r}.$$

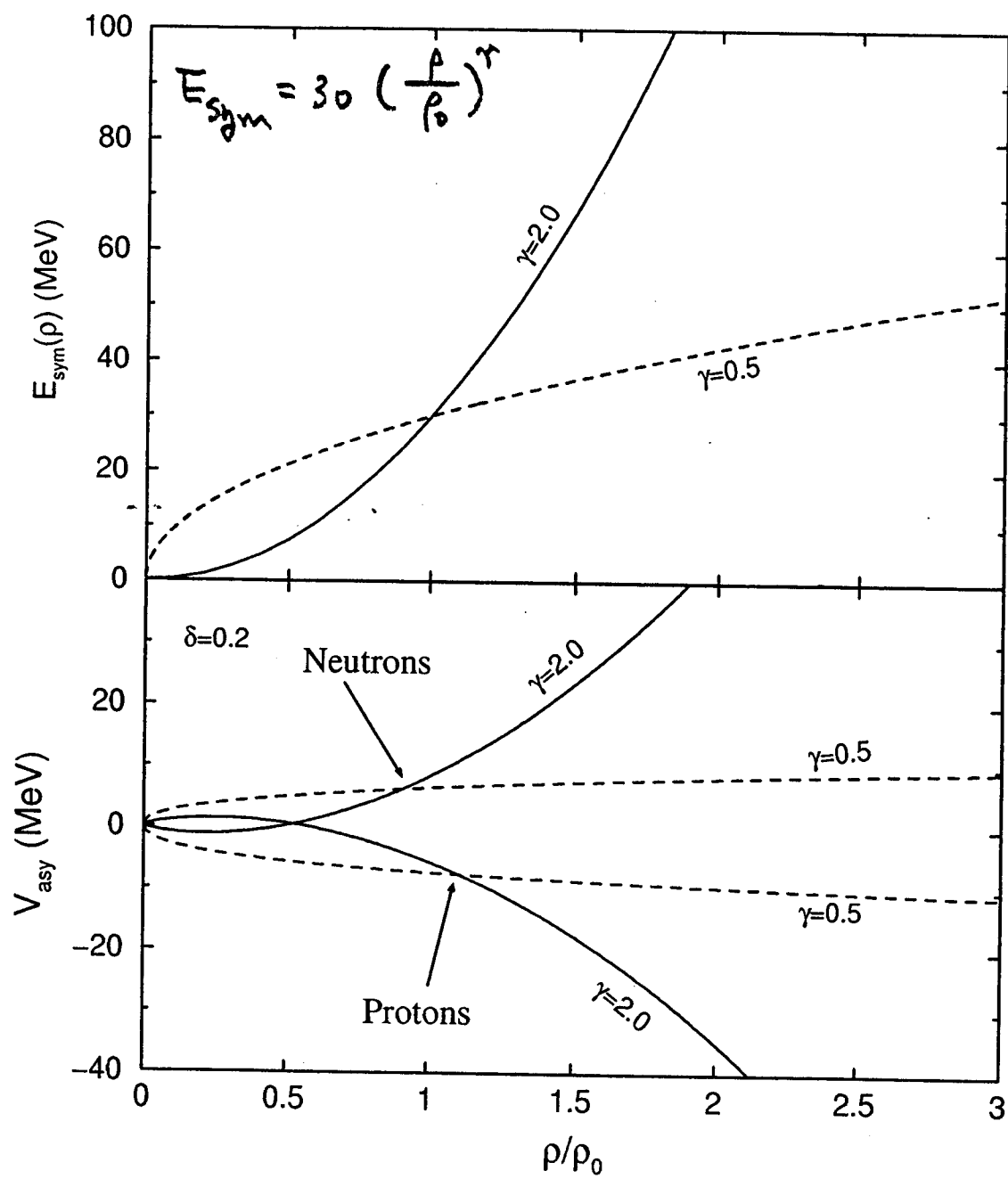
Isospin diffusion coef D_I



Diffusion

L. Shi

P. Danielewicz

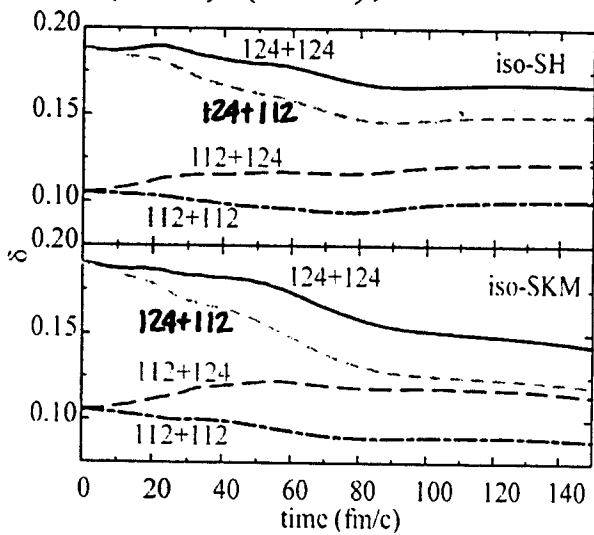


R_i from simulations by Shi and Danielewicz

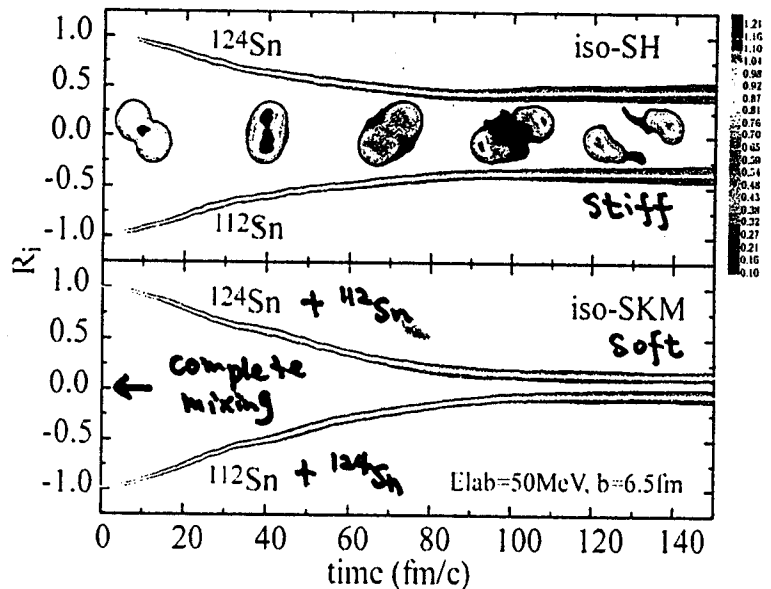
Projectile isospin asymmetry

δ from simulation

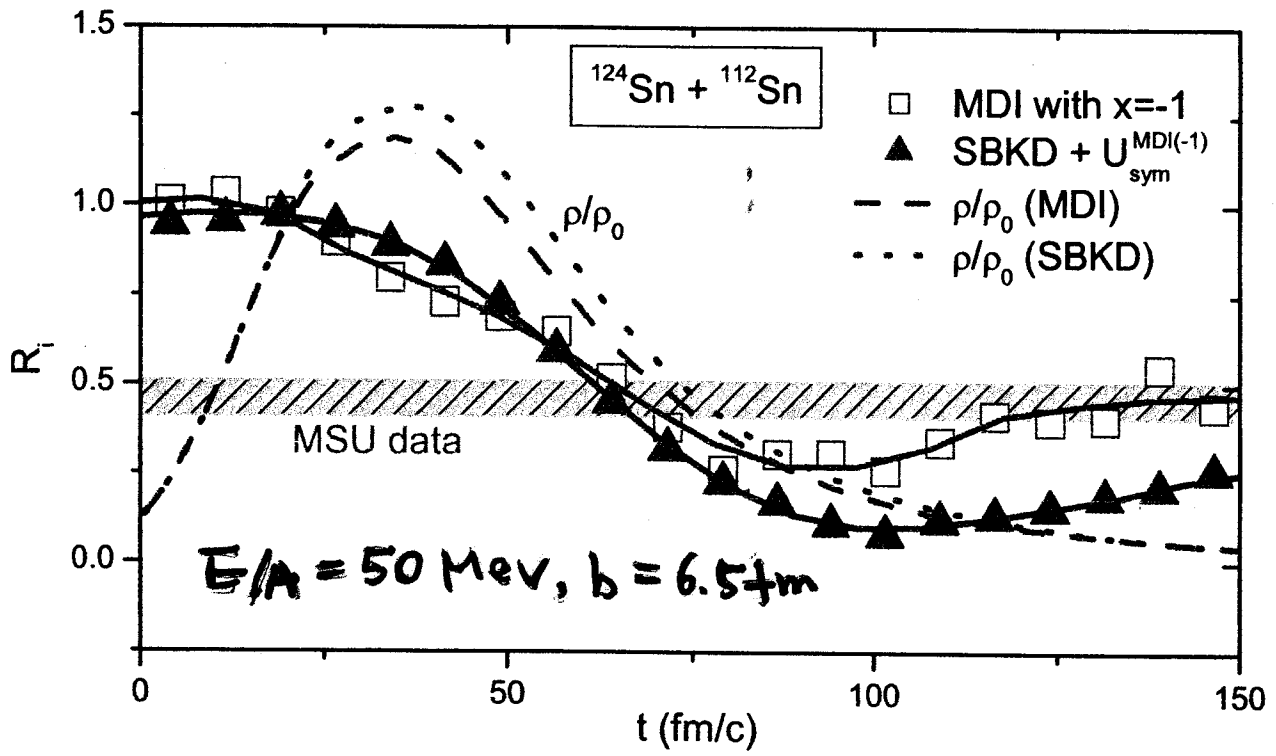
