

Philadelphia, 25.8.2004

Compact binary mergers:

the influence of the equation of state

S. Rosswog



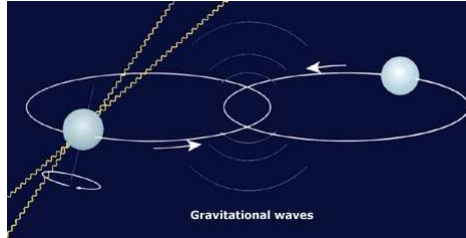
further information:

<http://www.faculty.iu-bremen.de/srosswog>

I.1 Why compact binary systems?

- Physics of Nuclear Matter:
 - * high T \Rightarrow heavy ion collisions
 - * low T \Rightarrow neutron stars
- Nucleosynthesis:
 - * formation of neutron-rich, rapid neutron capture elements
- Gravitational Waves:
 - * large rates: $\sim 10^{-4}$ (year galaxy) $^{-1}$
 - * large, predictable (initial phase) GW-amplitudes
 \Rightarrow good detection prospects
- Gamma-ray Bursts:
 - * “NS-NS & NS-BH mergers
 \Rightarrow BH + massive accretion disk = GRB-engine”

I.2 Observed Neutron Star Binaries



- so far: **six observed systems**
- most precise observation: **PSR1913+16** (Hulse & Taylor 1974)

* **discovered in 1974** by R.Hulse and J.Taylor → NP 1993

* **masses:** 1.442 and $1.386 M_{\odot} \pm 0.0002 M_{\odot}$

* **orbital period:** $\tau_{orb} = 7.752 \text{ h}$ ($v \sim 10^{-3}c$)

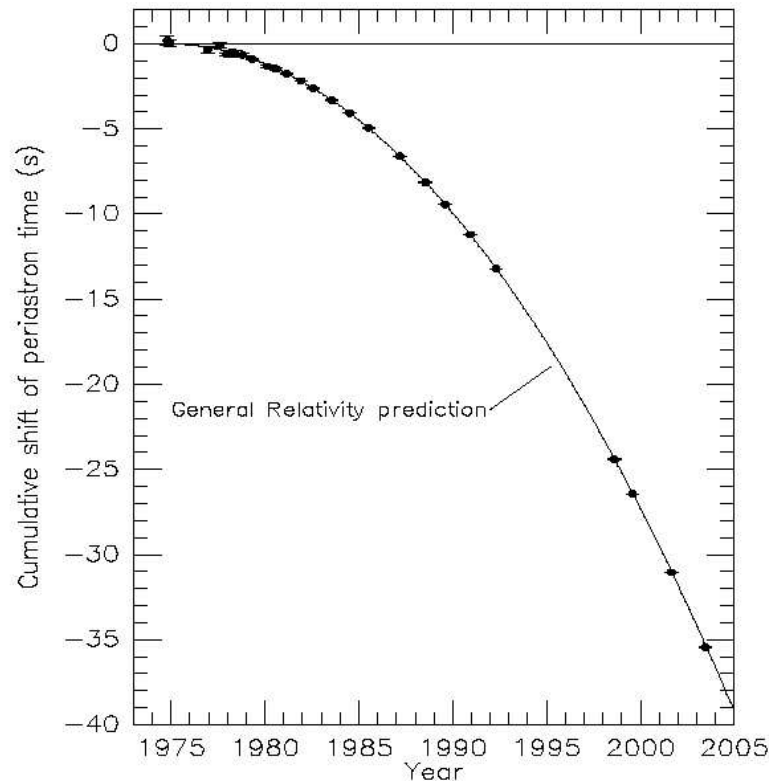
* **pulsar period:** $P_{PSR} = 59 \text{ ms}$

* **eccentricity:** $e = 0.617$

*periastron advance: $\dot{\omega} = 4.227^{\circ}y^{-1}$
($\gg (\dot{\omega})_{Mer} = 0.43''y^{-1}$)

*distance: ~ 10 kpc

*orbital decay: agreement with GR-prediction: **0.21 %** (2004)



*inspiral time: $\tau_{insp} = 2.97 \cdot 10^8$ years !

→ final coalescence !

- most relat. system: PSR J0737-3039A+B (Burgay et al.2003)

- * both are pulsars !

- * masses: 1.337 and 1.250 $M_{\odot} \pm 0.005 M_{\odot}$

- * orbital period: $\tau_{orb} = 2.4 \text{ h}$

- * pulsar periods: $P_A = 22.7 \text{ ms} \ \& \ P_B = 2.8 \text{ s}$

- * excentricity: $e = 0.09$

- * periastron advance: $\dot{\omega} = 17 \text{ }^{\circ}\text{y}^{-1}$

- * distance: $\sim 600 - 1000 \text{ pc}$

- * coalescence: in $8.5 \cdot 10^7 \text{ years}$

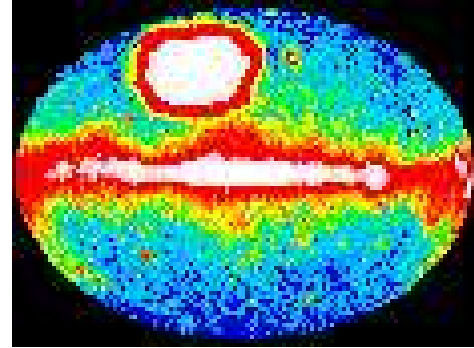
⇒ merger rate : $R_{\text{DNS}} \sim 10^{-4} \text{ (year galaxy)}^{-1}$

⇒ ground-based gravitational detectors (LIGO, GEO600, TAMA ...) should observe

⇒ one DNS merger event every few years !

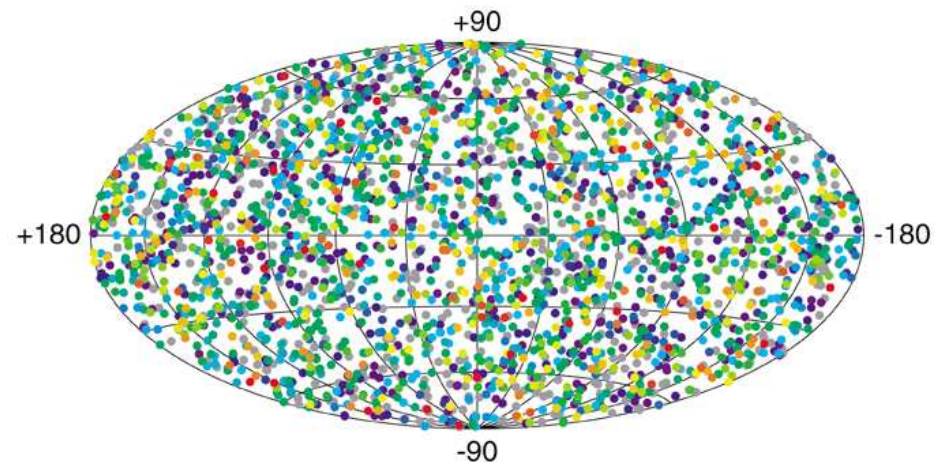
I.3 Gamma Ray Bursts (GRBs)

- * accidental discovery by satellites in the sixties



- * gamma ray sky:
- * rate: $\sim 1/\text{day}$ (BATSE)
- * isotropic distribution

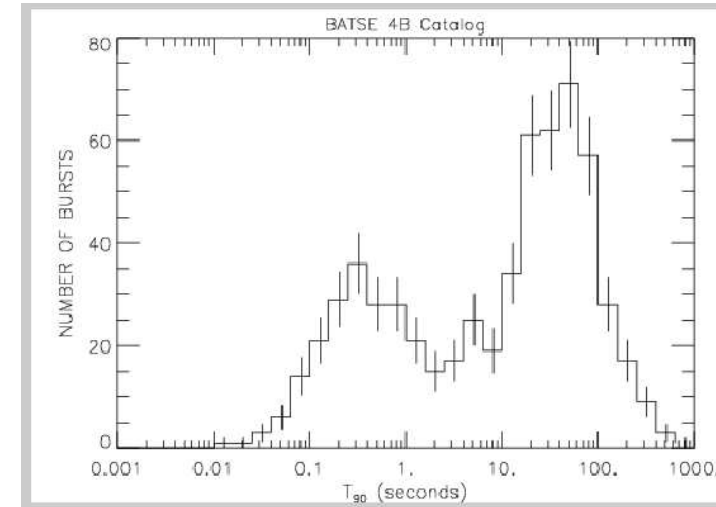
2704 BATSE Gamma-Ray Bursts



(ii) Duration

bimodal: (α) short Bursts ~ 0.2 s
compact binary mergers

(β) long Bursts ~ 30 s
collapsing stars (“collapsars”)