$$\delta = (N-Z)/(N+Z),$$



$$R_i = \frac{(2\delta_i - \delta_{^{124}\text{Sn}+^{124}\text{Sn}} - \delta_{^{112}\text{Sn}+^{112}\text{Sn}})}{(\delta_{^{124}\text{Sn}+^{124}\text{Sn}} - \delta_{^{112}\text{Sn}+^{112}\text{Sn}})}$$

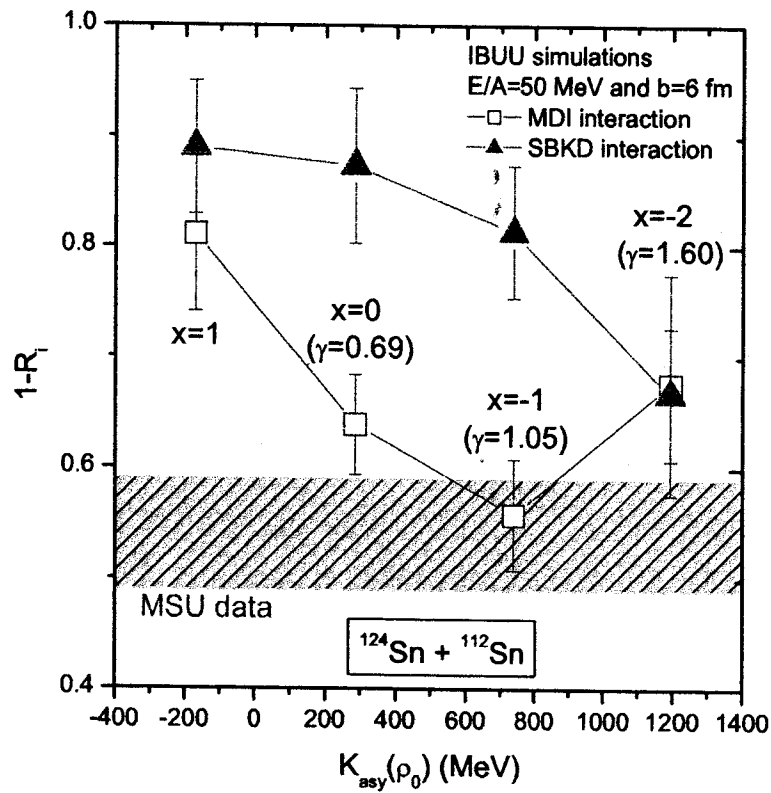


- R_i is a stable signal
- Non-diffusion effects: cancelled out
- $R_i \sim$ IEOS \rightarrow Diffusion effect



- Momentum-dependence is important for isospin diffusion
- For the Sn+Sn reactions at 50 MeV/A, the isospin diffusion happens in the density range of $0.3 < \rho/\rho_0 < 1.2$

Strength of Isospin diffusion



$$K(\text{asymmetric matter}) = K_0(\text{symmetric matter}) + K_{asy} \cdot \delta^2 \quad (1)$$