(iii) “standard” central engine:

BH + accretion disk

(iv) Most Popular Mechanisms (to produce beamed, relativistic outflow)

(α) Magnetohydrodynamics (MHD)

(β) $\nu_i + \bar{\nu}_i \rightarrow e^+ + e^-$

II. Modeling compact binary systems:

- intrinsically 3D process → numerical modelling
- high sound velocities: $c_s \sim 0.3c$

Courant-Friedrichs-Lewy stability criterion:

$$\Delta t < \frac{\Delta x}{c_s} = 10^{-6} \text{ s} \left(\frac{\Delta x}{1 \text{ km}} \right) \left(\frac{0.3c}{c_s} \right) \quad \text{short time steps !}$$

- > can only simulate “short physical time scales”
- > need powerful computer

- ideally:

- * “the true equation of state”
- * 3D neutrino transport
- * 3D general relativistic magneto-hydrodynamics (with time variable metrics)
- * numerical resolution of all relevant scales
- * ...

(current) Model ingredients

(current) Model ingredients

(i) Hydrodynamics:

3D Lagrangian particle scheme (SPH),
fully parallelized (Rosswog & Davies 2002)

(current) Model ingredients

(i) Hydrodynamics:

3D Lagrangian particle scheme (SPH),
fully parallelized (Rosswog & Davies 2002)

(ii) Gravity:

Newtonian self-gravity (e.g. Benz 1990) +
gravitational wave backreaction forces

(quadrupole approximation; for details Rosswog et al. 2002)

(current) Model ingredients

(i) Hydrodynamics:

3D Lagrangian particle scheme (SPH),
fully parallelized (Rosswog & Davies 2002)

(ii) Gravity:

Newtonian self-gravity (e.g. Benz 1990) +
gravitational wave backreaction forces

(quadrupole approximation; for details Rosswog et al. 2002)

(iii) Nuclear physics:

temp. and compos. dependent, nuclear EOS
(Relativistic Mean Field theory; Shen et al. 1998a,b)

(current) Model ingredients

(i) Hydrodynamics:

3D Lagrangian particle scheme (SPH),
fully parallelized (Rosswog & Davies 2002)

(ii) Gravity:

Newtonian self-gravity (e.g. Benz 1990) +
gravitational wave backreaction forces

(quadrupole approximation; for details Rosswog et al. 2002)

(iii) Nuclear physics:

temp. and compos. dependent, nuclear EOS
(Relativistic Mean Field theory; Shen et al. 1998a,b)

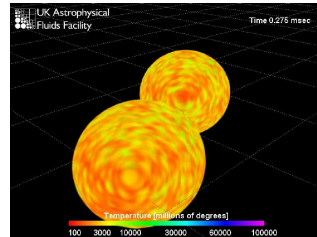
(iv) Neutrino physics:

multi-flavour neutrino leakage scheme
(Rosswog & Liebendörfer 2003)

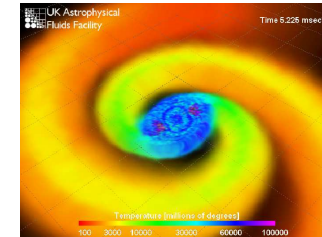
III. Results: the role of the EOS

III.1 Neutron Star Binaries

- Morphology:



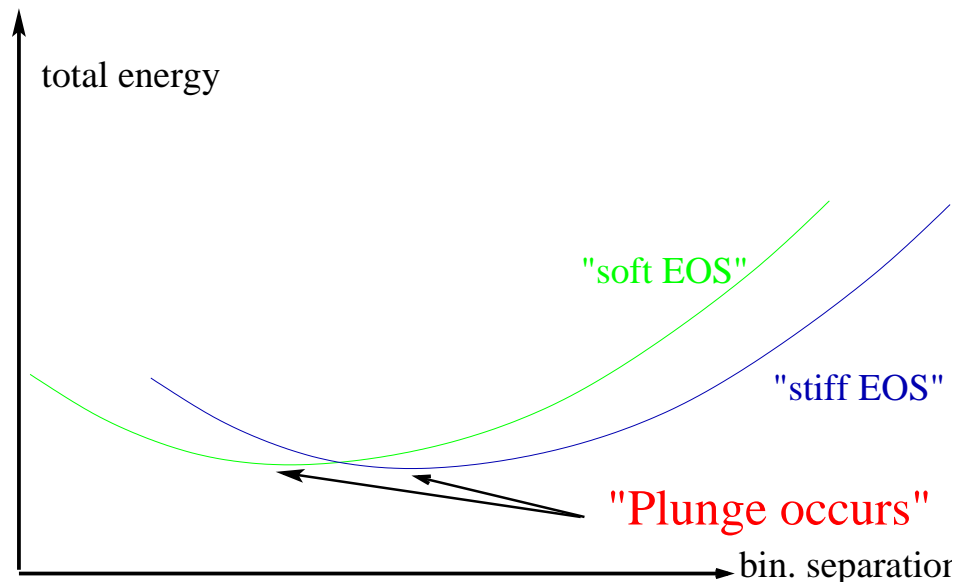
surf. temp.



orbital plane

- inspiral:

“plunging instability” (purely Newtonian discussion)