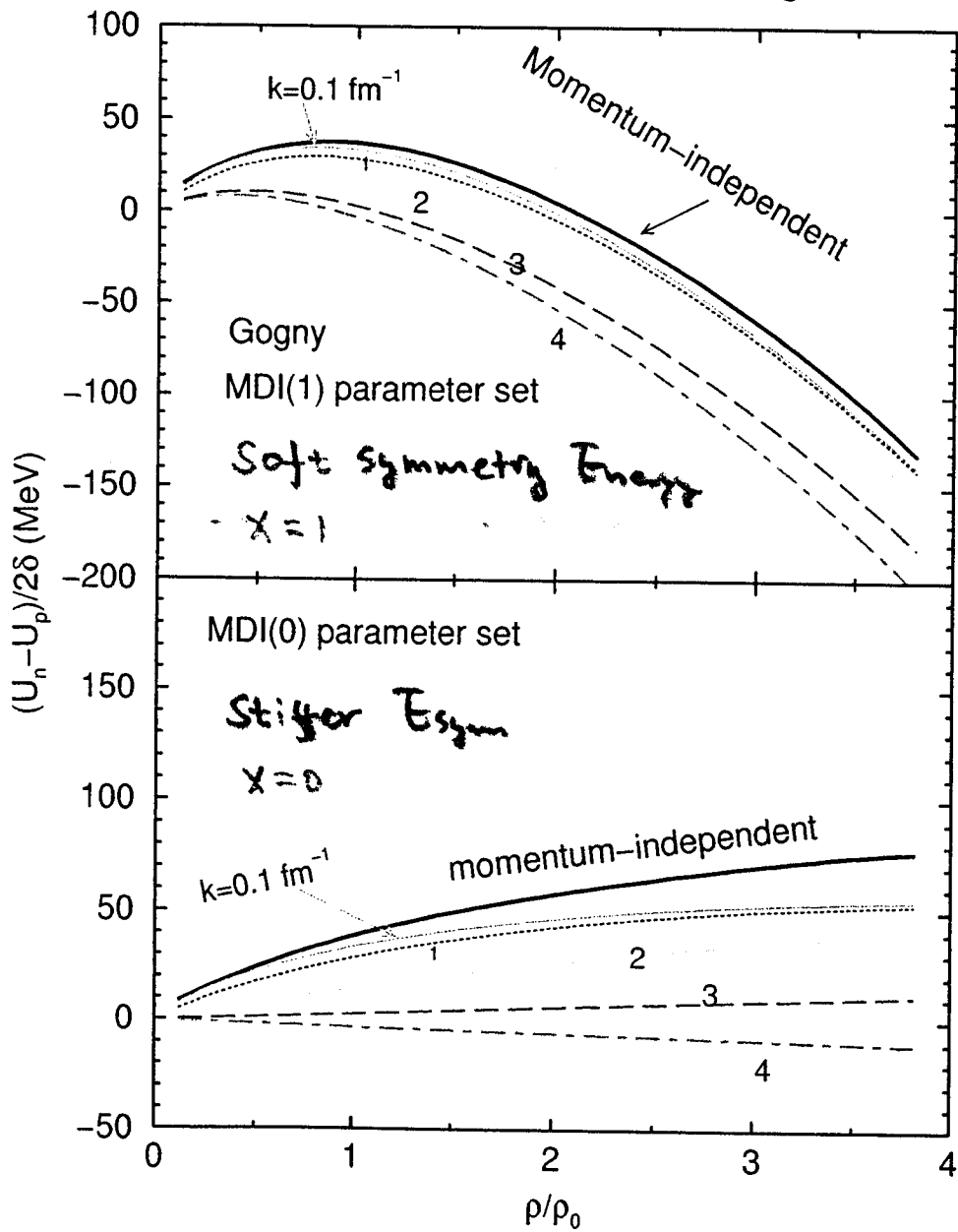
$$K_{asy}(\rho) = 18\rho \frac{dE_{sym}}{d\rho} + 9\rho^2 \frac{d^2E_{sym}}{d\rho^2} \quad (2)$$

The data favors $E_{sym} = 31.6(\rho/\rho_0)^{1.05}$ and $K_{asy}(\rho_0) \approx 738 \pm 150$ MeV.