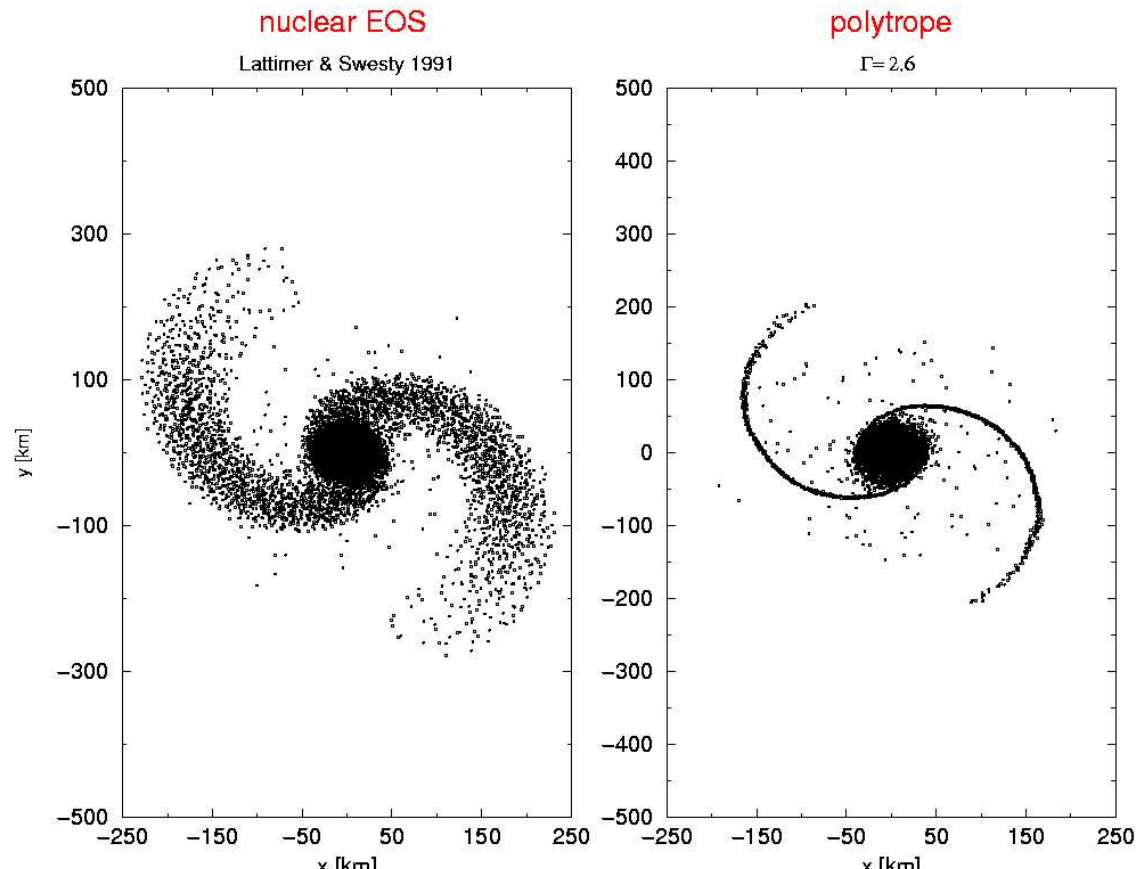


⇒ system becomes **dynamically** unstable

- “stiff EOS”: instability at large separations
- “soft EOS”: even “contact config’s” possible ($\Gamma \leq 2$)

⇒ **visible in GW-signal**

- spiral arms:



- central object: differentially rotating, $\sim 2.5 M_{\odot}$
 - “soft EOS”: immediate collapse to black hole
 - “stiff EOS”: metastable “super neutron star” possible

(complicated by time scale to remove diff. rot.; viscosity, GW, magnetic fields....)

- neutron-rich ejecta:
 - “stiff EOS”: $\sim 0.01 M_{\odot}$ ejected per event
(Shen et al. EOS 1998)
 - “soft EOS”: no resolvable mass loss
($\Gamma = 2$ -polytrope)

III.2 Neutron Star Black Hole Binaries

- supposed to yield “standard GRB central engine”:
BH + massive torus

- complex accretion dynamics (sensitive to EOS !)

determined by

- Mass transfer \Rightarrow increase orbital separation

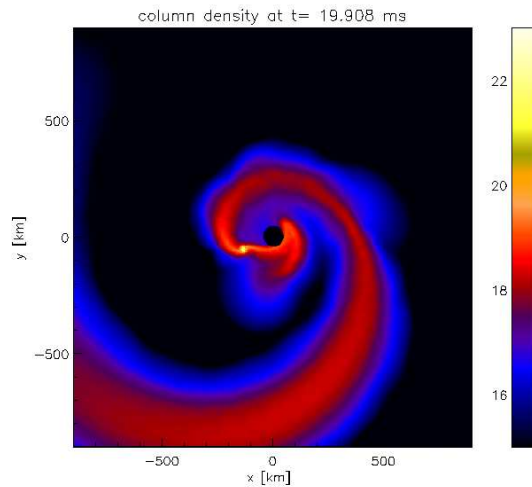
- GW emission \Rightarrow reduce orbital separation

- Reaction of NS to mass loss:

$$\frac{dR_{ns}}{dM} < 0 \Rightarrow \text{“ns expands”} \Rightarrow \text{increase mass transfer}$$

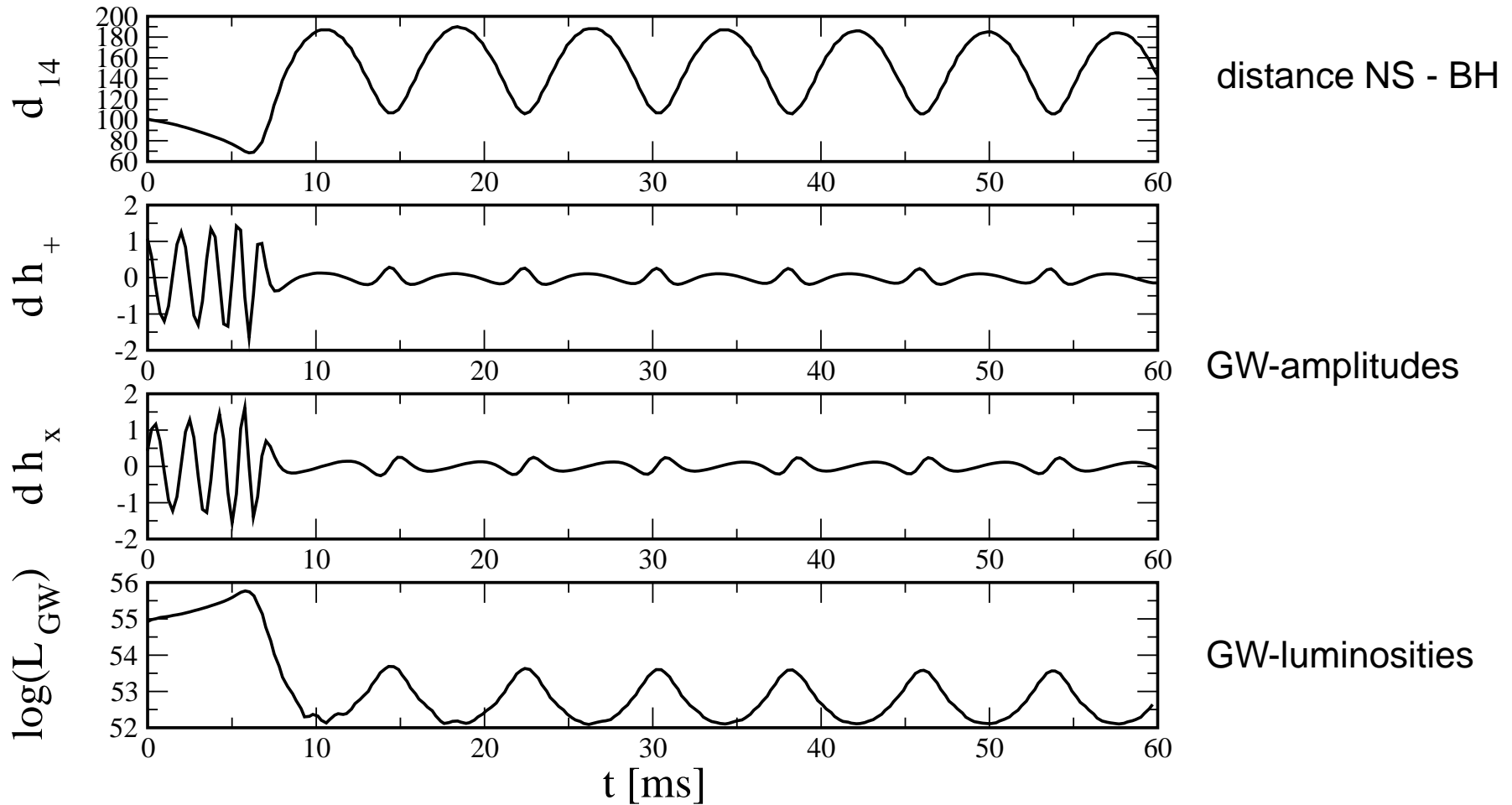
$$\frac{dR_{ns}}{dM} > 0 \Rightarrow \text{“ns shrinks”} \Rightarrow \text{decrease mass transfer}$$

Accretion Dynamics in Neutron Star Black Hole Binaries:



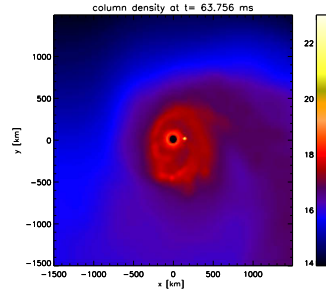
column density,
Newtonian gravity,
corotation,
mass ratio $q = 0.1$

* for shown run ($q=0.1$, corotation):



\Rightarrow “survival of mini-NS visible in GW-signal”

Implications for GRBs



disks NS-BH

NS-NS binaries

- masses, $< 10^{-2} M_{\odot}$ \Leftrightarrow $\sim 0.2 M_{\odot}$
- densities [g cm $^{-3}$] $10^8 < \rho < 10^{11}$ \Leftrightarrow $10^{11} < \rho < 10^{13}$
- temperatures, ~ 2.5 MeV \Leftrightarrow ~ 4 MeV
- neutrino emission, $\sim 10^{52}$ erg/s \Leftrightarrow $\sim 2 \cdot 10^{53}$ erg/s