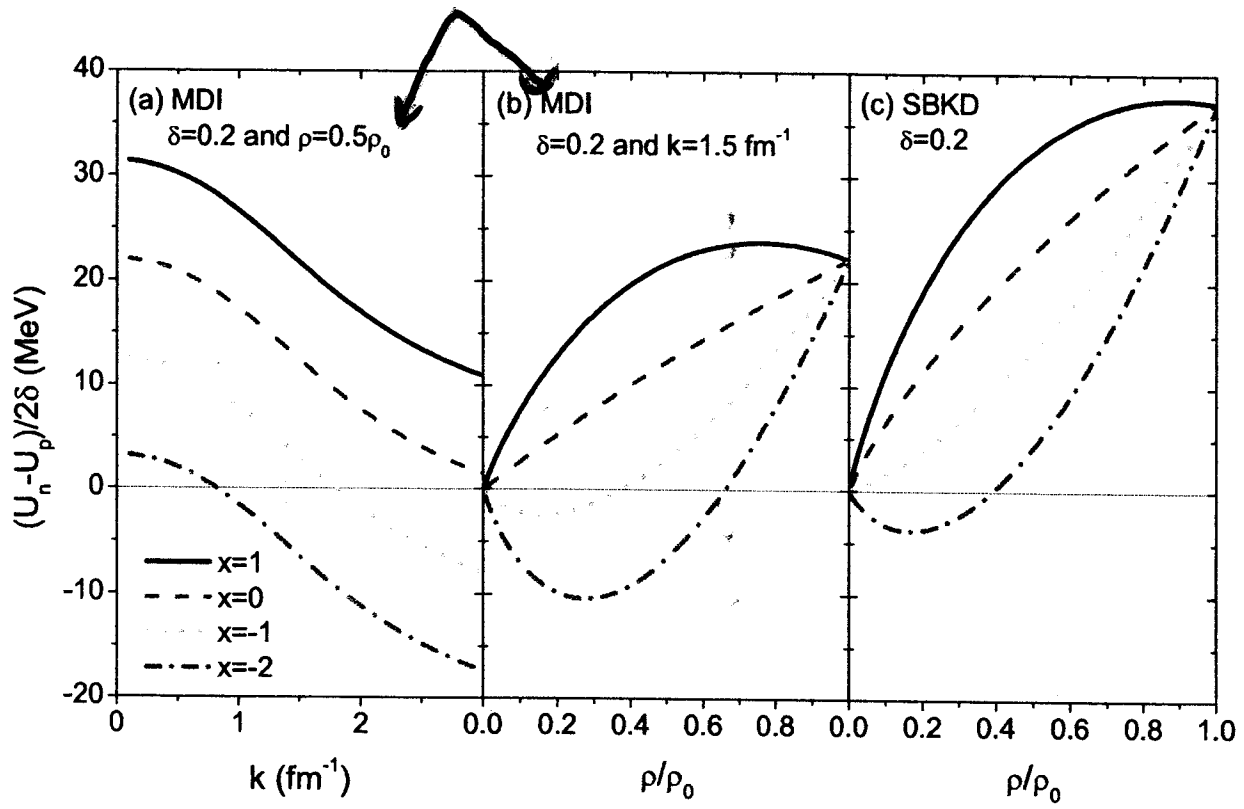
Giant Monopole resonance studies give:

-566 ± 1350 MeV $< K_{asy}(\rho_0) < 139 \pm 1617$ MeV by Shlomo and Youngblood in PRC 47, 529 (1993).

Strength of the symmetry potential

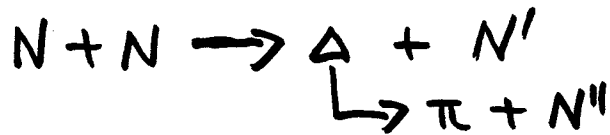


momentum-dependent



Strength of the symmetry potential

Pion probe of N/Z of high density participants



$\Delta(1232)$ resonance model:

	π^+	π^0	π^-
nn	0	1	5
pp	5	1	0
np(pn)	1	4	1

Pion asymmetry: before re-absorption and re-scattering

$$\frac{\pi^-}{\pi^+} = \frac{5N^2 + NZ}{5Z^2 + NZ} \approx \left(\frac{N}{Z}\right)^2$$