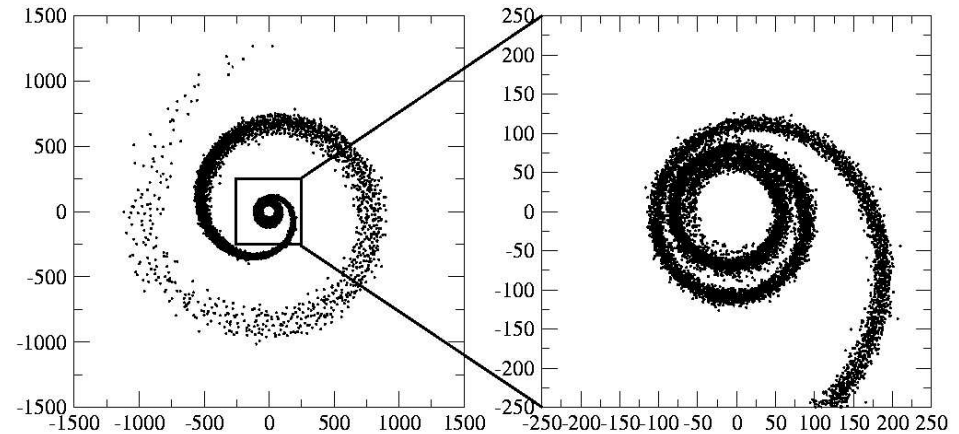
therefore:

- inefficient neutrino annihilation ($Q_{\nu\bar{\nu}} \propto L_{\nu_i} L_{\bar{\nu}_i}$)
- difficult to anchor strong magnetic fields in disk

\Rightarrow “pessimistic prospects for GRBs”

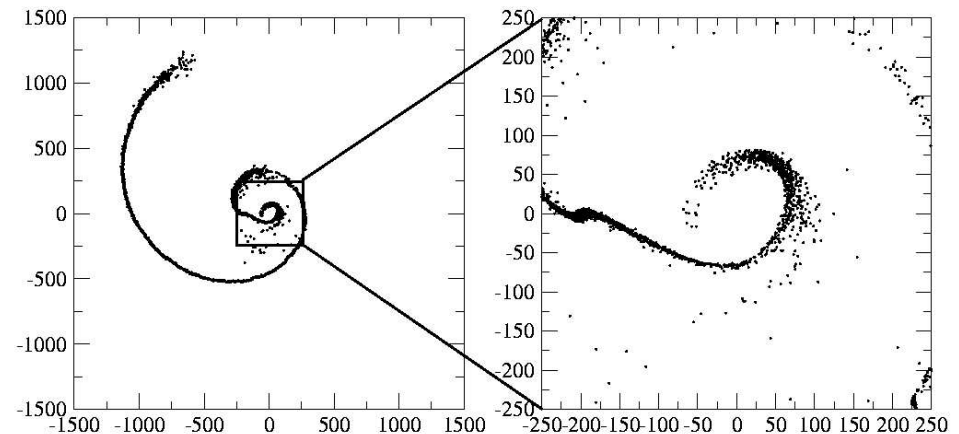
Sensitivity to EOS

* soft EOS (polytrope, $\Gamma = 2$)



⇒ complete disruption, massive disk

* stiff EOS (polytrope, $\Gamma = 3$)



⇒ mini-NS, low-mass disk

Summary

Neutron star mergers:

various aspects sensitive to EOS:

- morphology: spiral arms etc.
- stability central object
- amount neutron-rich ejecta
- neutrino emission
- ...

Neutron star black hole mergers:

- extremely complex accretion dynamics
 - dynamics very sensitive to EOS
 - “hard” EOS used (relativistic mean field, Shen et al. 1998a,b):
 - “mini-neutron” star survives
 - difficult to form accretion disk
- ⇒ good news for GW-detection,
bad news for GRBs

The Astrophysicist's wish list

for the EOS

- temperature dependence $0 < T < \sim 100 \text{ MeV}$
- NO β -equilibrium $0 < Y_e < \sim 0.5$
- large density range $\sim 10^3 < \rho < \sim 10^{15} \text{ g/cm}^3$