Thermal model: (G.F. Bertsch, Nature 283 (1980) 281.)

$$\frac{\pi^-}{\pi^+} \propto \exp[(\mu_n - \mu_p)/kT] \text{ (the same factor for n/p ratio !)}$$

$$\mu_n - \mu_p = 2V_{sym}(\rho)\delta - V_{Coul} + T \left\{ \ln \frac{\rho_n}{\rho_p} + \sum_m \frac{m+1}{m} b_m \left(\frac{1}{2}\lambda_T^3\right)^m (\rho_n^m - \rho_p^m) \right\}$$

Henry Jazaman
Aram Mekjian
PRC (1984)

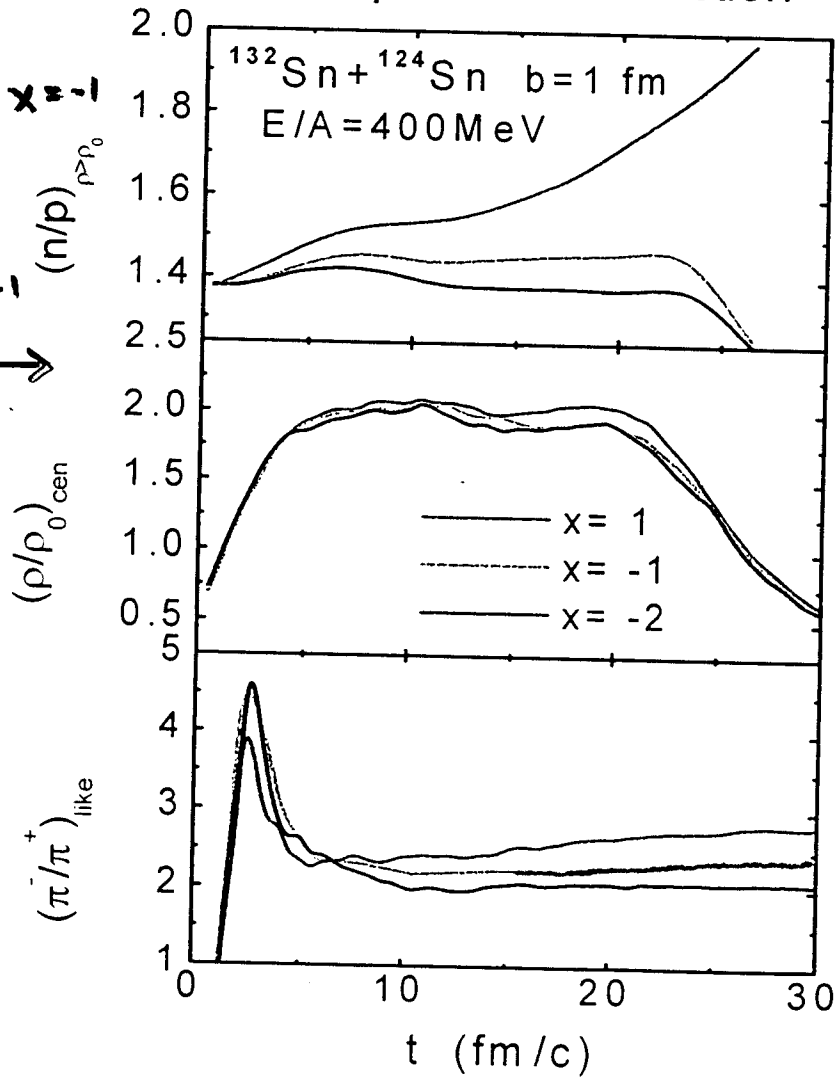
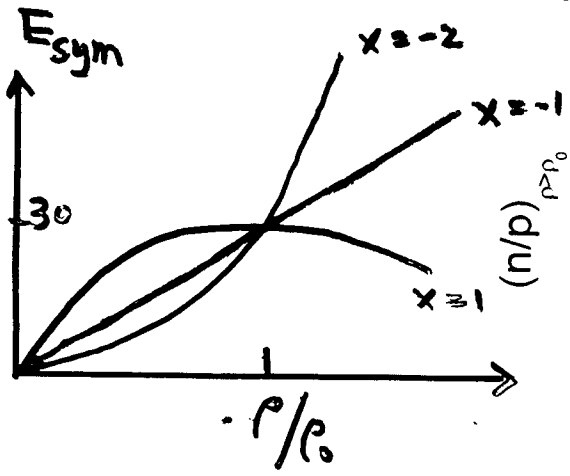
λ_T is the thermal wavelength and b_m are constants.

$$\rightarrow \frac{n}{p}, \frac{\pi^-}{\pi^+} \propto \exp[(2V_{sym}(\rho)\delta - V_{Coul})/kT]$$

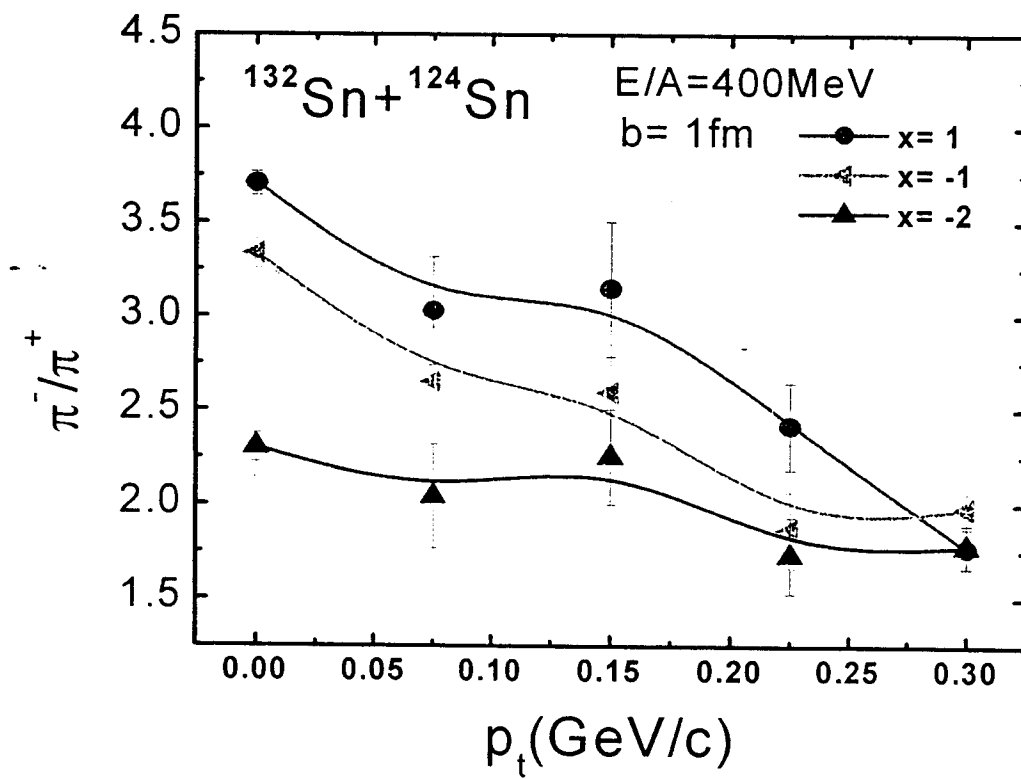
Pion-like ratio during the reaction:

$$\left(\frac{\pi^-}{\pi^+}\right)_{like} = \frac{\pi^- + \Delta^- + \frac{1}{3}\Delta^0 + \frac{2}{3}N^{*0}}{\pi^+ + \Delta^{++} + \frac{1}{3}\Delta^+ + \frac{2}{3}N^{*+}} \xrightarrow{t \rightarrow \infty} \frac{\pi^-}{\pi^+}$$

Pion production with momentum-dependent interaction



π^- / π^+ ratio at low transverse momentum



Summary

- $E_{\text{sym}}(\rho)$ is important for both astrophysics and heavy-ion reactions
- Isospin-diffusion data from NSCL/MSU favors
 $E_{\text{sym}}(\rho) = 31.6 (\rho/\rho_0)^{1.05}$
 $K_{\text{sym}}(\rho_0) = 738 \pm 150 \text{ MeV}$ for $\rho/\rho_0 < 1.2$
- π^-/π^+ ratio at low transverse momentum in central heavy-ion collisions is promising for determining the symmetry energy at higher densities.